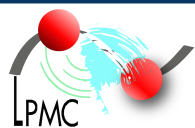


HBT type measurements in quantum optics experiments

Sébastien Tanzilli

Team “Quantum Information with Light & Matter”

UMR 7336



**Laboratoire
Physique de la Matière Condensée**



HBT Workshop, Observatoire de la Côte d'Azur, Nice, May 12-13 2014



Outline

1. Introduction

1. Standard HBT setup & classical description of $g^{(2)}(0)$
2. HBT in quantum optics / experimental tools
3. Quantum description of $g^{(2)}(0)$
4. The context of quantum communication

2. HBT for characterizing single photon sources (SPS)

3. HBT for characterizing photon pair sources

I.1 The HBT setup: a bright idea for Astrophysics

▶ HBT - original method

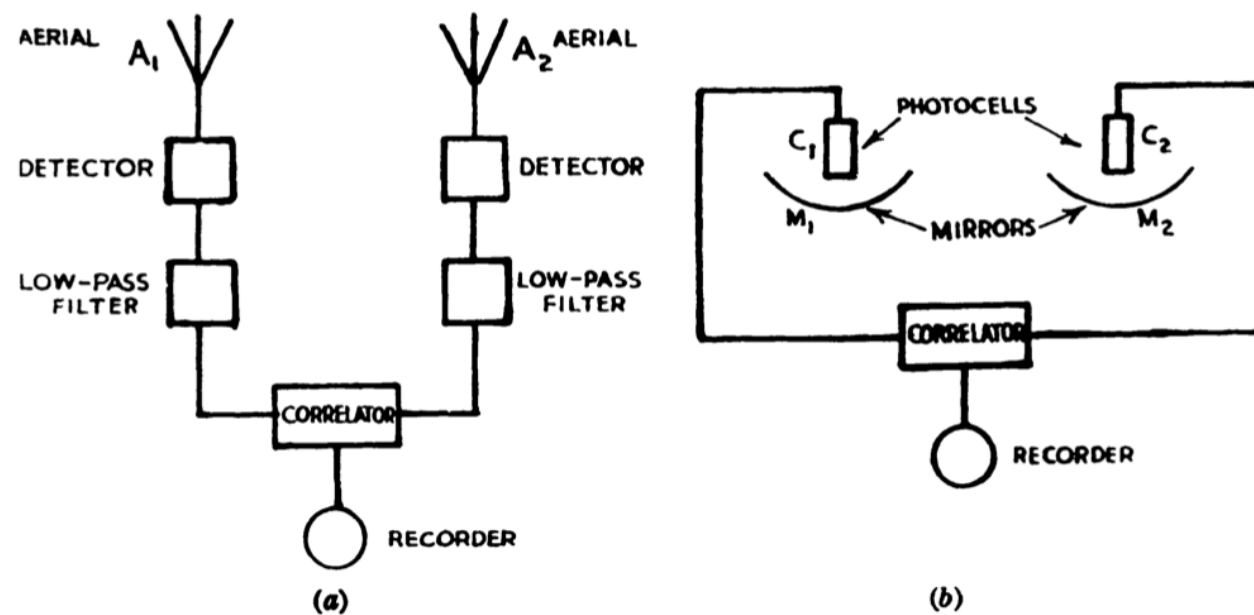


Fig. 1. A new type of radio interferometer (a), together with its analogue (b) at optical wave-lengths

I.I The HBT setup: a bright idea for Astrophysics

▶ HBT - original method

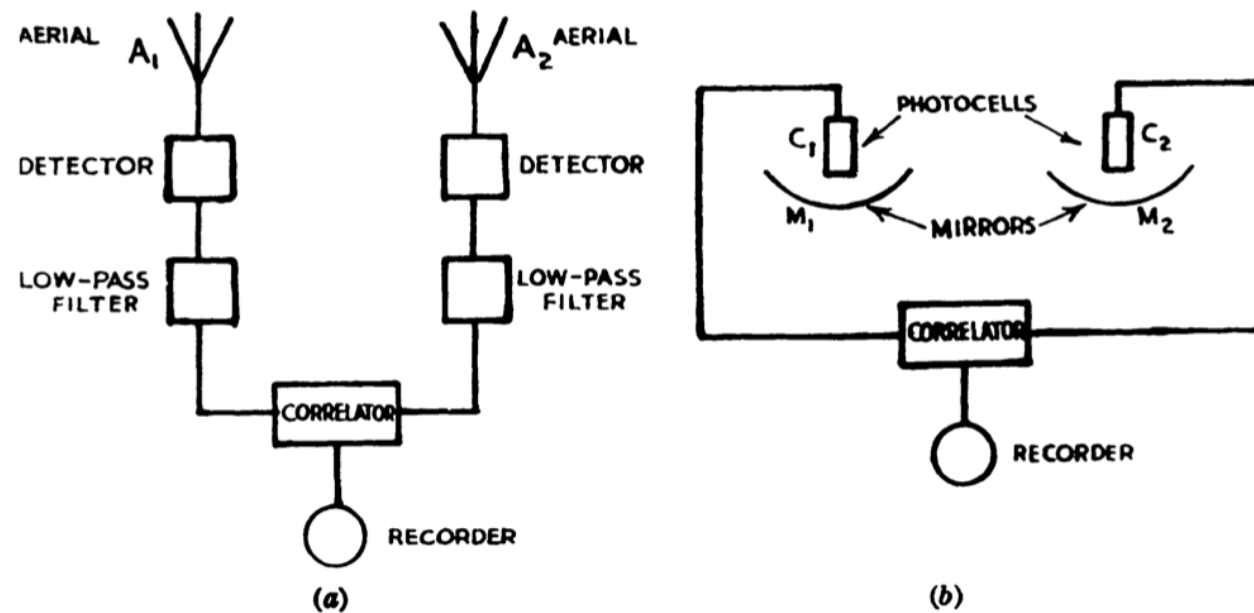


Fig. 1. A new type of radio interferometer (a), together with its analogue (b) at optical wave-lengths

▶ What for ?

Measurement of the spatial coherence area → angular diameter of a star

I.1 The HBT setup: a bright idea for Astrophysics

▶ HBT - original method

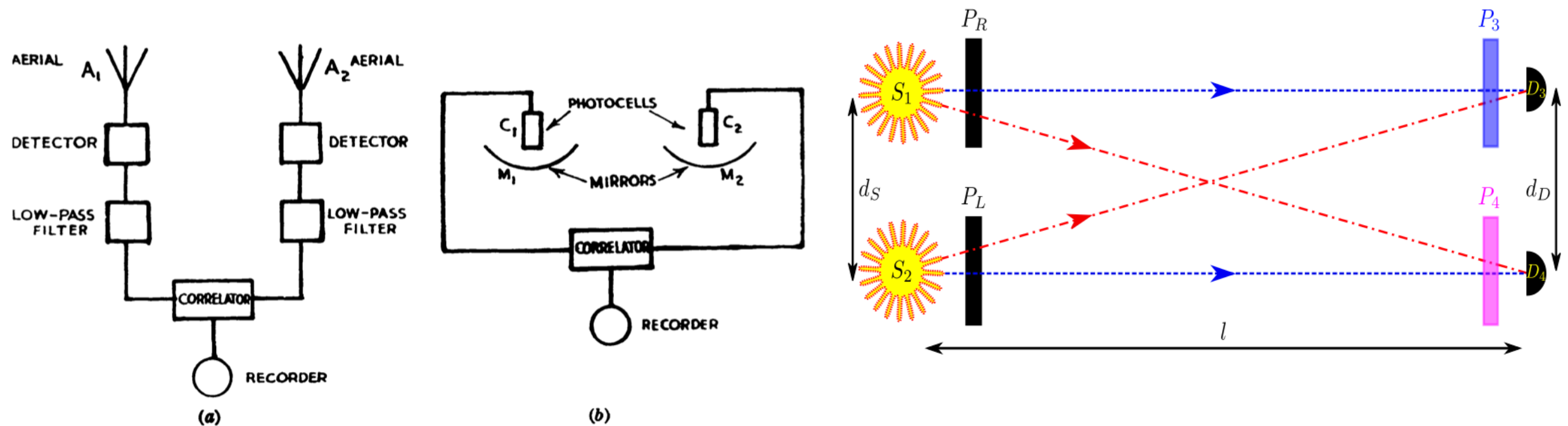


Fig. 1. A new type of radio interferometer (a), together with its analogue (b) at optical wave-lengths

▶ What for ?

Measurement of the spatial coherence area → angular diameter of a star

Interferometer-like configuration insensitive to atmospheric fluctuations !

1.1 Classical description of the autocorrelation function

▶ Classical intensity correlation in the time domain

$$g^{(2)}(t_1, t_2) = \frac{\langle E^*(t_1)E^*(t_2)E(t_2)E(t_1) \rangle}{\sqrt{\langle |E(t_1)|^2 \rangle \langle |E(t_2)|^2 \rangle}} = \frac{\langle I(t_1)I(t_2) \rangle}{\langle I(t_1) \rangle \langle I(t_2) \rangle}$$

with $I(t_i) = E^*(t_i)E(t_i) = |E(t_i)|^2$

I.1 Classical description of the autocorrelation function

▶ Classical intensity correlation in the time domain

$$g^{(2)}(t_1, t_2) = \frac{\langle E^*(t_1)E^*(t_2)E(t_2)E(t_1) \rangle}{\sqrt{\langle |E(t_1)|^2 \rangle \langle |E(t_2)|^2 \rangle}} = \frac{\langle I(t_1)I(t_2) \rangle}{\langle I(t_1) \rangle \langle I(t_2) \rangle}$$

with $I(t_i) = E^*(t_i)E(t_i) = |E(t_i)|^2$

▶ Properties

I.1 Classical description of the autocorrelation function

▶ Classical intensity correlation in the time domain

$$g^{(2)}(t_1, t_2) = \frac{\langle E^*(t_1)E^*(t_2)E(t_2)E(t_1) \rangle}{\sqrt{\langle |E(t_1)|^2 \rangle \langle |E(t_2)|^2 \rangle}} = \frac{\langle I(t_1)I(t_2) \rangle}{\langle I(t_1) \rangle \langle I(t_2) \rangle}$$

with $I(t_i) = E^*(t_i)E(t_i) = |E(t_i)|^2$

▶ Properties

Cauchy-Schwartz inequality $\mapsto \langle I(t_1)I(t_2) \rangle^2 \leq \langle I^2(t_1) \rangle \langle I^2(t_2) \rangle$

I.1 Classical description of the autocorrelation function

▶ Classical intensity correlation in the time domain

$$g^{(2)}(t_1, t_2) = \frac{\langle E^*(t_1)E^*(t_2)E(t_2)E(t_1) \rangle}{\sqrt{\langle |E(t_1)|^2 \rangle \langle |E(t_2)|^2 \rangle}} = \frac{\langle I(t_1)I(t_2) \rangle}{\langle I(t_1) \rangle \langle I(t_2) \rangle}$$

with $I(t_i) = E^*(t_i)E(t_i) = |E(t_i)|^2$

▶ Properties

Cauchy-Schwartz inequality $\mapsto \langle I(t_1)I(t_2) \rangle^2 \leq \langle I^2(t_1) \rangle \langle I^2(t_2) \rangle$

Stationary regime & $\tau = t_2 - t_1 \mapsto \langle I(t)I(t + \tau) \rangle^2 \leq \langle I^2(t) \rangle^2$

1.1 Classical description of the autocorrelation function

▶ Classical intensity correlation in the time domain

$$g^{(2)}(t_1, t_2) = \frac{\langle E^*(t_1)E^*(t_2)E(t_2)E(t_1) \rangle}{\sqrt{\langle |E(t_1)|^2 \rangle \langle |E(t_2)|^2 \rangle}} = \frac{\langle I(t_1)I(t_2) \rangle}{\langle I(t_1) \rangle \langle I(t_2) \rangle}$$

with $I(t_i) = E^*(t_i)E(t_i) = |E(t_i)|^2$

▶ Properties

Cauchy-Schwartz inequality $\mapsto \langle I(t_1)I(t_2) \rangle^2 \leq \langle I^2(t_1) \rangle \langle I^2(t_2) \rangle$

Stationary regime & $\tau = t_2 - t_1 \mapsto \langle I(t)I(t + \tau) \rangle^2 \leq \langle I^2(t) \rangle^2$

$g^{(2)}(0)$ is max in 0

→ so-called photon bunching

$$\mapsto g^{(2)}(\tau) \leq g^{(2)}(0)$$

1.1 Classical description of the autocorrelation function

▶ Classical intensity correlation in the time domain

$$g^{(2)}(t_1, t_2) = \frac{\langle E^*(t_1)E^*(t_2)E(t_2)E(t_1) \rangle}{\sqrt{\langle |E(t_1)|^2 \rangle \langle |E(t_2)|^2 \rangle}} = \frac{\langle I(t_1)I(t_2) \rangle}{\langle I(t_1) \rangle \langle I(t_2) \rangle}$$

with $I(t_i) = E^*(t_i)E(t_i) = |E(t_i)|^2$

▶ Properties

Cauchy-Schwartz inequality $\mapsto \langle I(t_1)I(t_2) \rangle^2 \leq \langle I^2(t_1) \rangle \langle I^2(t_2) \rangle$

Stationary regime & $\tau = t_2 - t_1 \mapsto \langle I(t)I(t + \tau) \rangle^2 \leq \langle I^2(t) \rangle^2$

$g^{(2)}(0)$ is max in 0

→ so-called photon bunching

$$\mapsto g^{(2)}(\tau) \leq g^{(2)}(0)$$

Remarking $\langle I^2(t) \rangle \geq \langle I(t) \rangle^2$

1.1 Classical description of the autocorrelation function

▶ Classical intensity correlation in the time domain

$$g^{(2)}(t_1, t_2) = \frac{\langle E^*(t_1)E^*(t_2)E(t_2)E(t_1) \rangle}{\sqrt{\langle |E(t_1)|^2 \rangle \langle |E(t_2)|^2 \rangle}} = \frac{\langle I(t_1)I(t_2) \rangle}{\langle I(t_1) \rangle \langle I(t_2) \rangle}$$

with $I(t_i) = E^*(t_i)E(t_i) = |E(t_i)|^2$

▶ Properties

Cauchy-Schwartz inequality $\mapsto \langle I(t_1)I(t_2) \rangle^2 \leq \langle I^2(t_1) \rangle \langle I^2(t_2) \rangle$

Stationary regime & $\tau = t_2 - t_1 \mapsto \langle I(t)I(t + \tau) \rangle^2 \leq \langle I^2(t) \rangle^2$

$g^{(2)}(0)$ is max in 0

→ so-called photon bunching

$$\mapsto g^{(2)}(\tau) \leq g^{(2)}(0)$$

Remarking $\langle I^2(t) \rangle \geq \langle I(t) \rangle^2$

$$\mapsto g^{(2)}(0) \geq 1$$

Classical $g^{(2)}(0)$ is always greater than 1 !

1.1 Classical description of the autocorrelation function

- ▶ More properties & practical examples

1.1 Classical description of the autocorrelation function

▶ **More properties & practical examples**

$g^{(2)}(0) \geq 1$ → no photon antibunching expected from this description

1.1 Classical description of the autocorrelation function

▶ More properties & practical examples

$g^{(2)}(0) \geq 1$ → no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

1.1 Classical description of the autocorrelation function

▶ More properties & practical examples

$g^{(2)}(0) \geq 1$ → no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$g^{(2)}(0) = 1$ → no correlation between subsequent photo-detections as is the case for any standard laser

1.1 Classical description of the autocorrelation function

► More properties & practical examples

$g^{(2)}(0) \geq 1$ → no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$g^{(2)}(0) = 1$ → no correlation between subsequent photo-detections as is the case for any standard laser

For small $\bar{n} \ll 1$ → $P_2 = \frac{P_1^2}{2}$

1.1 Classical description of the autocorrelation function

► More properties & practical examples

$g^{(2)}(0) \geq 1$ → no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$g^{(2)}(0) = 1$ → no correlation between subsequent photo-detections as is the case for any standard laser

For small $\bar{n} \ll 1$ → $P_2 = \frac{P_1^2}{2}$

The case of thermal light $P_T(n, \bar{n}) = \frac{1}{(1 + \bar{n}) \left(1 + \frac{1}{\bar{n}}\right)^n}$

1.1 Classical description of the autocorrelation function

► More properties & practical examples

$g^{(2)}(0) \geq 1$ → no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$g^{(2)}(0) = 1$ → no correlation between subsequent photo-detections as is the case for any standard laser

For small $\bar{n} \ll 1$ → $P_2 = \frac{P_1^2}{2}$

The case of thermal light $P_T(n, \bar{n}) = \frac{1}{(1 + \bar{n}) \left(1 + \frac{1}{\bar{n}}\right)^n}$

$$g^{(2)}(\tau) = 1 + |g^{(1)}(\tau)|^2$$

Corr. function
of the EM field

1.1 Classical description of the autocorrelation function

▶ More properties & practical examples

$g^{(2)}(0) \geq 1 \rightarrow$ no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$g^{(2)}(0) = 1 \rightarrow$ no correlation between subsequent photo-detections as is the case for any standard laser

For small $\bar{n} \ll 1 \rightarrow P_2 = \frac{P_1^2}{2}$

The case of thermal light $P_T(n, \bar{n}) = \frac{1}{(1 + \bar{n}) \left(1 + \frac{1}{\bar{n}}\right)^n}$

$g^{(2)}(\tau) = 1 + |g^{(1)}(\tau)|^2$ with $g^{(1)}(\tau) = e^{-\gamma^2 \tau^2}$ Doppler broadening
 Corr. function of the EM field

1.1 Classical description of the autocorrelation function

► More properties & practical examples

$g^{(2)}(0) \geq 1 \rightarrow$ no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$g^{(2)}(0) = 1 \rightarrow$ no correlation between subsequent photo-detections as is the case for any standard laser

For small $\bar{n} \ll 1 \rightarrow P_2 = \frac{P_1^2}{2}$

The case of thermal light $P_T(n, \bar{n}) = \frac{1}{(1 + \bar{n}) \left(1 + \frac{1}{\bar{n}}\right)^n}$

$g^{(2)}(\tau) = 1 + |g^{(1)}(\tau)|^2$
 Corr. function of the EM field

with $g^{(1)}(\tau) = e^{-\gamma^2 \tau^2}$ Doppler broadening

$$\mapsto g^{(2)}(0) = 2$$

1.1 Classical description of the autocorrelation function

► More properties & practical examples

$g^{(2)}(0) \geq 1$ → no photon antibunching expected from this description

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$g^{(2)}(0) = 1$ → no correlation between subsequent photo-detections as is the case for any standard laser

For small $\bar{n} \ll 1$ → $P_2 = \frac{P_1^2}{2}$

The case of thermal light $P_T(n, \bar{n}) = \frac{1}{(1 + \bar{n}) \left(1 + \frac{1}{\bar{n}}\right)^n}$

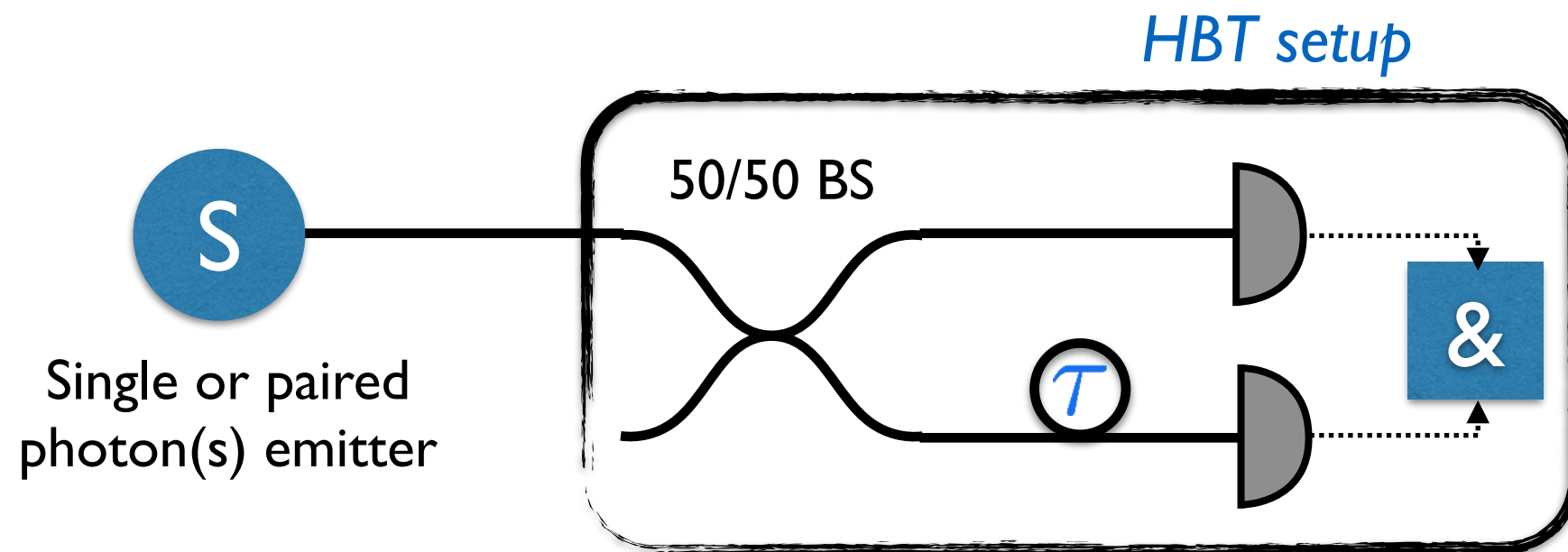
$g^{(2)}(\tau) = 1 + |g^{(1)}(\tau)|^2$ with $g^{(1)}(\tau) = e^{-\gamma^2 \tau^2}$ Doppler broadening

Corr. function of the EM field

$$\mapsto g^{(2)}(0) = 2$$

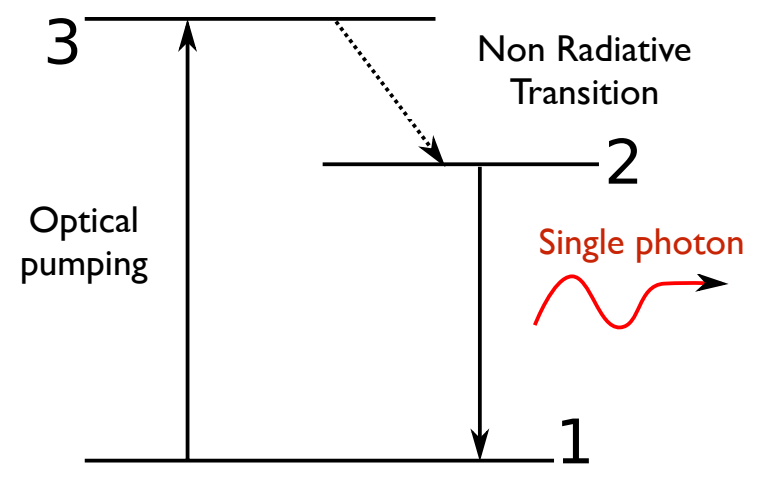
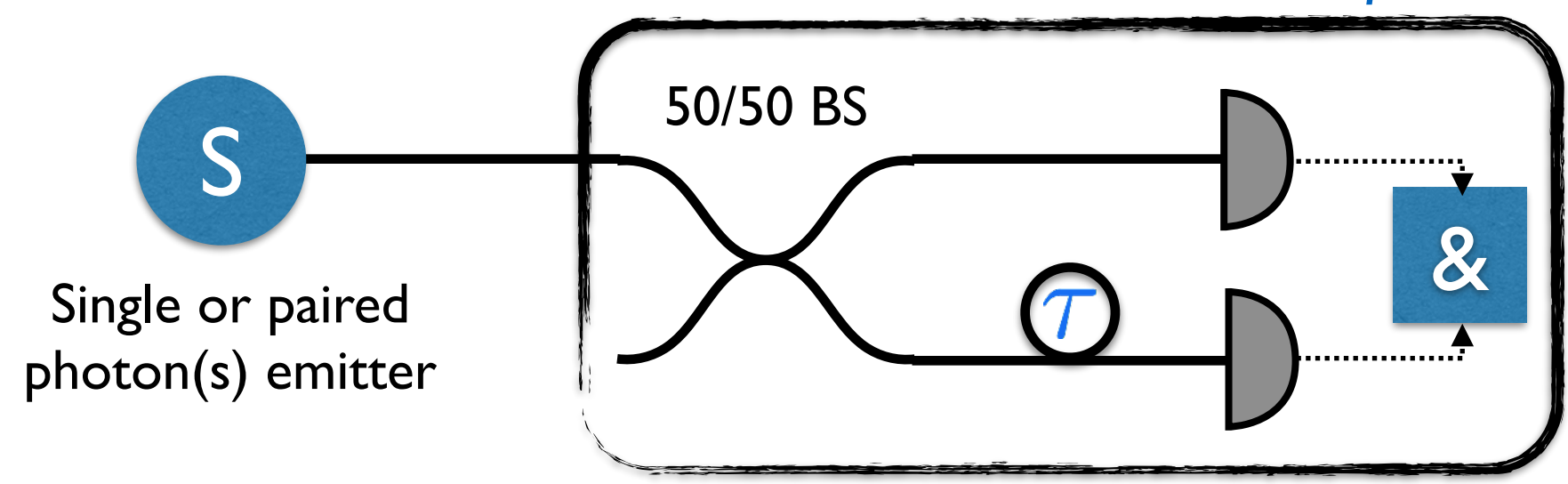
For small $\bar{n} \ll 1$ → $P_2 = P_1^2$

1.2 Experimental tools for HBT in quantum optics

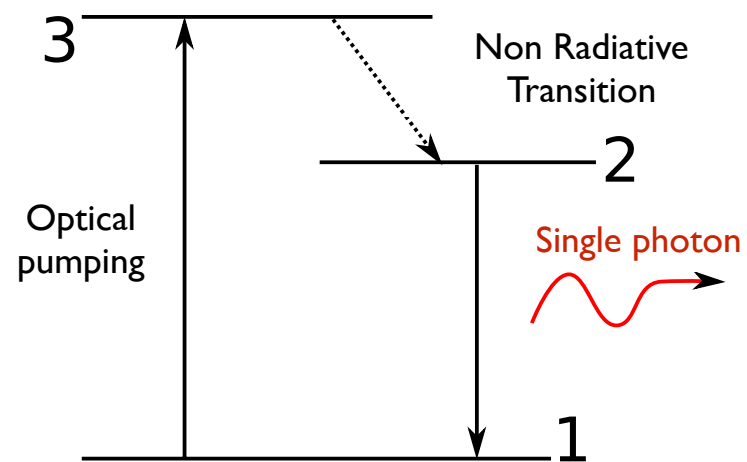
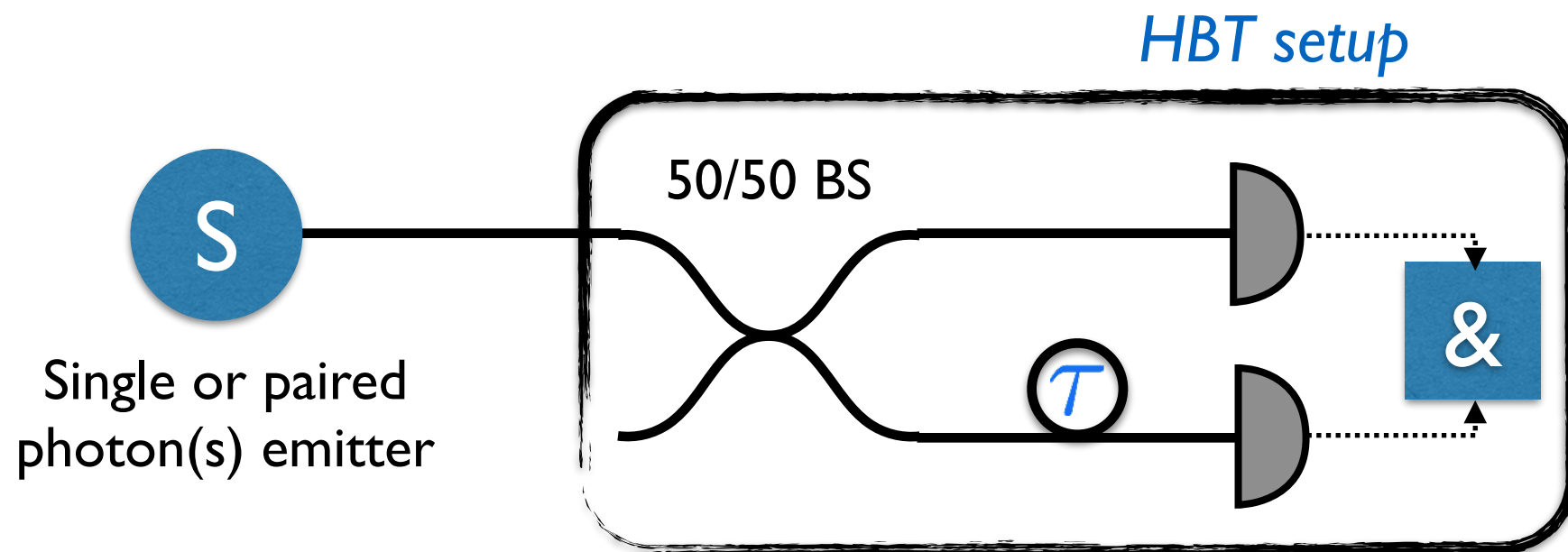


1.2 Experimental tools for HBT in quantum optics

HBT setup



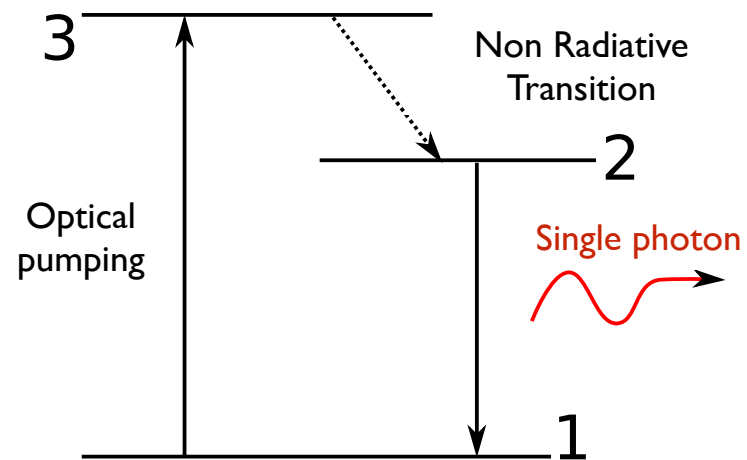
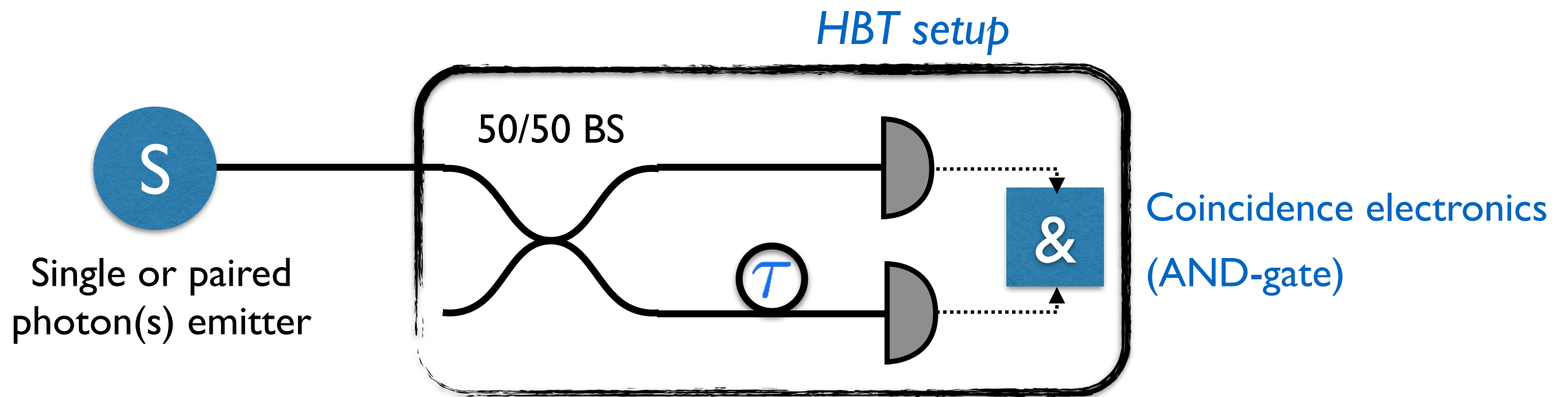
1.2 Experimental tools for HBT in quantum optics



Relevant features

Efficiency	η	“Poissonization”
Dark counts	p_{dc}	noise level
Timing jitter	τ_d	resolution

1.2 Experimental tools for HBT in quantum optics



Relevant features

Efficiency	η	“Poissonization”
Dark counts	p_{dc}	noise level
Timing jitter	τ_d	resolution

I.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

1.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

$$g^{(2)}(\tau) = \frac{\langle E^\dagger(t) E^\dagger(t + \tau) E(t + \tau) E(t) \rangle}{\langle E^\dagger(t) E(t) \rangle \langle E^\dagger(t + \tau) E(t + \tau) \rangle}$$

1.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

$$g^{(2)}(\tau) = \frac{\langle E^\dagger(t) E^\dagger(t + \tau) E(t + \tau) E(t) \rangle}{\langle E^\dagger(t) E(t) \rangle \langle E^\dagger(t + \tau) E(t + \tau) \rangle}$$

$$\mapsto g^{(2)}(\tau) = \frac{\langle : I(t + \tau) I(t) : \rangle}{\langle I(t) \rangle^2}$$

I.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

$$g^{(2)}(\tau) = \frac{\langle E^\dagger(t) E^\dagger(t + \tau) E(t + \tau) E(t) \rangle}{\langle E^\dagger(t) E(t) \rangle \langle E^\dagger(t + \tau) E(t + \tau) \rangle}$$

$$\mapsto g^{(2)}(\tau) = \frac{\langle : I(t + \tau) I(t) : \rangle}{\langle I(t) \rangle^2}$$

- ▶ **Take care !**

I.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

$$g^{(2)}(\tau) = \frac{\langle E^\dagger(t) E^\dagger(t + \tau) E(t + \tau) E(t) \rangle}{\langle E^\dagger(t) E(t) \rangle \langle E^\dagger(t + \tau) E(t + \tau) \rangle}$$

$$\mapsto g^{(2)}(\tau) = \frac{\langle : I(t + \tau) I(t) : \rangle}{\langle I(t) \rangle^2}$$

- ▶ **Take care !**

$\langle : : \rangle \rightarrow$ the operators are normally ordered

I.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

$$g^{(2)}(\tau) = \frac{\langle E^\dagger(t) E^\dagger(t + \tau) E(t + \tau) E(t) \rangle}{\langle E^\dagger(t) E(t) \rangle \langle E^\dagger(t + \tau) E(t + \tau) \rangle}$$

$$\mapsto g^{(2)}(\tau) = \frac{\langle : I(t + \tau) I(t) : \rangle}{\langle I(t) \rangle^2}$$

- ▶ **Take care !**

$\langle : : \rangle \rightarrow$ the operators are normally ordered

$E^\dagger(t)$ and $E^\dagger(\tau)$ do not commute \rightarrow Cauchy-Schwartz cannot be applied anymore
 $\rightarrow g^{(2)}(0)$ can drop to 0

1.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

$$g^{(2)}(\tau) = \frac{\langle E^\dagger(t) E^\dagger(t + \tau) E(t + \tau) E(t) \rangle}{\langle E^\dagger(t) E(t) \rangle \langle E^\dagger(t + \tau) E(t + \tau) \rangle}$$

$$\mapsto g^{(2)}(\tau) = \frac{\langle : I(t + \tau) I(t) : \rangle}{\langle I(t) \rangle^2}$$

- ▶ **Take care !**

$\langle : : \rangle \rightarrow$ the operators are normally ordered

$E^\dagger(t)$ and $E^\dagger(\tau)$ do not commute \rightarrow Cauchy-Schwartz cannot be applied anymore
 $\rightarrow g^{(2)}(0)$ can drop to 0

- ▶ **Single mode operation**

$$\mapsto g^{(2)}(\tau) = \frac{\langle a^\dagger a^\dagger a a \rangle}{\langle a^\dagger a \rangle^2} = \frac{\langle \hat{n}(\hat{n} - 1) \rangle}{\langle \hat{n} \rangle^2}$$

1.3 Quantum description of the autocorrelation function

- ▶ **Quantum description of the EM field** $E(z, t) = E^\dagger(z, t) + E(z, t)$

$$g^{(2)}(\tau) = \frac{\langle E^\dagger(t) E^\dagger(t + \tau) E(t + \tau) E(t) \rangle}{\langle E^\dagger(t) E(t) \rangle \langle E^\dagger(t + \tau) E(t + \tau) \rangle}$$

$$\mapsto g^{(2)}(\tau) = \frac{\langle : I(t + \tau) I(t) : \rangle}{\langle I(t) \rangle^2}$$

- ▶ **Take care !**

$\langle : : \rangle \rightarrow$ the operators are normally ordered

$E^\dagger(t)$ and $E^\dagger(\tau)$ do not commute \rightarrow Cauchy-Schwartz cannot be applied anymore
 $\rightarrow g^{(2)}(0)$ can drop to 0

- ▶ **Single mode operation**

$$\mapsto g^{(2)}(\tau) = \frac{\langle a^\dagger a^\dagger a a \rangle}{\langle a^\dagger a \rangle^2} = \frac{\langle \hat{n} (\hat{n} - 1) \rangle}{\langle \hat{n} \rangle^2}$$

with $\hat{n} = a^\dagger a$ photon number operator

$\langle \hat{n} \rangle$ mean number of photons in the mode

1.3 Quantum description of the autocorrelation function

- ▶ More properties

I.3 Quantum description of the autocorrelation function

► More properties

The case of coherent-state (poissonnian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

1.3 Quantum description of the autocorrelation function

► More properties

The case of coherent-state (poissonnian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$$a|\alpha\rangle = \alpha|\alpha\rangle \mapsto g^{(2)}(t_i) = 1$$

- no correlation between subsequent photo-detections
- same result as before

I.3 Quantum description of the autocorrelation function

► More properties

The case of coherent-state (poissonnian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$$a|\alpha\rangle = \alpha|\alpha\rangle \mapsto g^{(2)}(t_i) = 1$$

- no correlation between subsequent photo-detections
- same result as before

The case of a single photon emitter → simple 2-level system

I.3 Quantum description of the autocorrelation function

► More properties

The case of coherent-state (poissonian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$$a|\alpha\rangle = \alpha|\alpha\rangle \mapsto g^{(2)}(t_i) = 1$$

- no correlation between subsequent photo-detections
- same result as before

The case of a single photon emitter → simple 2-level system

Rate equations show $g^{(2)}(\tau) \sim 1 - e^{-(\Omega + \Gamma)\tau}$

with Ω the pump rate driving the transition, and Γ^{-1} the lifetime of the excited state

I.3 Quantum description of the autocorrelation function

► More properties

The case of coherent-state (poissonnian) light $P_P(n, \bar{n}) = \frac{\bar{n}^n \cdot e^{-\bar{n}}}{n!}$

$$a|\alpha\rangle = \alpha|\alpha\rangle \mapsto g^{(2)}(t_i) = 1$$

- no correlation between subsequent photo-detections
- same result as before

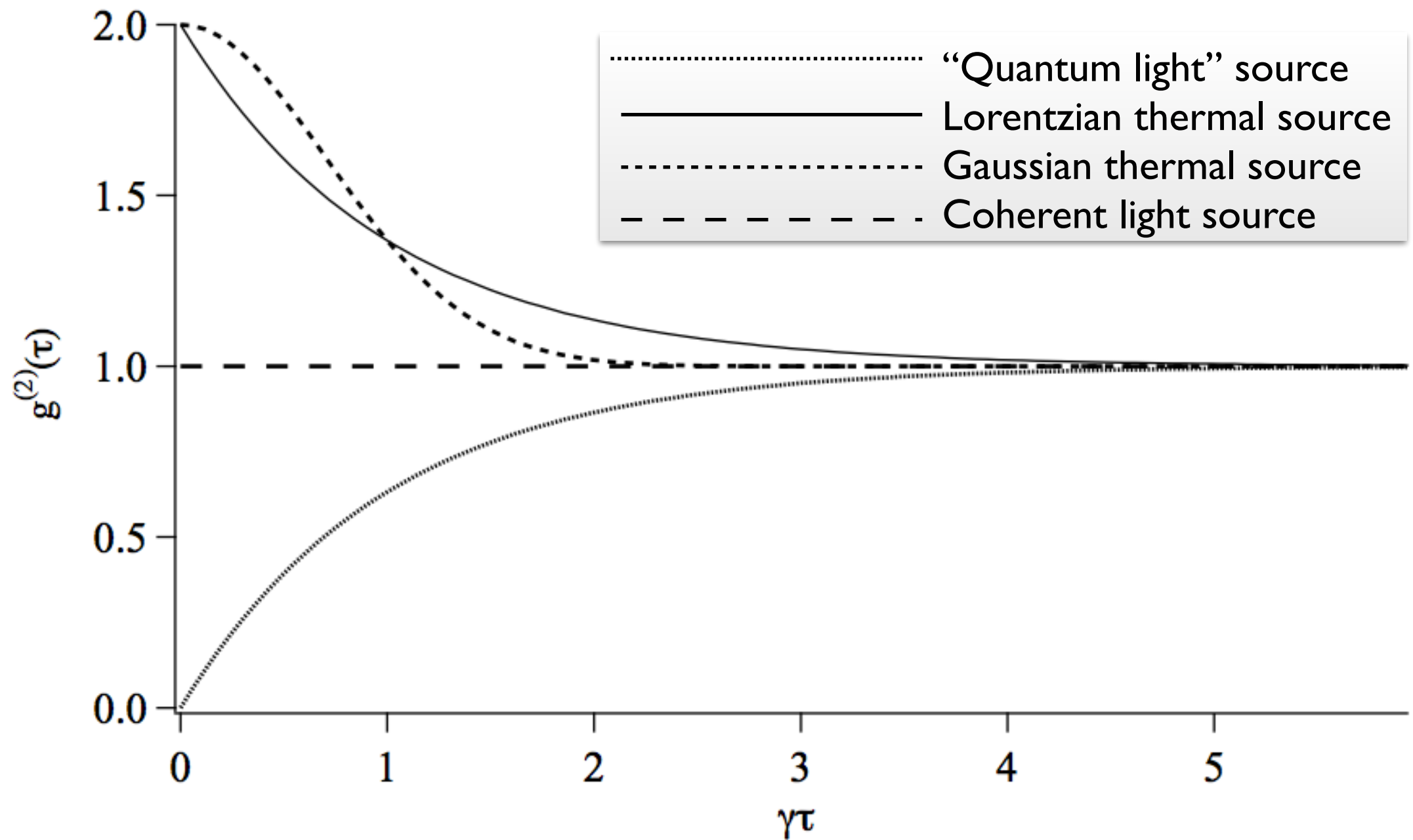
The case of a single photon emitter → simple 2-level system

Rate equations show $g^{(2)}(\tau) \sim 1 - e^{-(\Omega + \Gamma)\tau}$

with Ω the pump rate driving the transition, and Γ^{-1} the lifetime of the excited state

$$\rightarrow \begin{cases} g^{(2)}(0) = 0 \\ g^{(2)}(0) \leq g^{(2)}(\tau) \end{cases} \quad \text{safe !}$$

I.3 Little summary



I.2 HBT in quantum optics

- ▶ The context of today's quantum communication

I.2 HBT in quantum optics

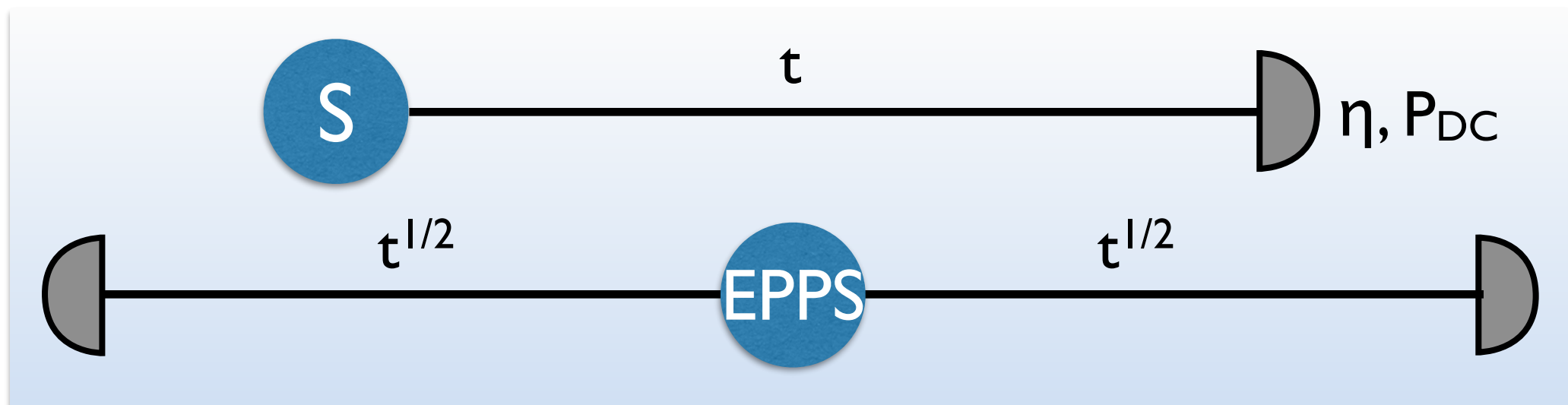
▶ The context of today's quantum communication

Distribution of quantum bits of information
using single photon and entangled photon pair sources

1.2 HBT in quantum optics

► The context of today's quantum communication

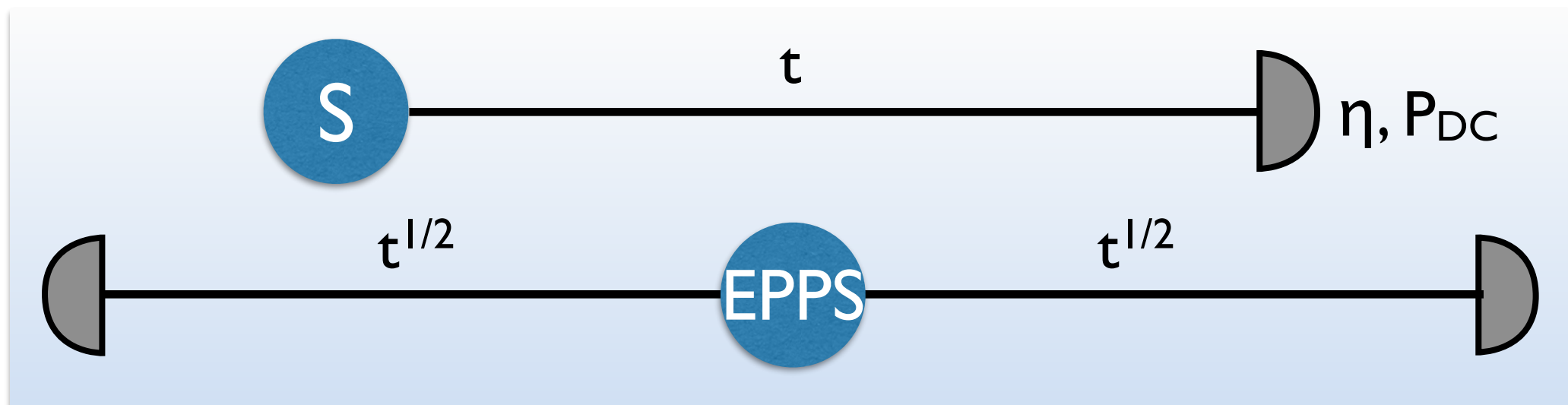
Distribution of quantum bits of information
using single photon and entangled photon pair sources



I.2 HBT in quantum optics

► The context of today's quantum communication

Distribution of quantum bits of information
using single photon and entangled photon pair sources



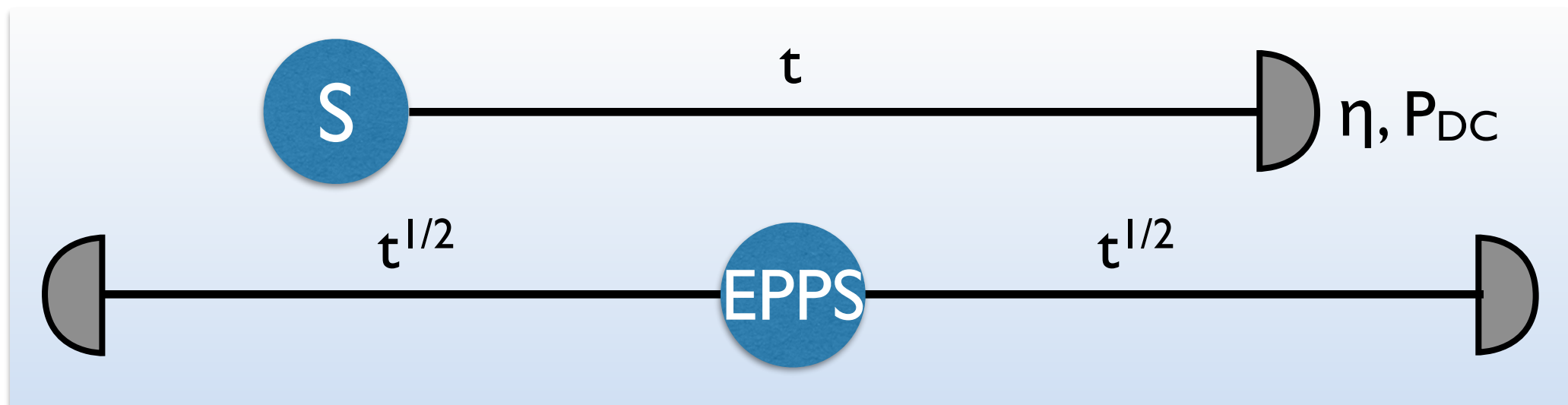
► Application to quantum key distribution (QKD)

Single photon (pair) emission character → security condition

I.2 HBT in quantum optics

► The context of today's quantum communication

Distribution of quantum bits of information using single photon and entangled photon pair sources



► Application to quantum key distribution (QKD)

Single photon (pair) emission character → security condition

HBT: an experimental resource for characterization

- Sources' emission statistics: Poissonian, Thermal (Bose-Einstein)
- Photons' coherence time

Outline

1. Introduction

2. HBT for characterizing single photon sources (SPS)

1. What is an SPS ?
2. “True” single photon sources (NV-centers, Qdots)
3. Heralded single photon sources (HSPS)
4. Connecting HSPSs in a quantum network

3. HBT for characterizing photon pair sources

2.1 Single photon sources

- ▶ What is a single photon source ?



2.1 Single photon sources

► What is a single photon source ?



P0: prob. of having no photon at all
 P1: prob. of having exactly 1 photon
 P2: prob. of having 2 photons
 ...

2.1 Single photon sources

▶ What is a single photon source ?



P_0 : prob. of having no photon at all
 P_1 : prob. of having exactly 1 photon
 P_2 : prob. of having 2 photons
 ...

▶ Different types of single photon sources (SPS)

2.1 Single photon sources

▶ What is a single photon source ?



P_0 : prob. of having no photon at all
 P_1 : prob. of having exactly 1 photon
 P_2 : prob. of having 2 photons
 ...

▶ Different types of single photon sources (SPS)

Molecule

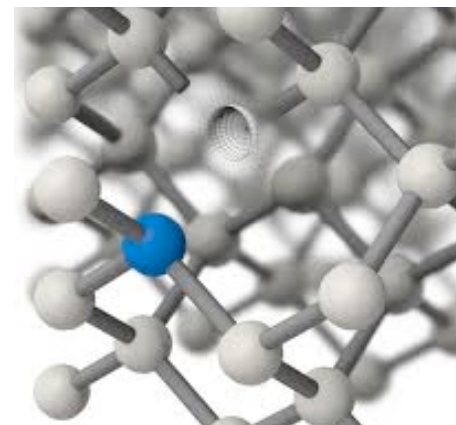
Semiconductor devices

NV center in diamond

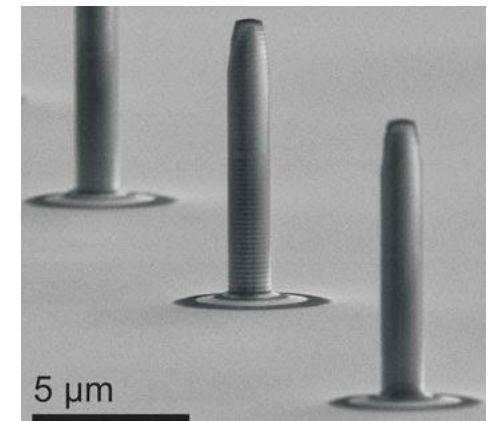
Isolated ion/atom

→ Any 2 energy-level isolated system

NV centers in diamond



Q-dots



2.1 Single photon sources

▶ What is a single photon source ?

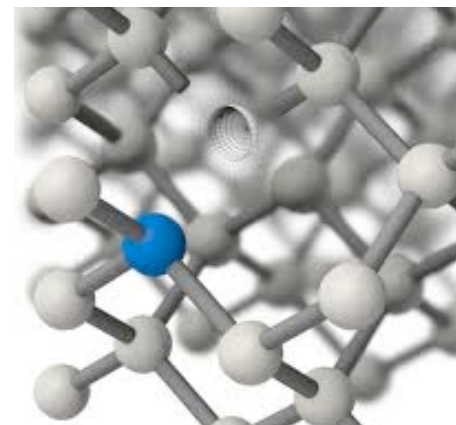


P_0 : prob. of having no photon at all
 P_1 : prob. of having exactly 1 photon
 P_2 : prob. of having 2 photons
 ...

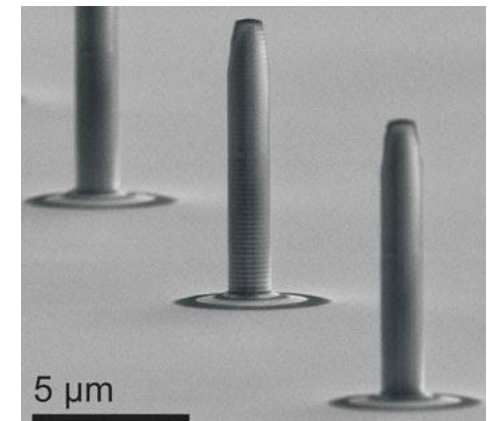
▶ Different types of single photon sources (SPS)

- Molecule
- Semiconductor devices
- NV center in diamond
- Isolated ion/atom
- Any 2 energy-level isolated system

NV centers in diamond



Q-dots



▶ What are they used for ?

2.1 Single photon sources

▶ What is a single photon source ?

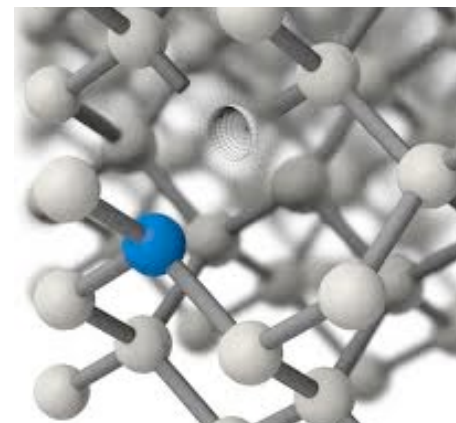


P_0 : prob. of having no photon at all
 P_1 : prob. of having exactly 1 photon
 P_2 : prob. of having 2 photons
 ...

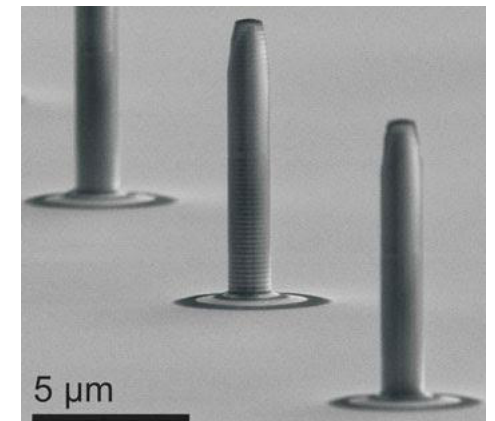
▶ Different types of single photon sources (SPS)

- Molecule
- Semiconductor devices
- NV center in diamond
- Isolated ion/atom
- Any 2 energy-level isolated system

NV centers in diamond



Q-dots



▶ What are they used for ?

- Quantum cryptography (QKD)
- Fundamental tests in quantum physics (Wheeler's delayed-choice exp.)
- (Quantum) metrology (phase measurements, magnetometry, etc.)

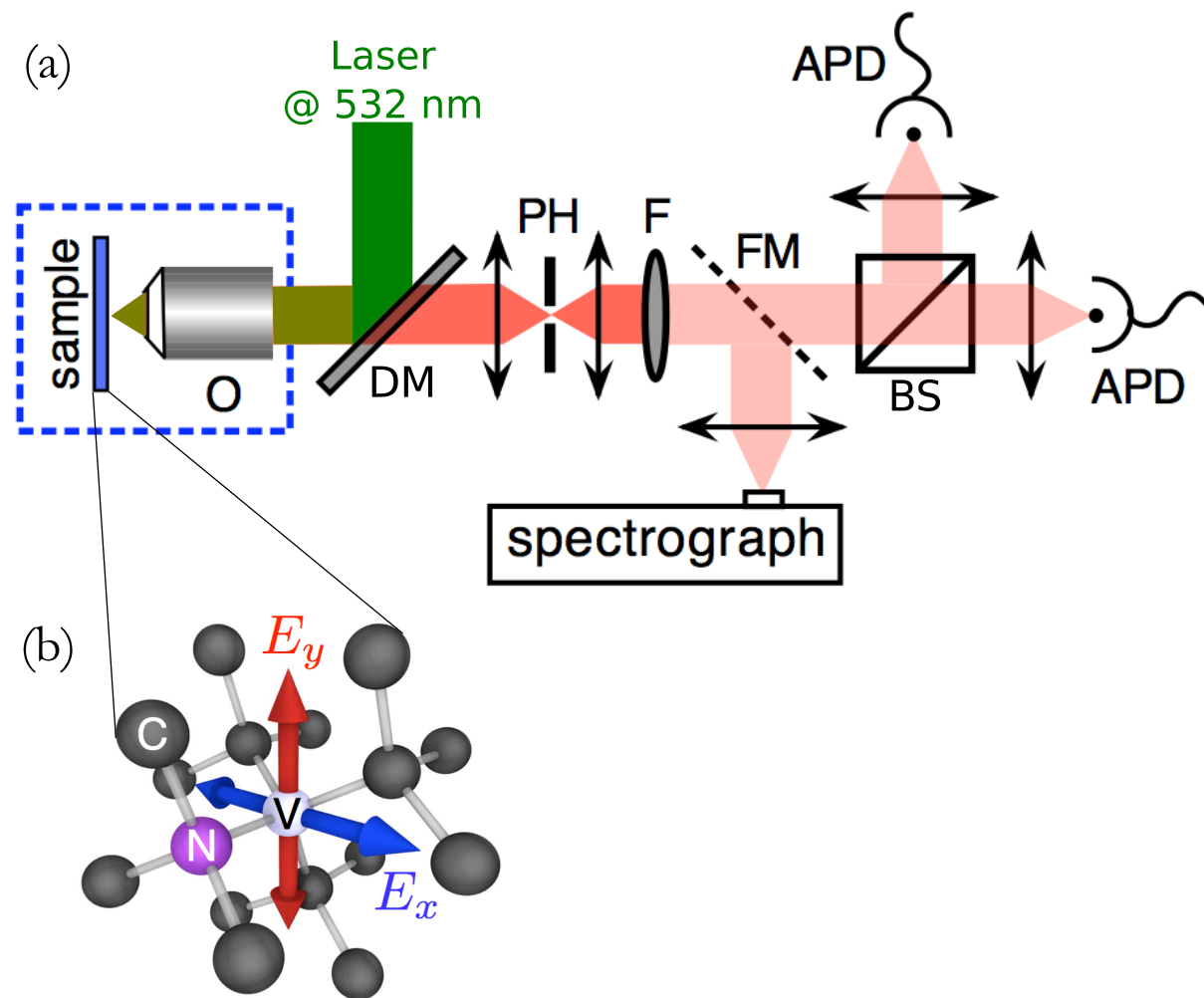
2.2 True SPS based on a single NV center in diamond

► Confocal microscopy based setup

[Wrachtrup's group, Stuttgart] P. Siyushev *et al.*, New J. Phys. 11, 113029 (2009)

2.2 True SPS based on a single NV center in diamond

► Confocal microscopy based setup

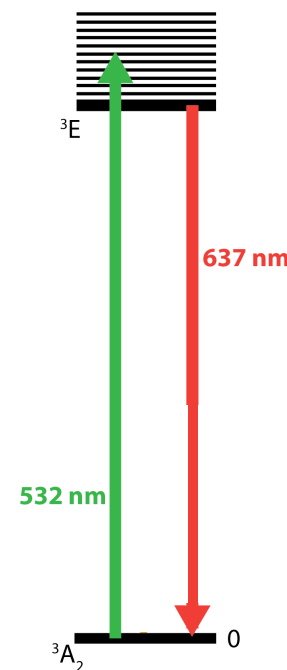
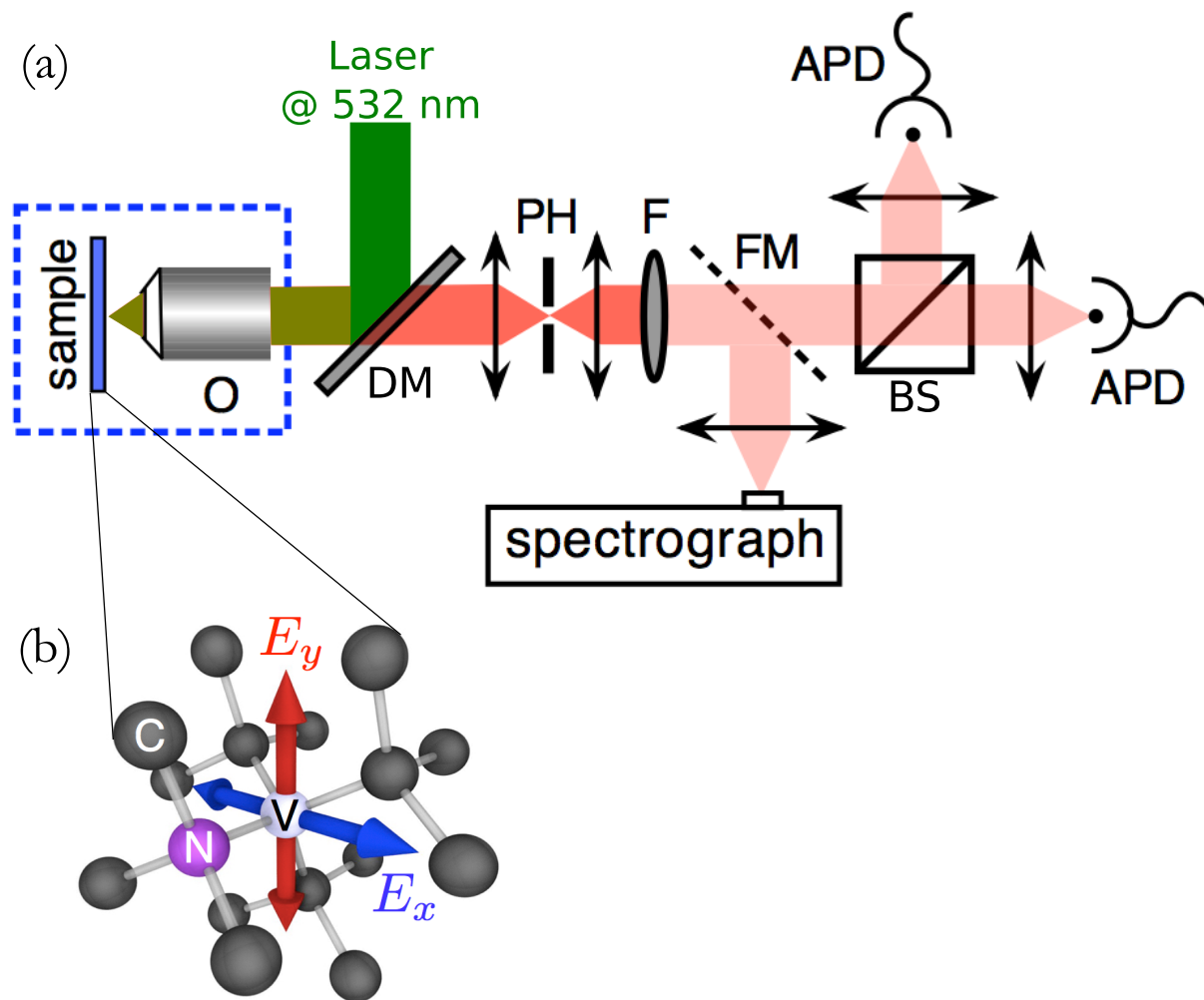


[Wrachtrup's group, Stuttgart] P. Siyushev *et al.*, New J. Phys. 11, 113029 (2009)

2.2 True SPS based on a single NV center in diamond

► Confocal microscopy based setup

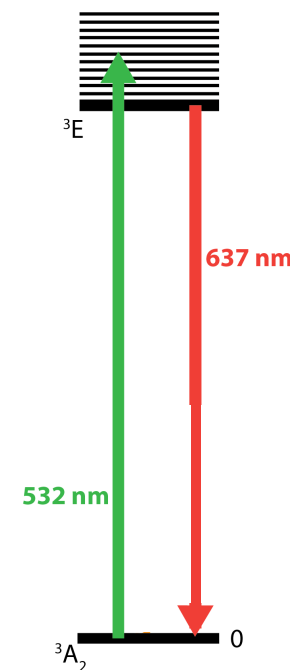
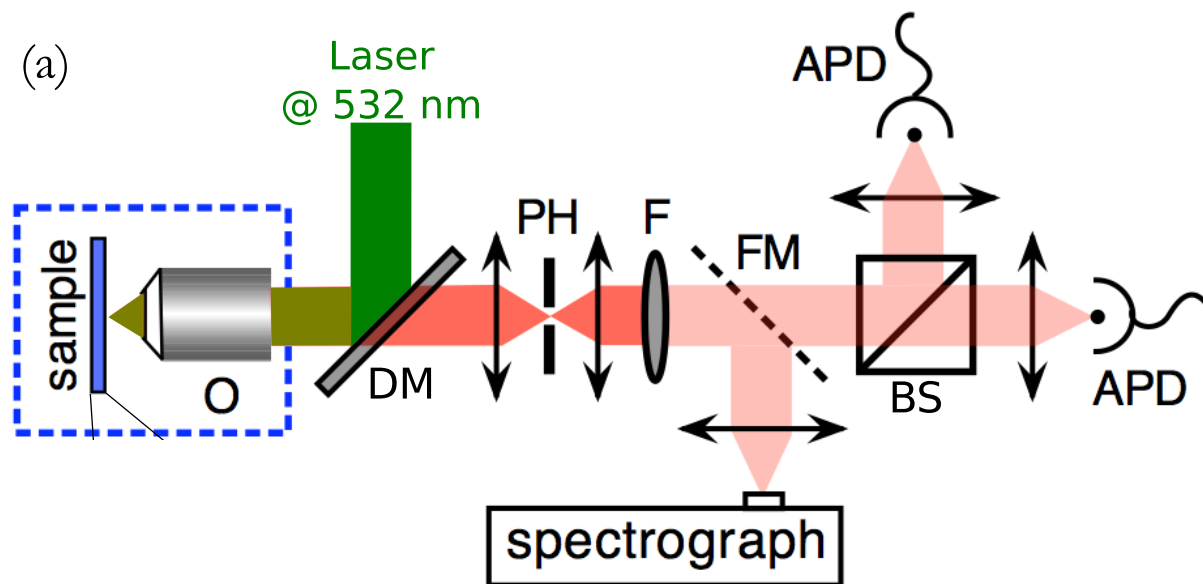
► “Artificial atom” energy levels



2.2 True SPS based on a single NV center in diamond

► Confocal microscopy based setup

► “Artificial atom” energy levels



► Diamond based nanostructures

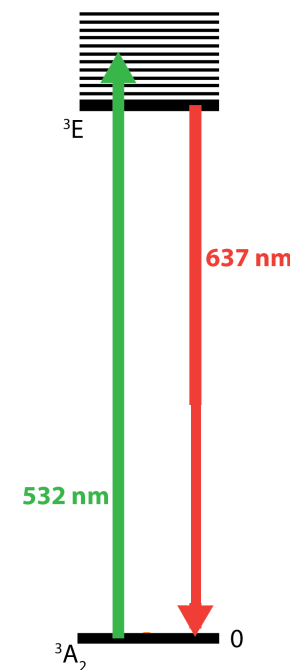
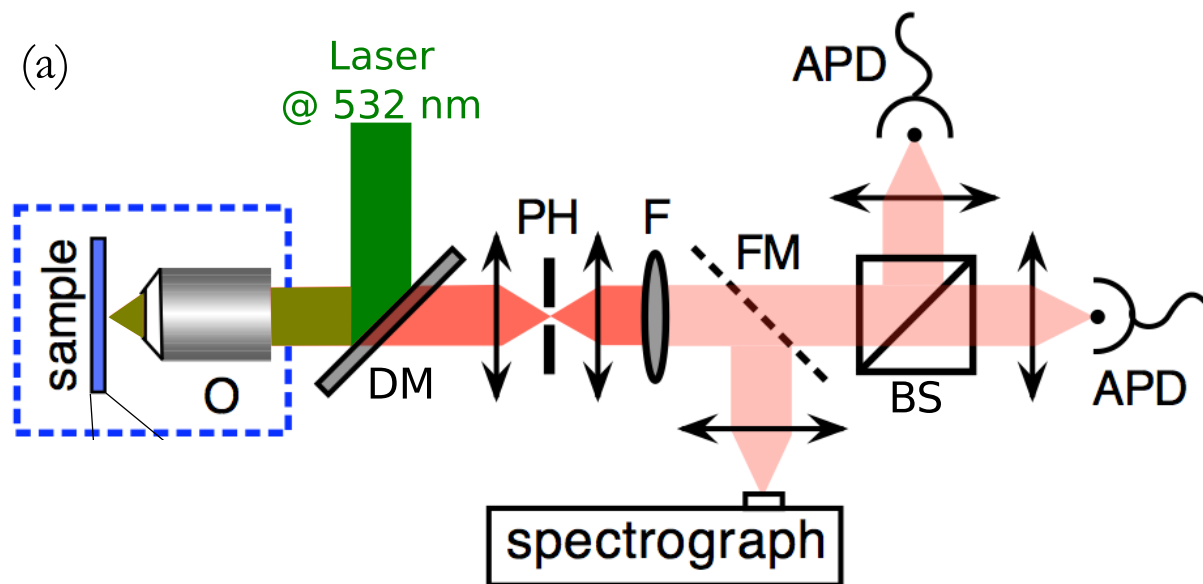
[Wrachtrup's group, Stuttgart] P. Siyushev *et al.*, New J. Phys. 11, 113029 (2009)

[Lukin/Loncar's group, Cambridge] B. Hausmann *et al.*, New J. Phys. 13, 045004 (2011)

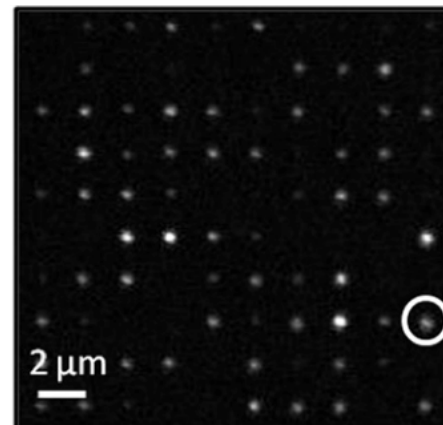
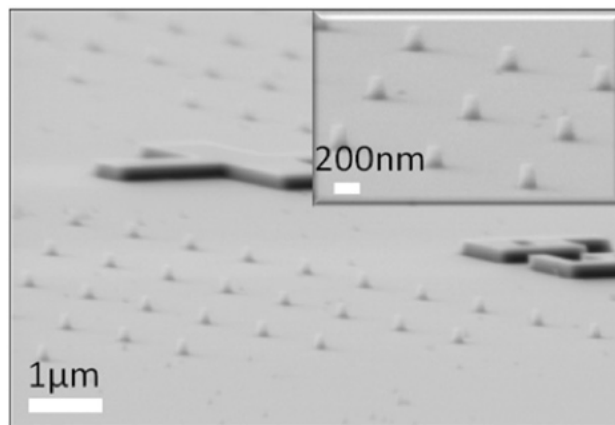
2.2 True SPS based on a single NV center in diamond

► Confocal microscopy based setup

► “Artificial atom” energy levels



► Diamond based nanostructures



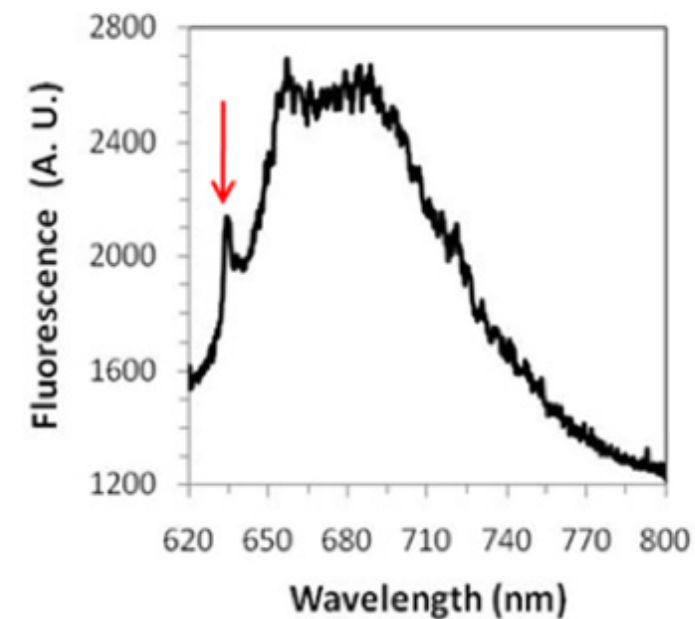
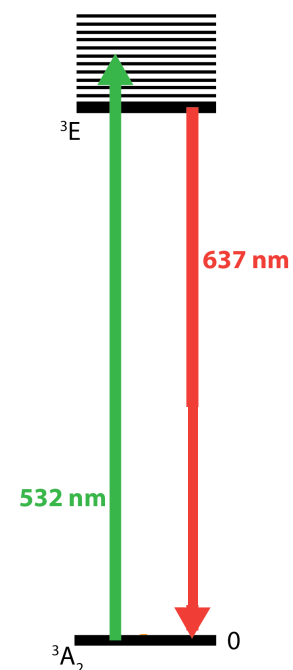
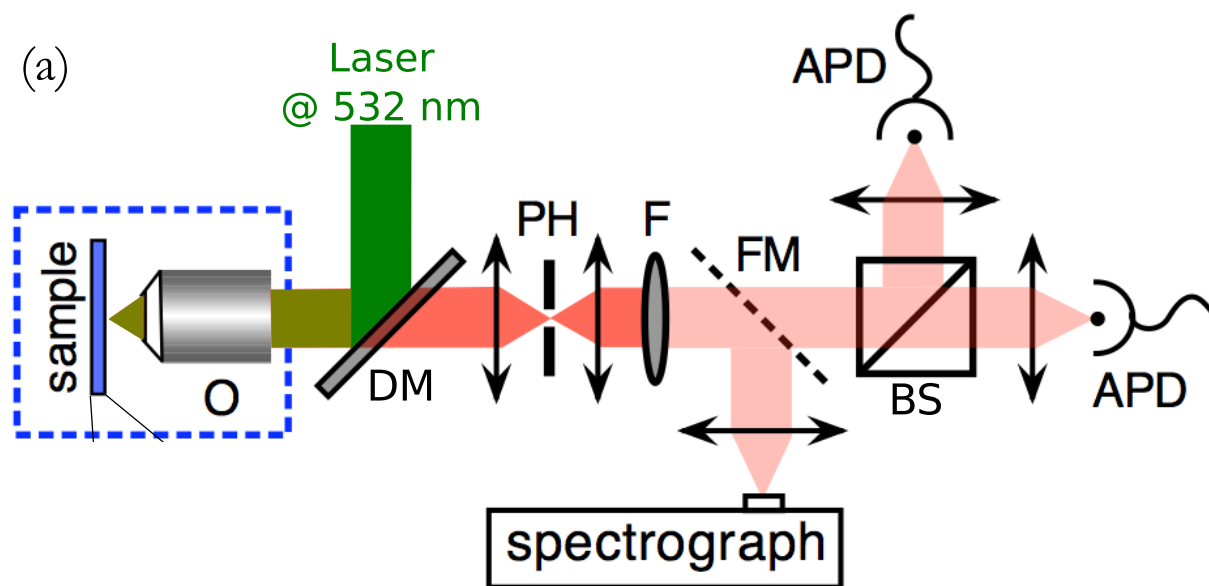
[Wrachtrup's group, Stuttgart] P. Siyushev *et al.*, New J. Phys. 11, 113029 (2009)

[Lukin/Loncar's group, Cambridge] B. Hausmann *et al.*, New J. Phys. 13, 045004 (2011)

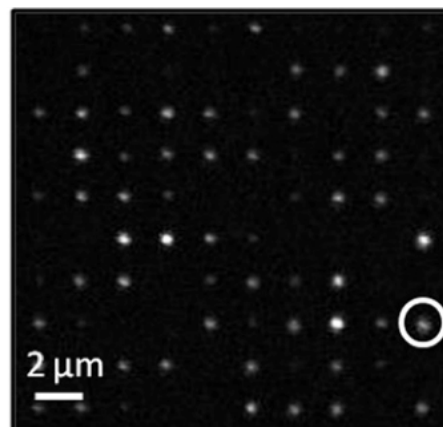
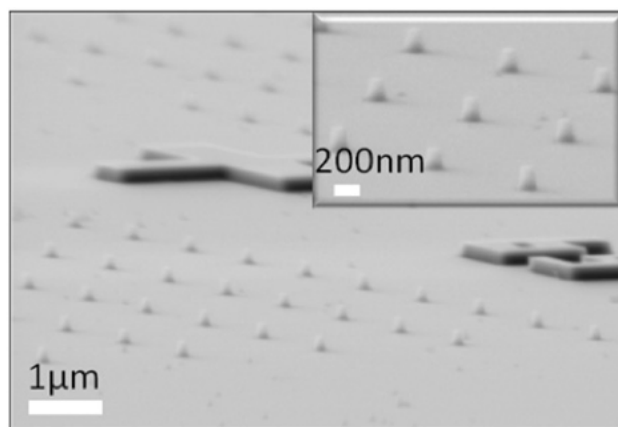
2.2 True SPS based on a single NV center in diamond

► Confocal microscopy based setup

► “Artificial atom” energy levels



► Diamond based nanostructures



Zero photon line

[Wrachtrup's group, Stuttgart] P. Siyushev *et al.*, *New J. Phys.* 11, 113029 (2009)

[Lukin/Loncar's group, Cambridge] B. Hausmann *et al.*, *New J. Phys.* 13, 045004 (2011)

2.2 True SPS based on a single NV center in diamond

Antibunching measurements

[Lukin/Loncar's group, Cambridge]

B. Hausmann *et al.*, New J. Phys. **13**, 045004 (2011) 14

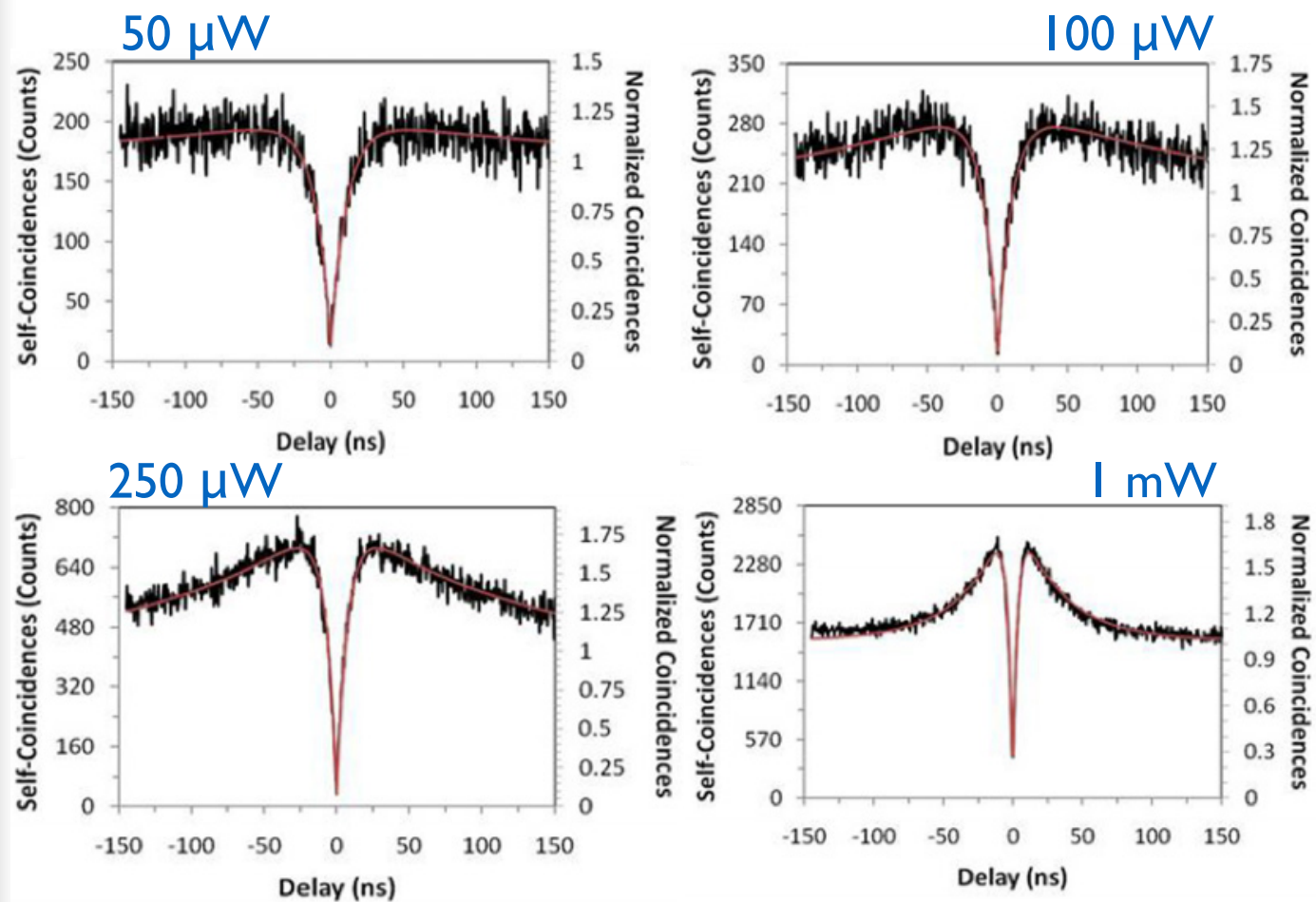
2.2 True SPS based on a single NV center in diamond

Antibunching measurements

Continuous wave excitation

2.2 True SPS based on a single NV center in diamond Antibunching measurements

Continuous wave excitation

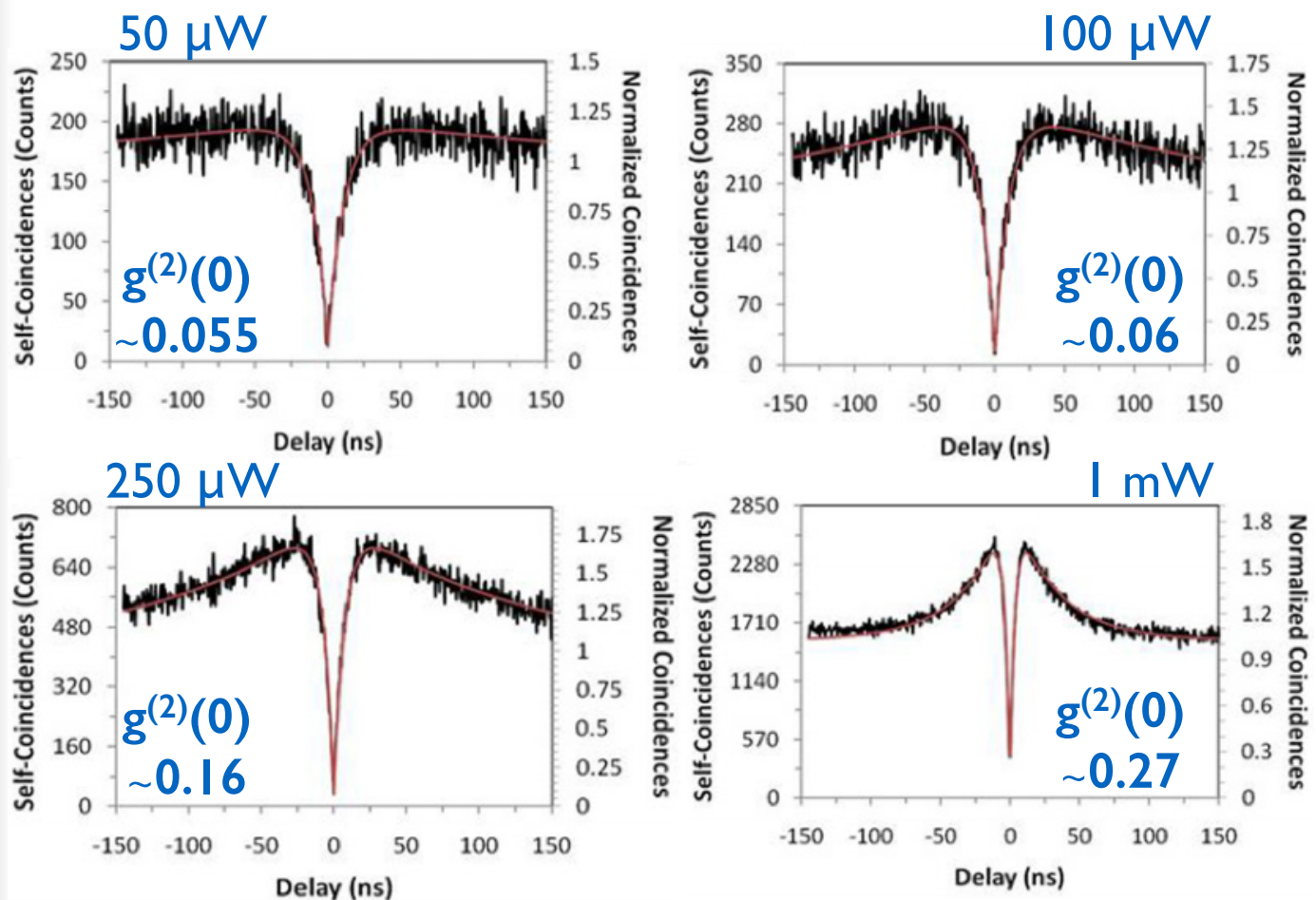


[Lukin/Loncar's group, Cambridge]

B. Hausmann *et al.*, New J. Phys. 13, 045004 (2011) 14

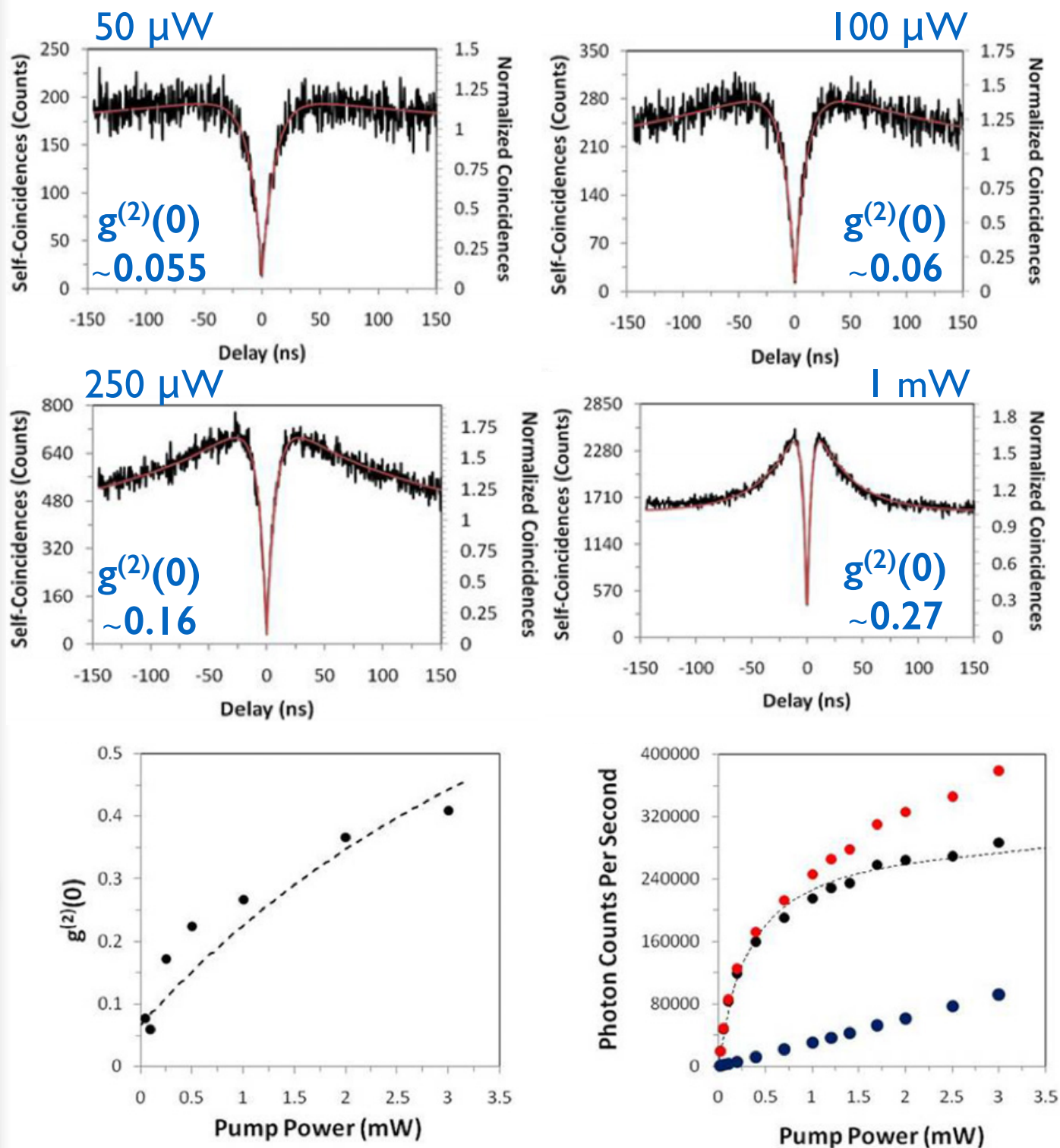
2.2 True SPS based on a single NV center in diamond Antibunching measurements

Continuous wave excitation



2.2 True SPS based on a single NV center in diamond Antibunching measurements

Continuous wave excitation

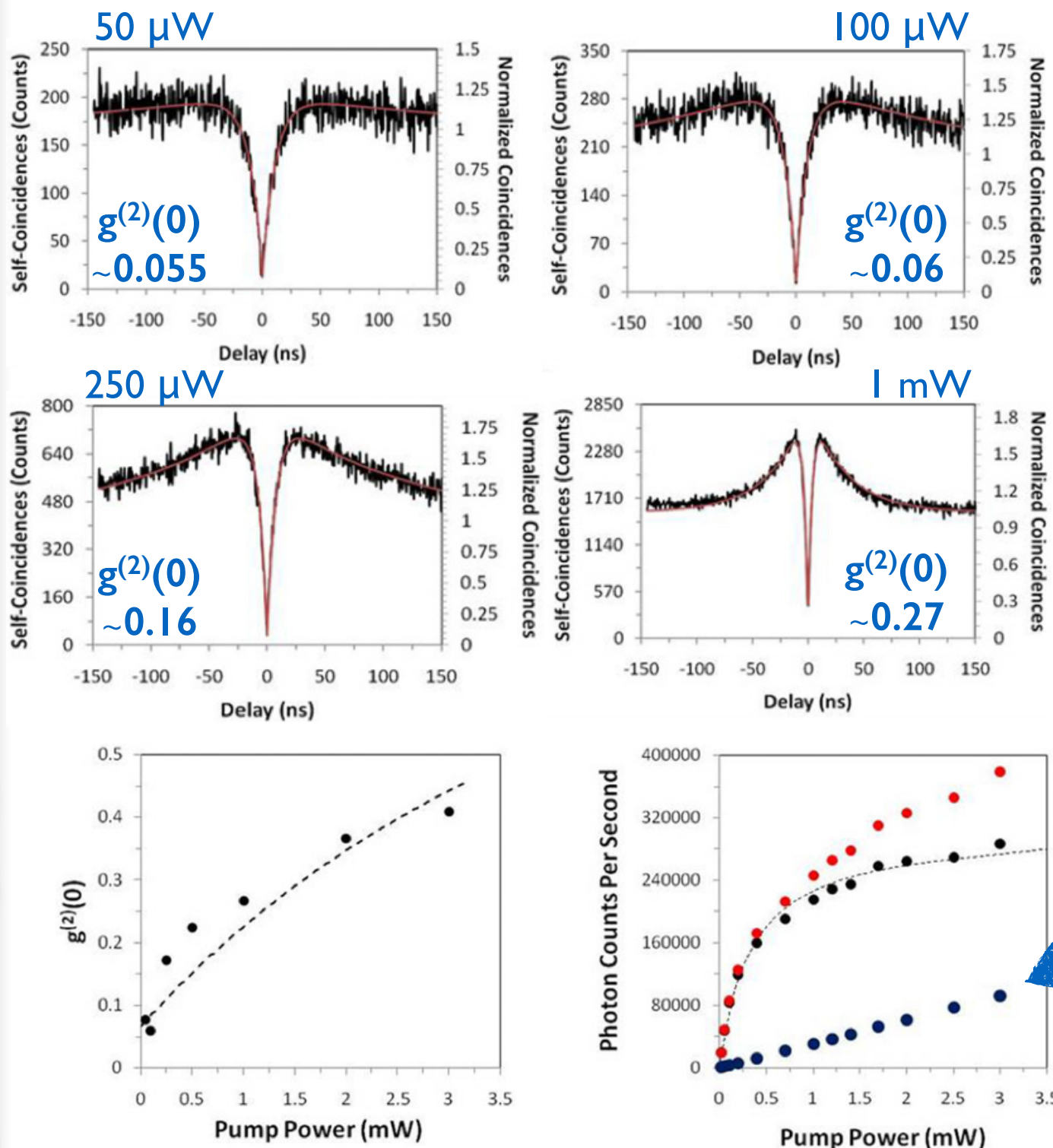


[Lukin/Loncar's group, Cambridge]

B. Hausmann *et al.*, New J. Phys. 13, 045004 (2011) 14

2.2 True SPS based on a single NV center in diamond Antibunching measurements

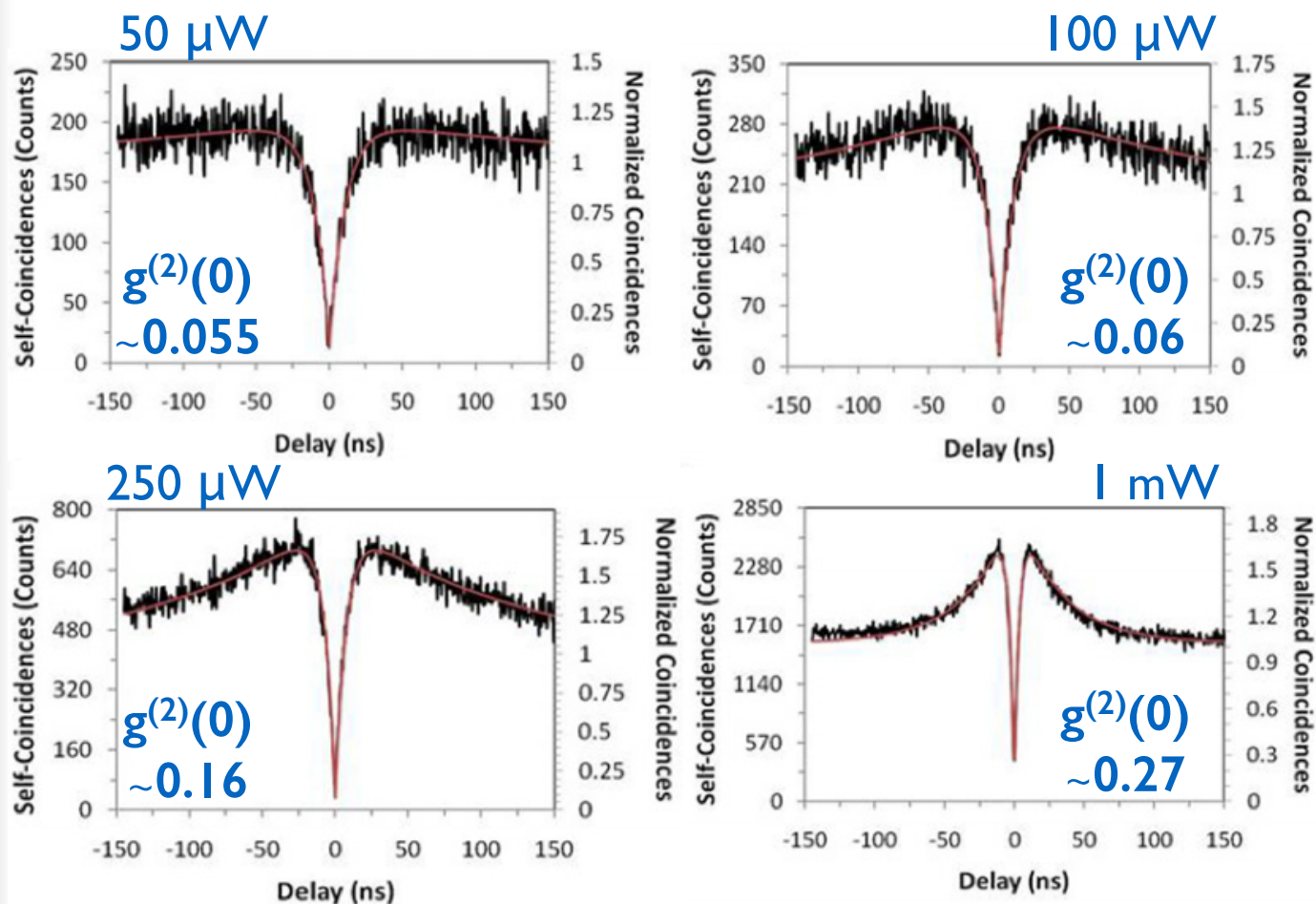
Continuous wave excitation



Poissonian
background photons

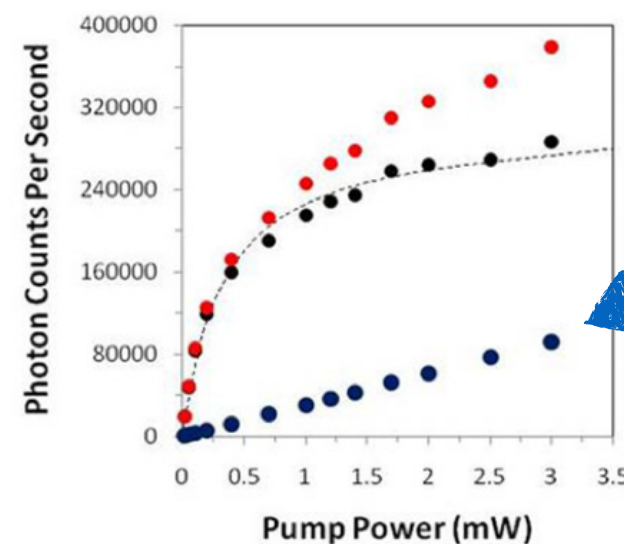
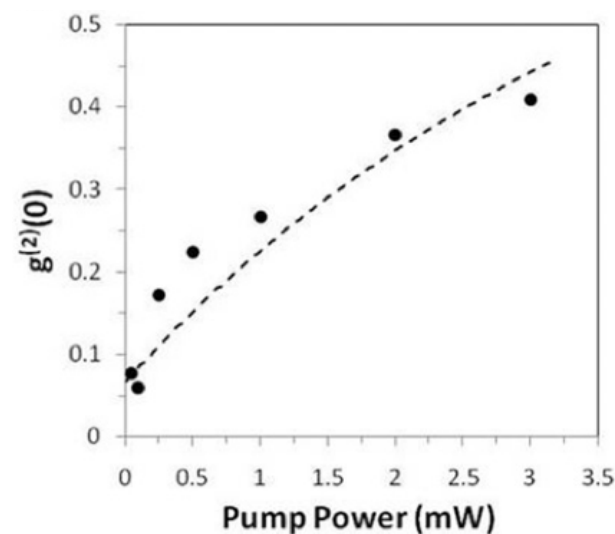
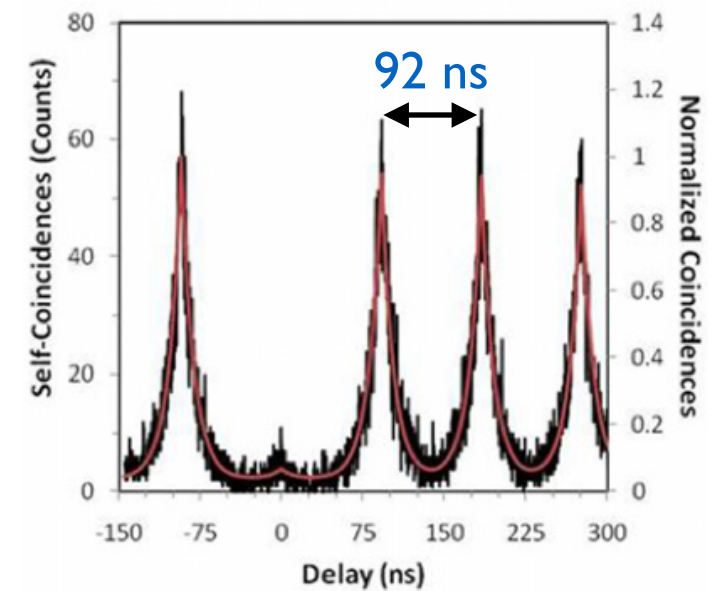
2.2 True SPS based on a single NV center in diamond Antibunching measurements

Continuous wave excitation



Pulsed excitation

$\langle 65 \rangle \mu\text{W} \rightarrow g^{(2)}(0) \sim 0.16$



[Lukin/Loncar's group, Cambridge]

B. Hausmann *et al.*, New J. Phys. 13, 045004 (2011) 14

2.2 True SPS based on a single quantum dot

- ▶ Many (many) types and fab. processes

2.2 True SPS based on a single quantum dot

► Many (many) types and fab. processes

→ an entire zoology of properties

Material system	λ (nm)	Tmax (K)
InAs/GaAs	~850 - 1000	50
InGaAs/InAs/GaAs	~1300	90
InP/InGaP	~650 - 700	50
InAs/InP	~1550	50
GaN/AlN	~250 - 500	200
CdSe/ZnSSe	~500 - 550	200
...

2.2 True SPS based on a single quantum dot

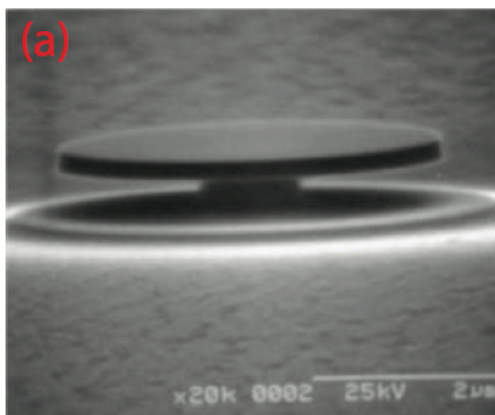
▶ Many (many) types and fab. processes

→ an entire zoology of properties

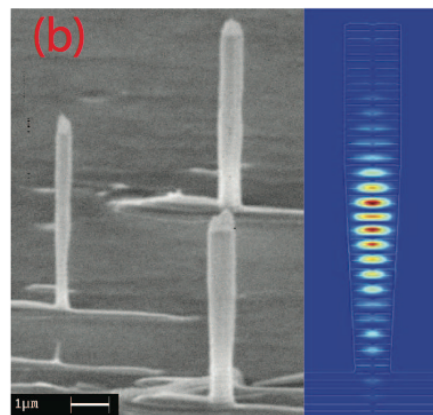
Material system	λ (nm)	Tmax (K)
InAs/GaAs	~850 - 1000	50
InGaAs/InAs/GaAs	~1300	90
InP/InGaP	~650 - 700	50
InAs/InP	~1550	50
GaN/AlN	~250 - 500	200
CdSe/ZnSSe	~500 - 550	200
...

▶ Surrounded by a cavity and search for the single exciton line (X-polarized)

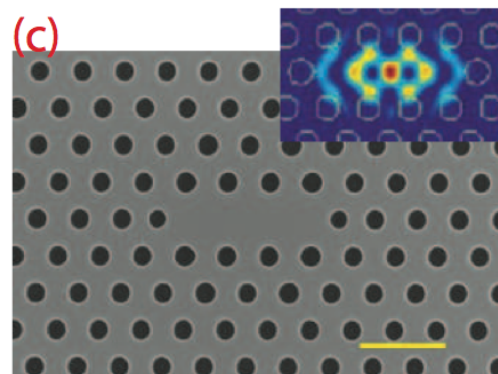
μ -disk



μ -pillar



photonic Xtal



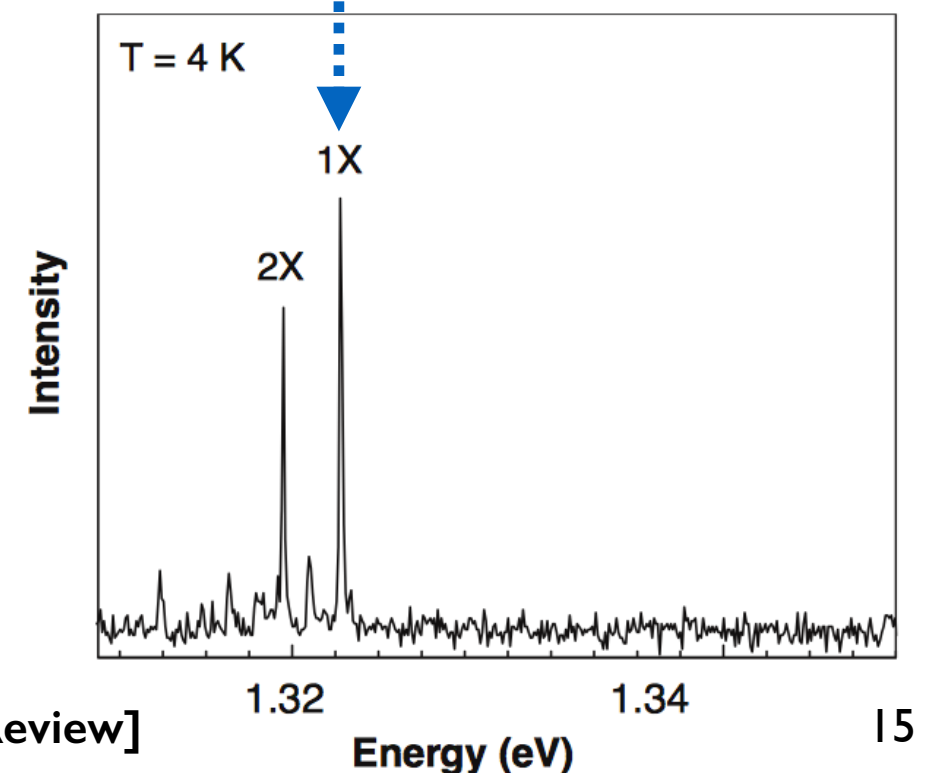
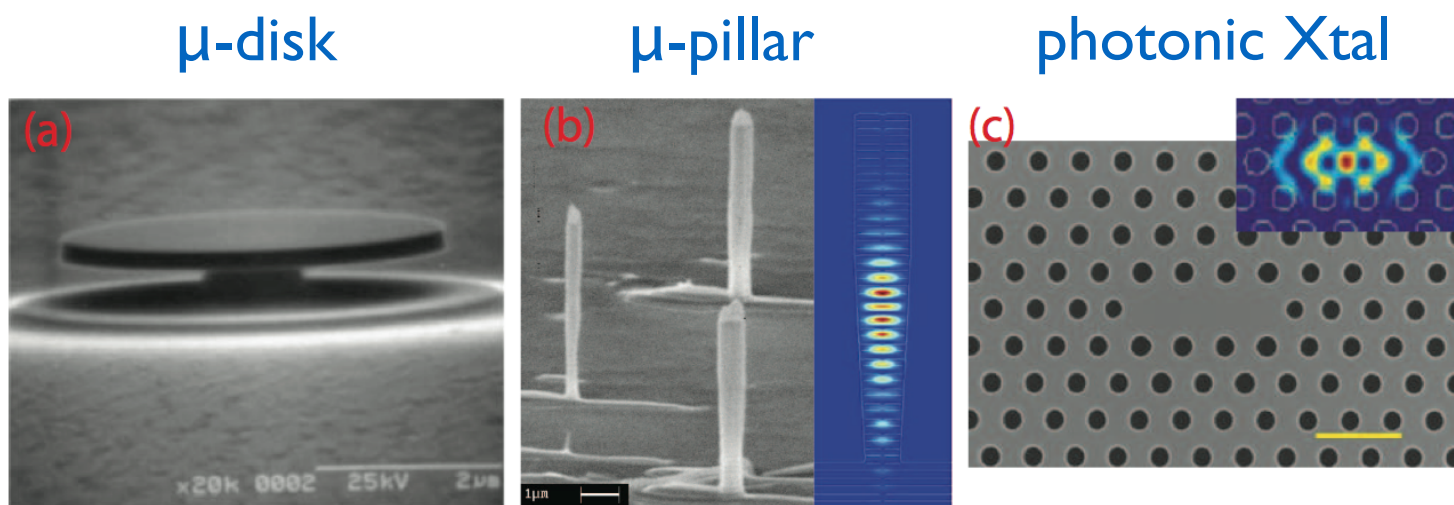
2.2 True SPS based on a single quantum dot

▶ Many (many) types and fab. processes

→ an entire zoology of properties

Material system	λ (nm)	Tmax (K)
InAs/GaAs	~850 - 1000	50
InGaAs/InAs/GaAs	~1300	90
InP/InGaP	~650 - 700	50
InAs/InP	~1550	50
GaN/AlN	~250 - 500	200
CdSe/ZnSSe	~500 - 550	200
...

▶ Surrounded by a cavity and search for the single exciton line (X-polarized)



[Imamoglu's group, ETH] C. Becher *et al.*, Physica E 13, 412 (2002)

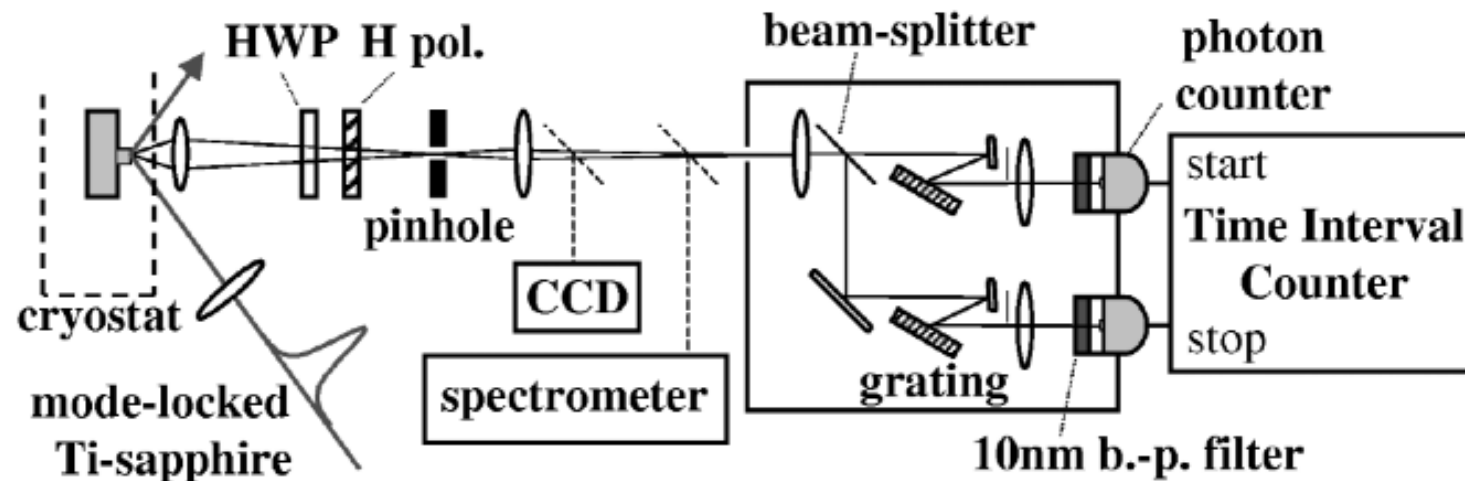
[Vuckovic's group, Stanford] J.Vuckovic *et al.*, Rep. Prog. Phys. 75 126503 (2012) [Review]

2.2 True SPS based on a single quantum dot

- ▶ **Standard setup using a cryostat**

2.2 True SPS based on a single quantum dot

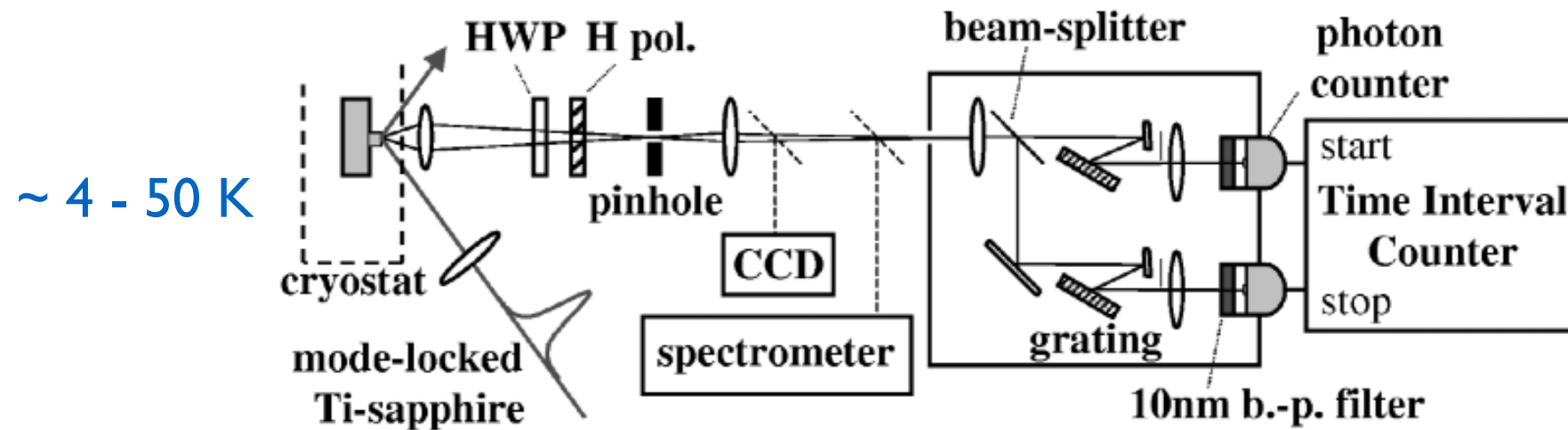
► Standard setup using a cryostat



[Yamamoto's group, Stanford] C. Santori et al., PRL 86, 1502 (2001)

2.2 True SPS based on a single quantum dot

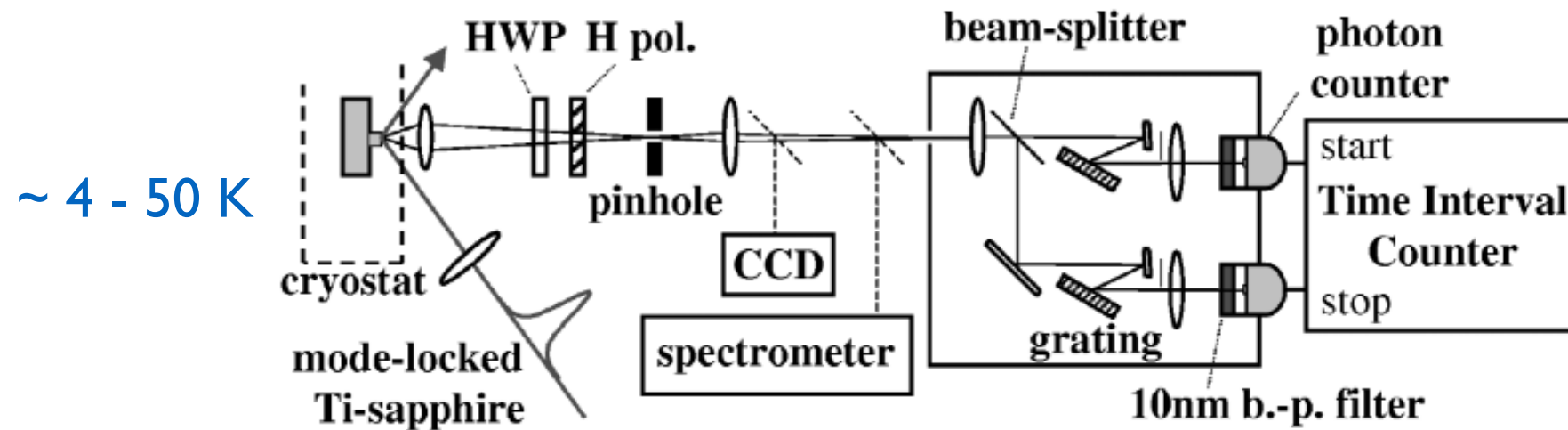
► Standard setup using a cryostat



[Yamamoto's group, Stanford] C. Santori *et al.*, PRL **86**, 1502 (2001)

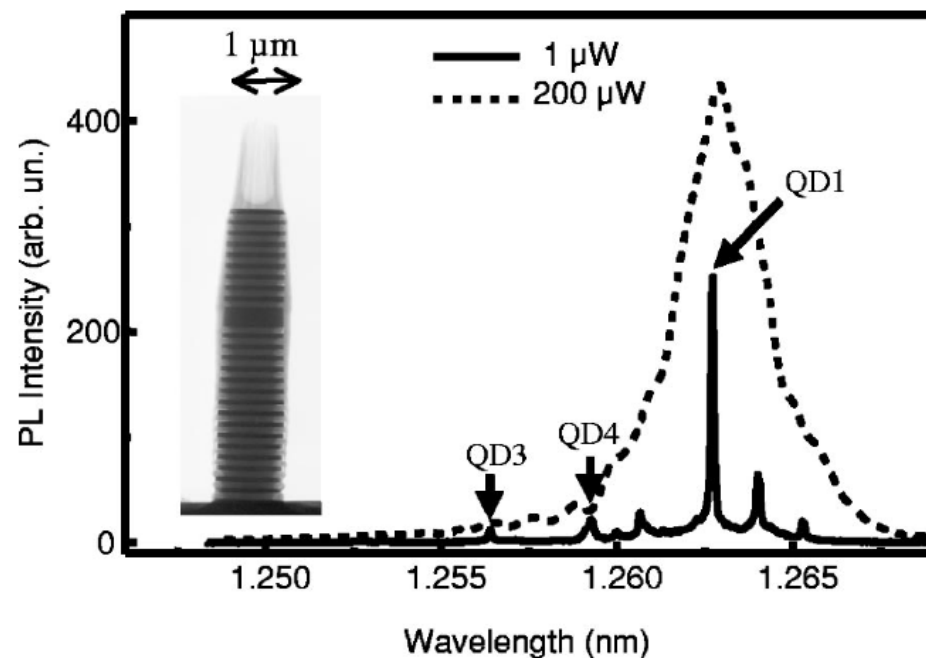
2.2 True SPS based on a single quantum dot

► Standard setup using a cryostat



[Yamamoto's group, Stanford] C. Santori *et al.*, PRL 86, 1502 (2001)

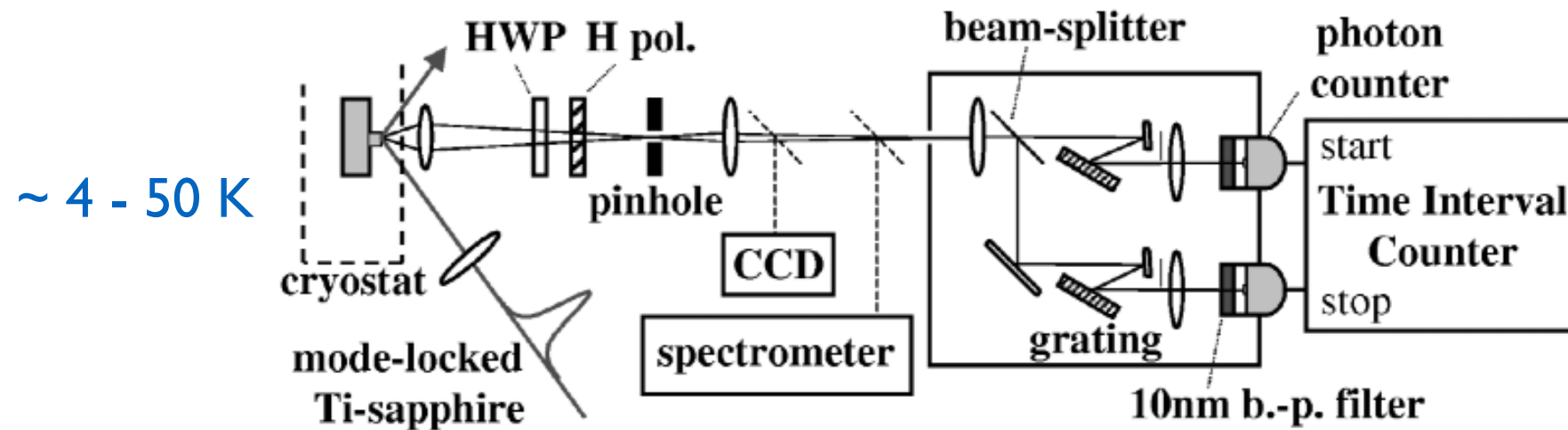
► Weak vs strong excitation



[Gérard's group, CEA] E. Moreau *et al.*, APL 79, 2865 (2001)

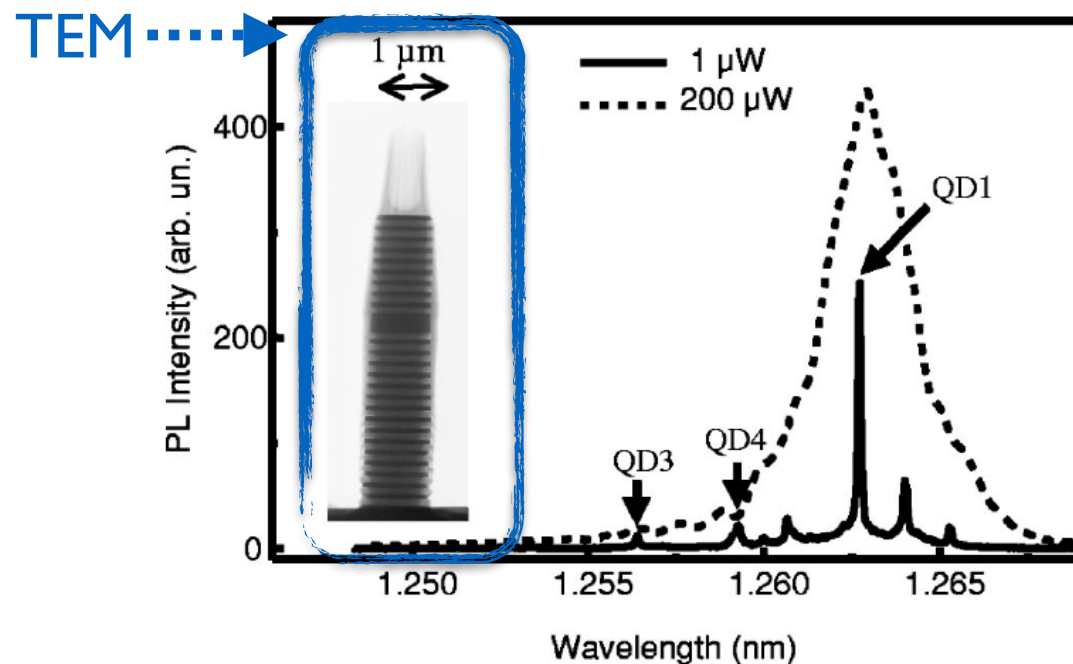
2.2 True SPS based on a single quantum dot

► Standard setup using a cryostat



[Yamamoto's group, Stanford] C. Santori *et al.*, PRL 86, 1502 (2001)

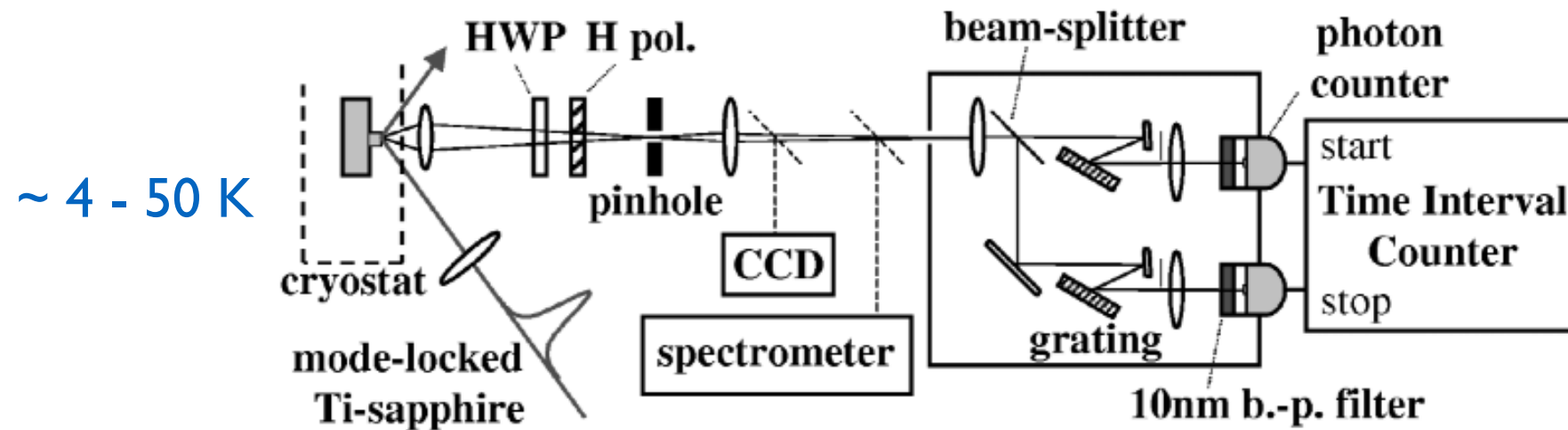
► Weak vs strong excitation



[Gérard's group, CEA] E. Moreau *et al.*, APL 79, 2865 (2001)

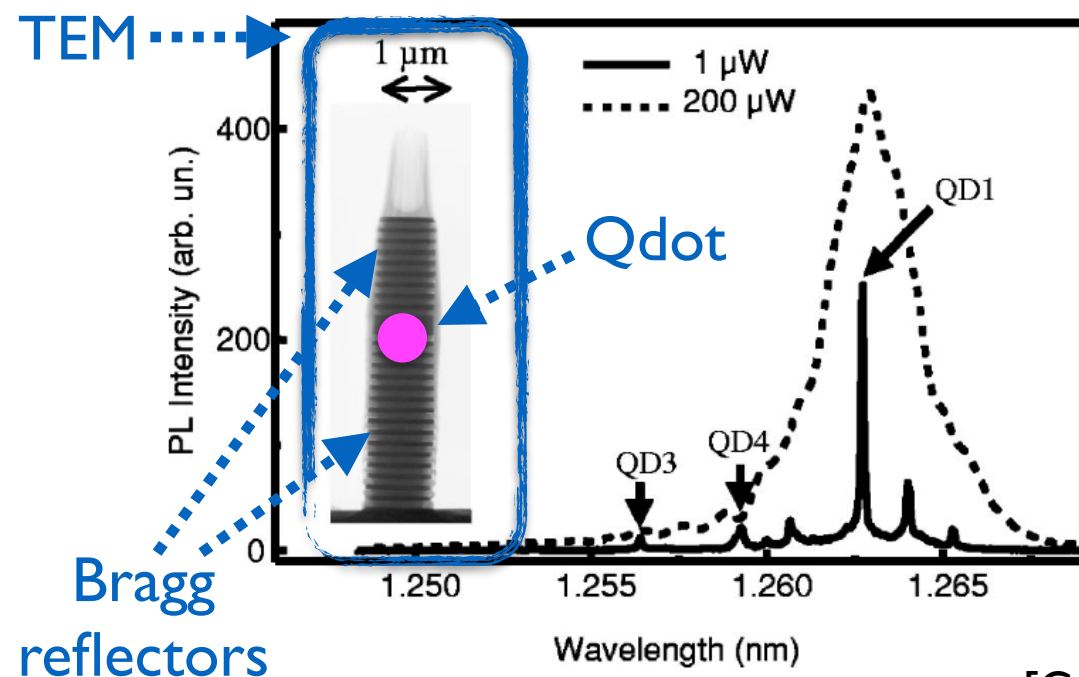
2.2 True SPS based on a single quantum dot

► Standard setup using a cryostat



[Yamamoto's group, Stanford] C. Santori *et al.*, PRL 86, 1502 (2001)

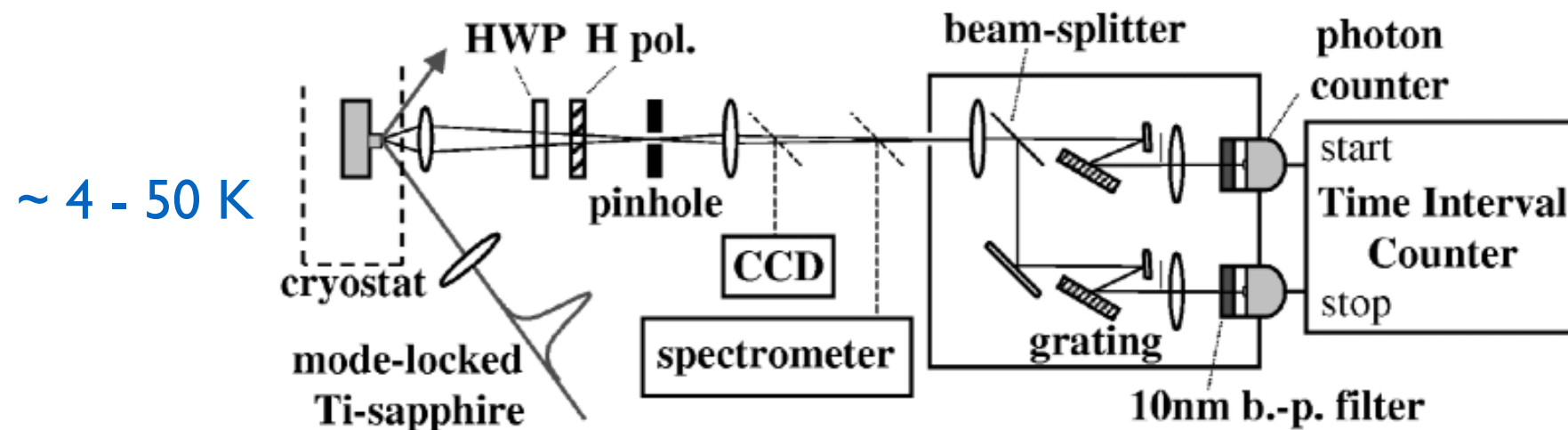
► Weak vs strong excitation



[Gérard's group, CEA] E. Moreau *et al.*, APL 79, 2865 (2001)

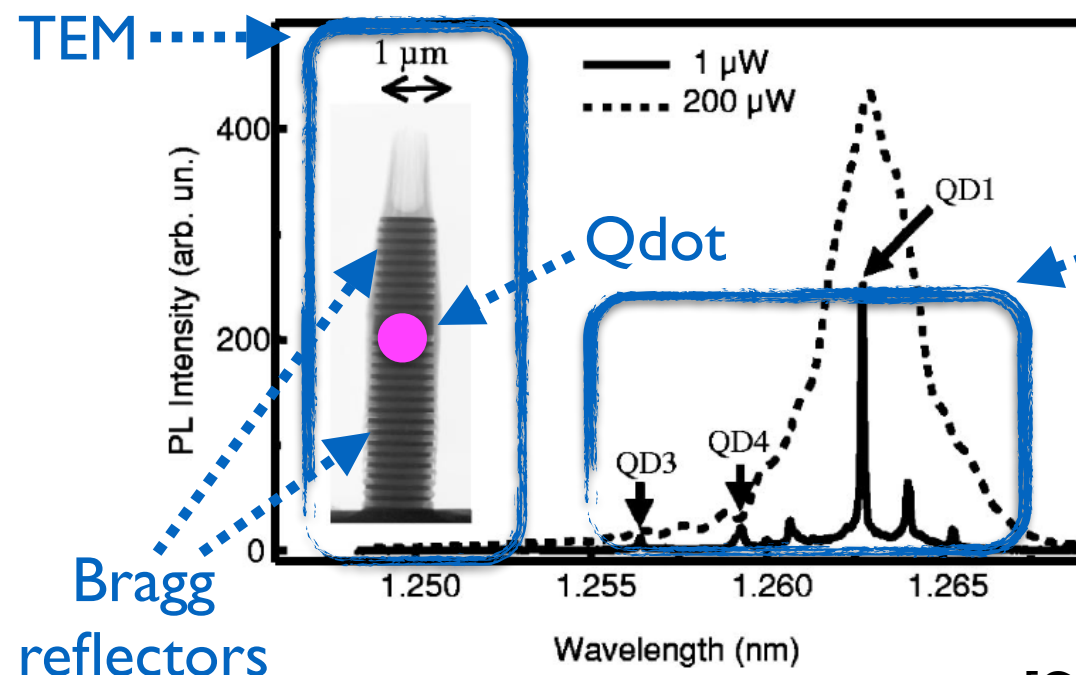
2.2 True SPS based on a single quantum dot

► Standard setup using a cryostat



[Yamamoto's group, Stanford] C. Santori *et al.*, PRL 86, 1502 (2001)

► Weak vs strong excitation



A few sharp emission lines

→ different Qdots having a single exciton

→ so-called "X-line"

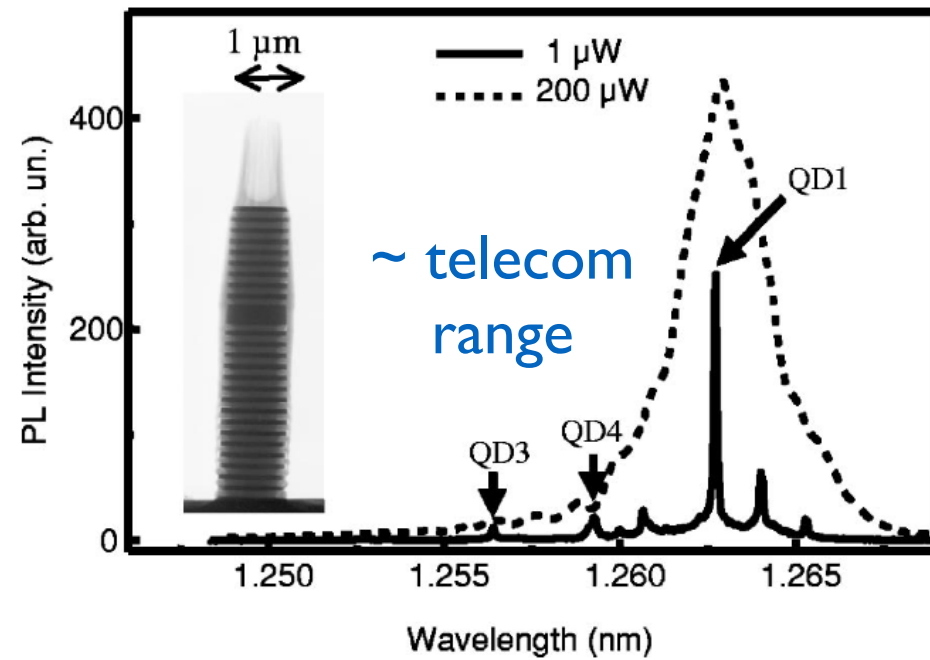
[Gérard's group, CEA] E. Moreau *et al.*, APL 79, 2865 (2001)

2.2 True SPS based on a single quantum dot

Antibunching measurements

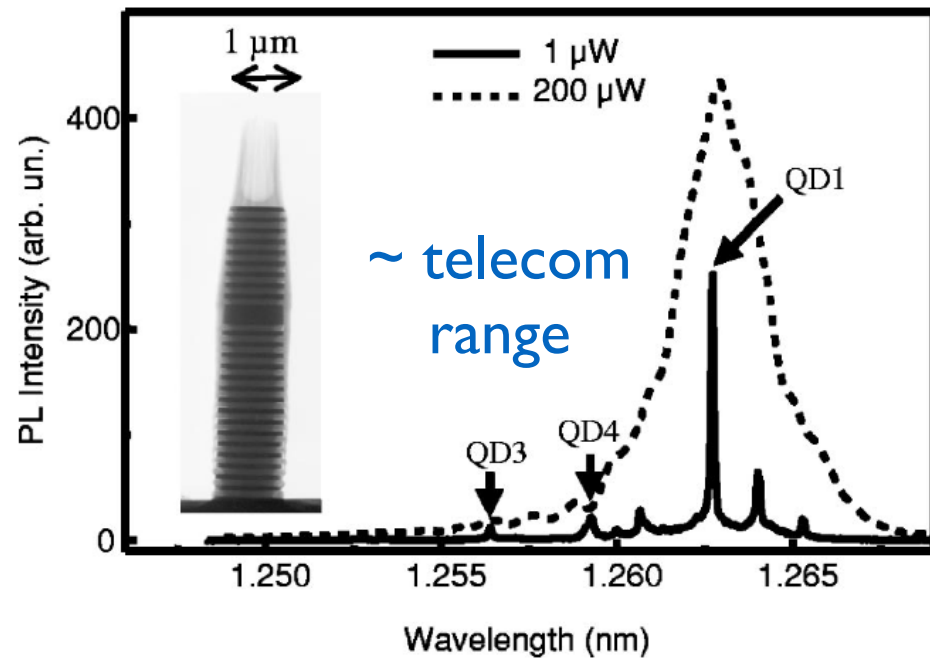
2.2 True SPS based on a single quantum dot Antibunching measurements

In 2001, Gérard's group, CEA

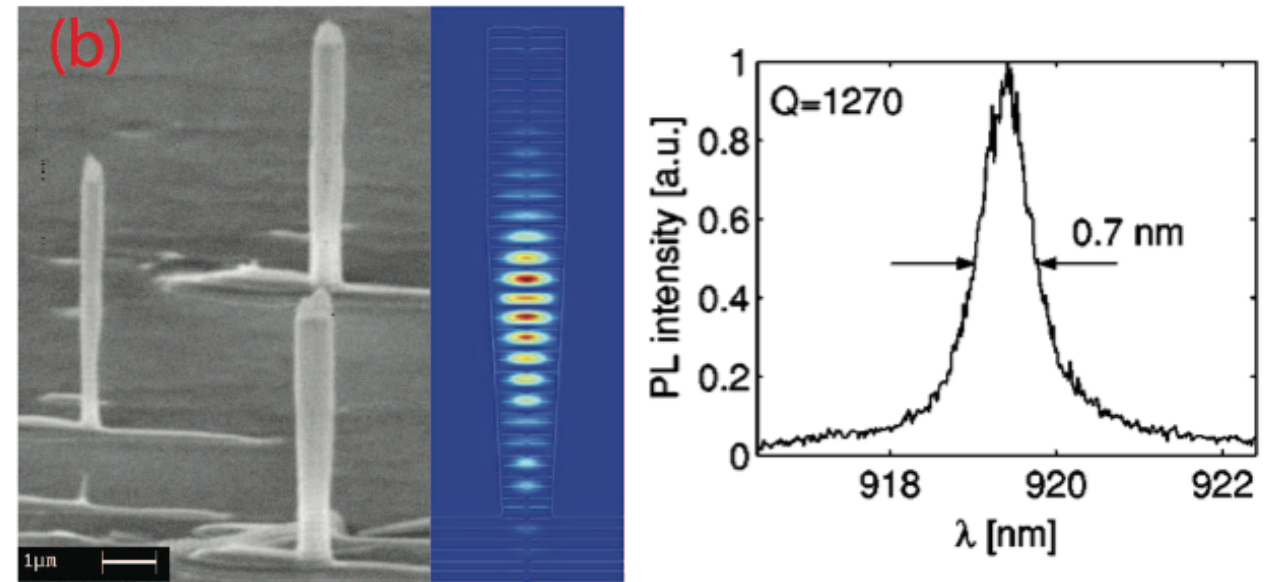


2.2 True SPS based on a single quantum dot Antibunching measurements

In 2001, Gérard's group, CEA

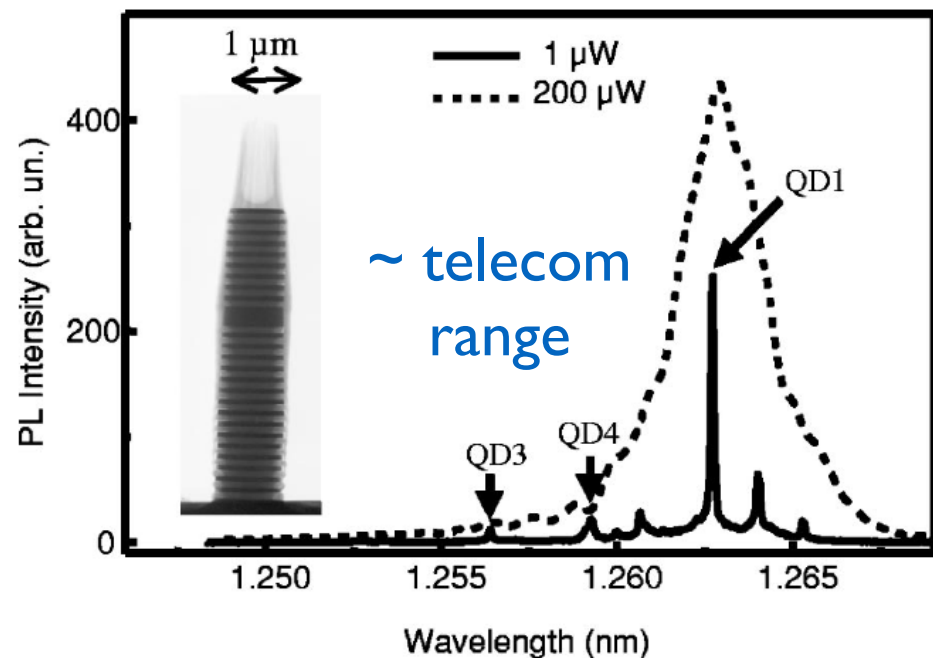


In 2003, Vuckovic's group, Stanford

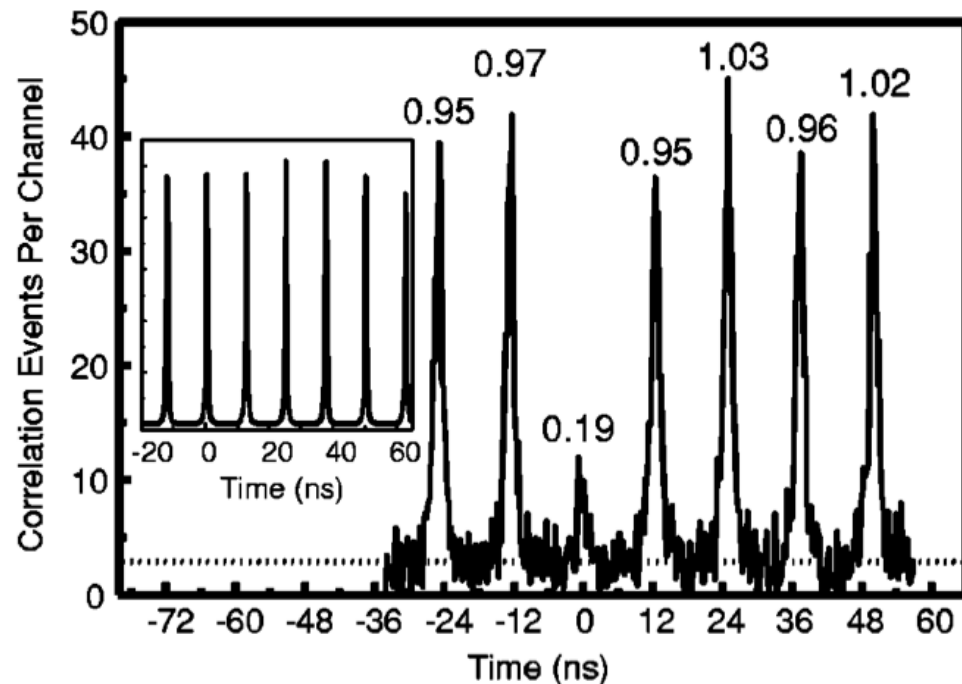


2.2 True SPS based on a single quantum dot Antibunching measurements

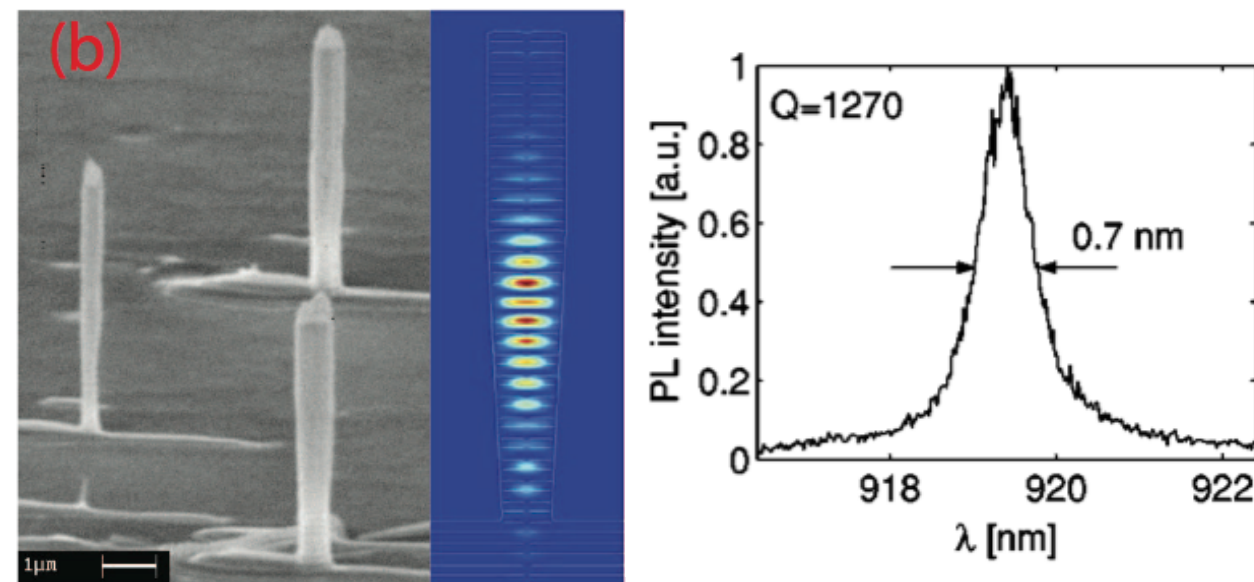
In 2001, Gérard's group, CEA



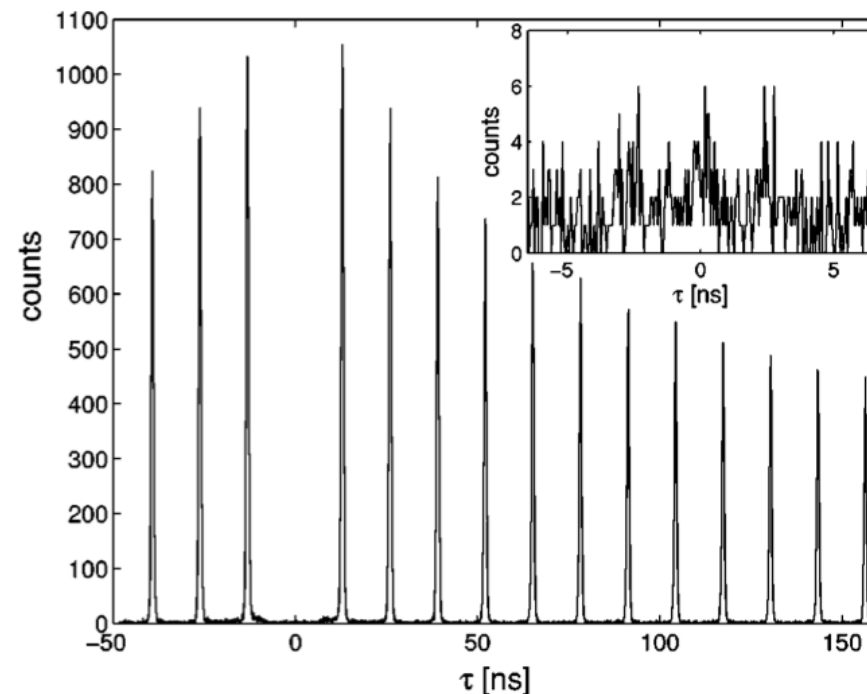
$$g^{(2)}(0) \sim 0.19, Q \sim 400, \eta_{\text{coll}} \sim 10^{-3}$$



In 2003, Vuckovic's group, Stanford



$$g^{(2)}(0) \sim 0.02, Q \sim 4000, \eta_{\text{coll}} \sim 10^{-2}$$

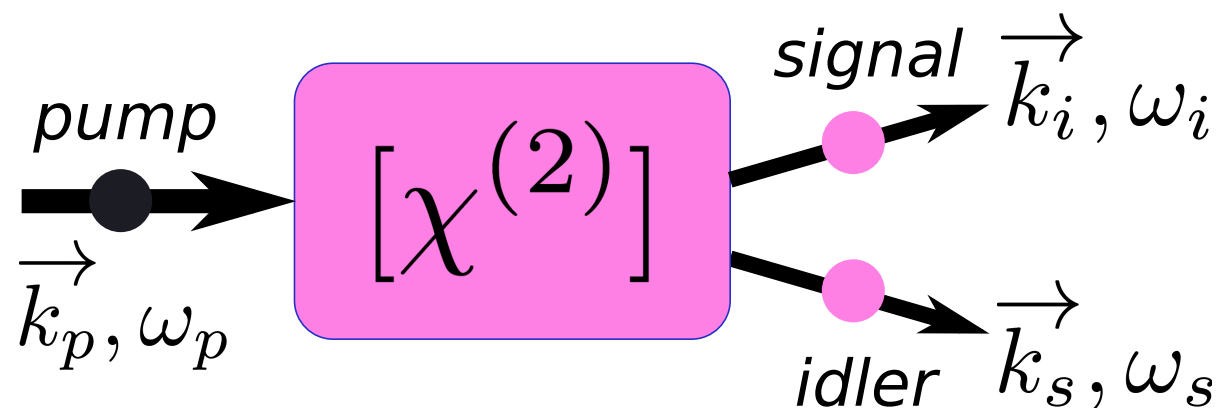


2.3 Heralded SPS based on nonlinear optics

- ▶ Spontaneous parametric downconversion (SPDC) in nonlinear optics

2.3 Heralded SPS based on nonlinear optics

► Spontaneous parametric downconversion (SPDC) in nonlinear optics

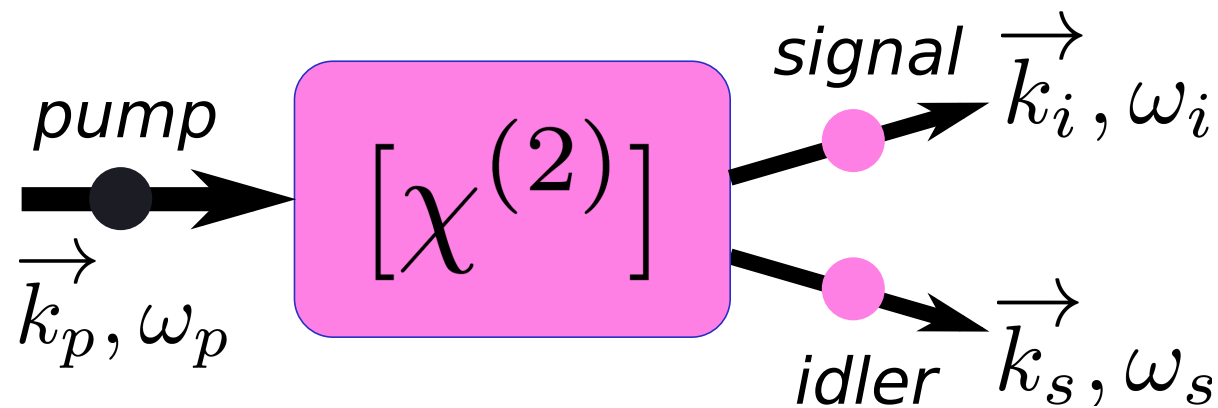


$$\omega_p = \omega_s + \omega_i \quad \text{energy}$$

$$\vec{k}_p = \vec{k}_s + \vec{k}_i \quad \text{phase matching}$$

2.3 Heralded SPS based on nonlinear optics

► Spontaneous parametric downconversion (SPDC) in nonlinear optics



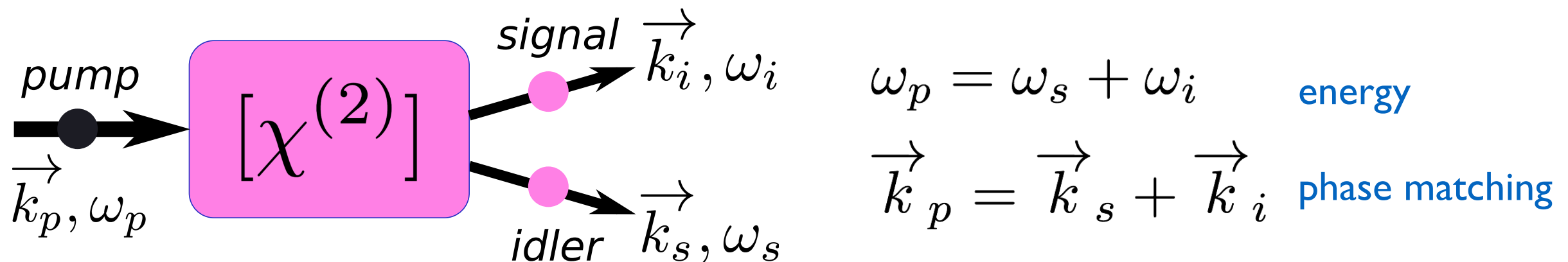
$$\omega_p = \omega_s + \omega_i \quad \text{energy}$$

$$\vec{k}_p = \vec{k}_s + \vec{k}_i \quad \text{phase matching}$$

Photon-pair emission \rightarrow Poissonian distribution (pump laser)

2.3 Heralded SPS based on nonlinear optics

► Spontaneous parametric downconversion (SPDC) in nonlinear optics



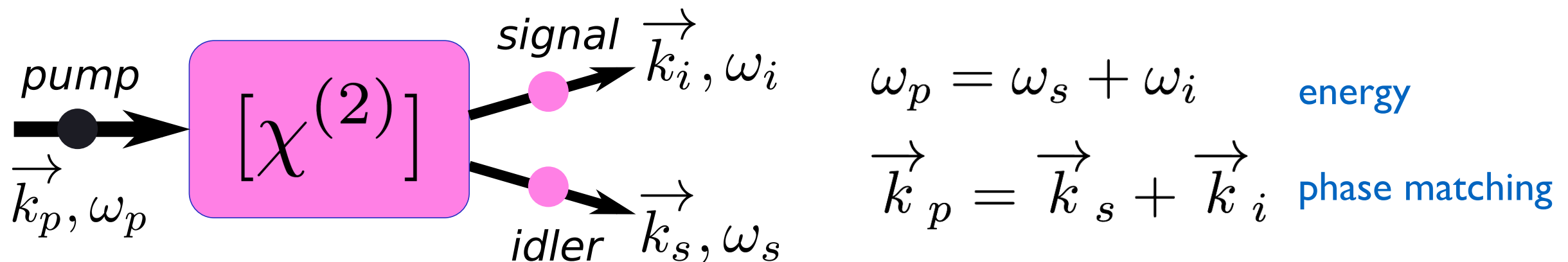
Photon-pair emission → Poissonian distribution (pump laser)

► Single photon emission character ?

- Emission of non degenerate paired photons by SPDC emitted simultaneously
- Detection of the signal photon (shorter λ_s) for heralding the idler photon (λ_i)

2.3 Heralded SPS based on nonlinear optics

▶ Spontaneous parametric downconversion (SPDC) in nonlinear optics



Photon-pair emission \rightarrow Poissonian distribution (pump laser)

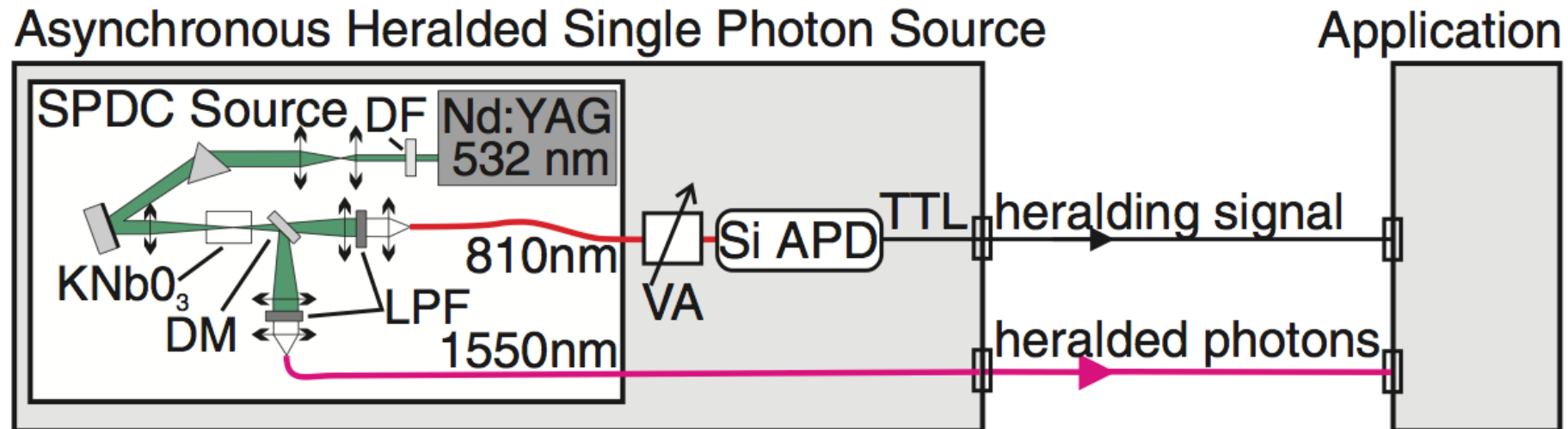
▶ Single photon emission character ?

- Emission of non degenerate paired photons by SPDC emitted simultaneously
- Detection of the signal photon (shorter λ_s) for heralding the idler photon (λ_i)

- \rightarrow Reduces both empty pulses and detection noise
- \rightarrow Reduces the emission statistics to **sub-Poissonian distributions (antibunching)**
- \rightarrow Emission regime is asynchronous !
- \rightarrow Prevents from observing standard antibunching curves !

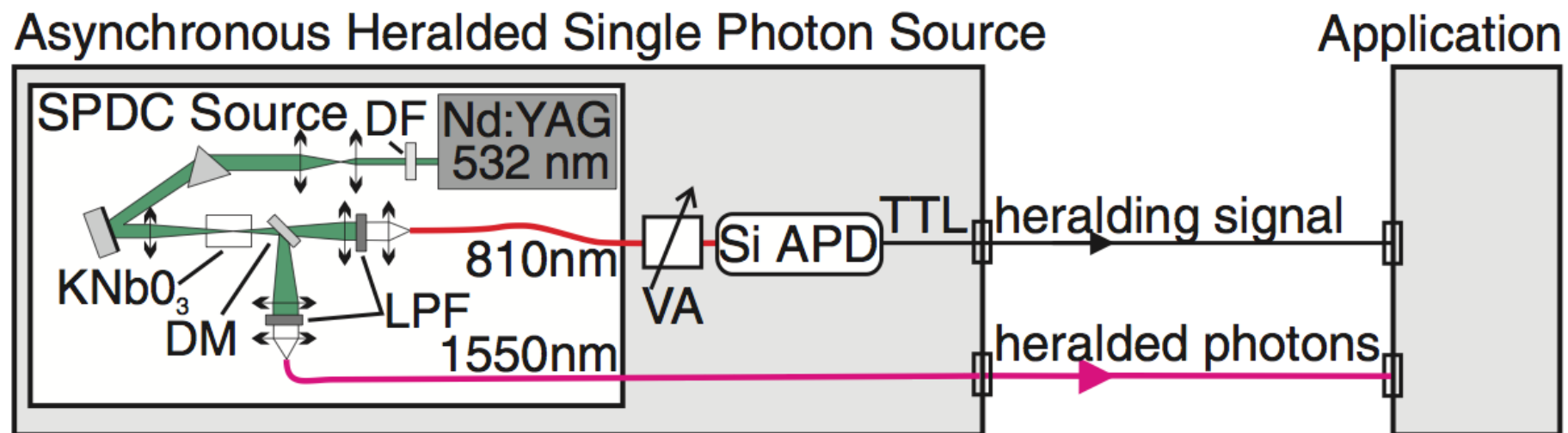
2.3 Heralded SPS based on nonlinear optics

► Standard setup based on a bulk crystal

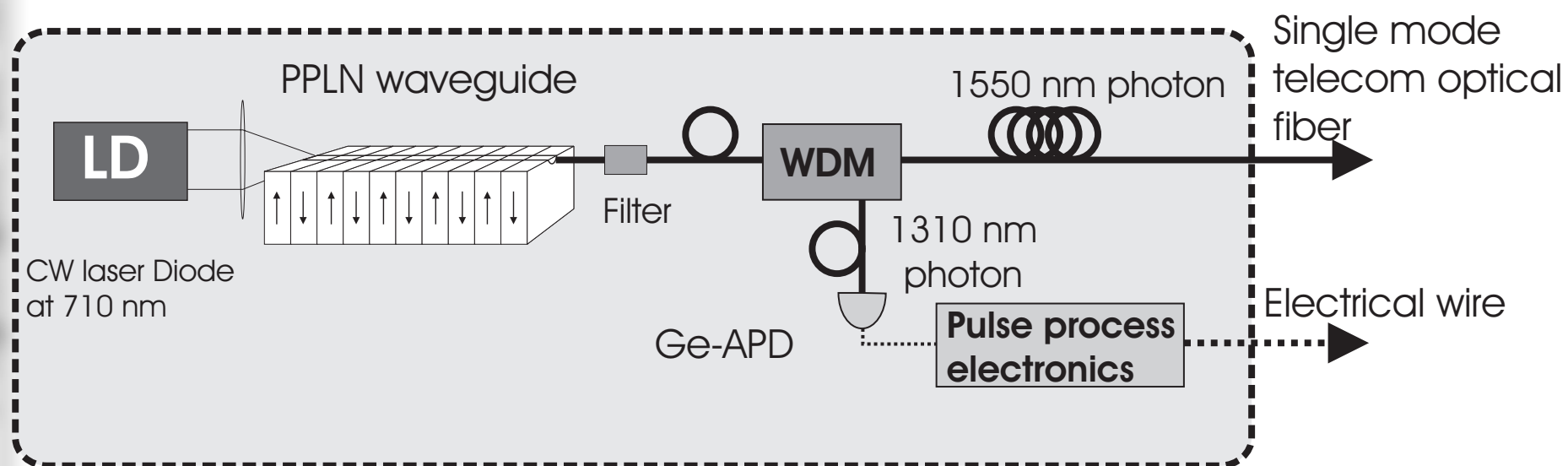


2.3 Heralded SPS based on nonlinear optics

▶ Standard setup based on a bulk crystal



▶ Standard setup based on a waveguide crystal

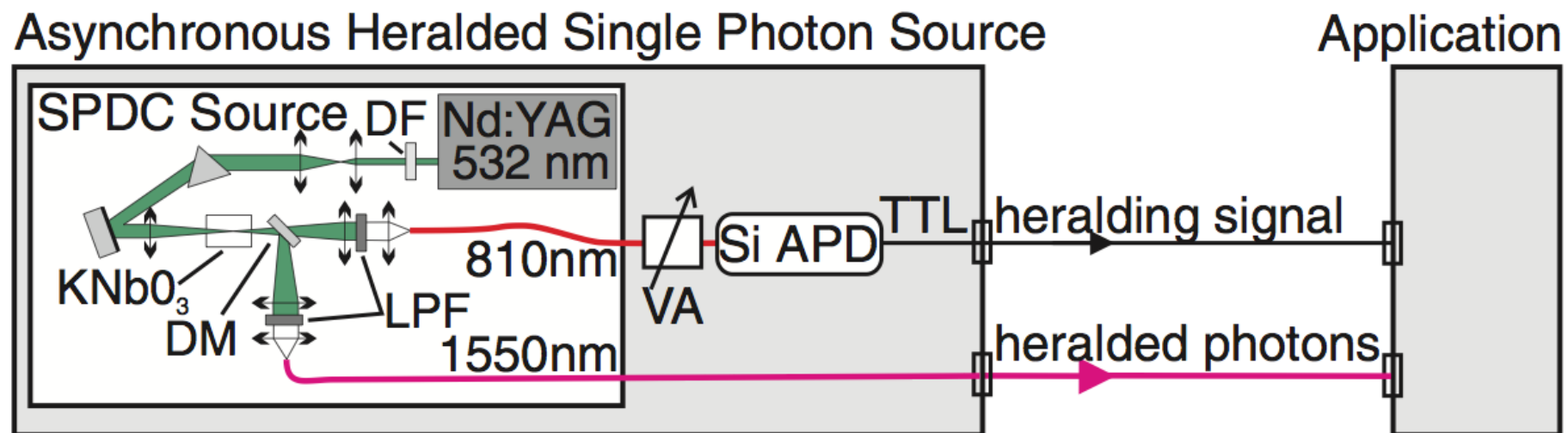


[Gisin's group, Geneva] S. Fasel *et al.*, NJP 6, 163 (2004)

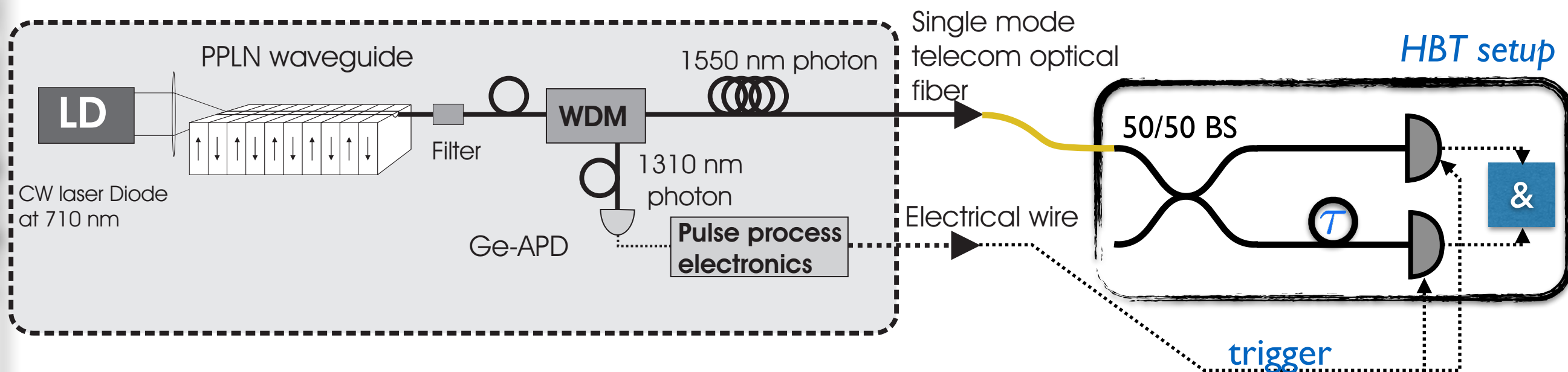
[Tanzilli's group, Nice] O. Alibart *et al.*, Opt. Lett. 30, 1539 (2005)

2.3 Heralded SPS based on nonlinear optics

▶ Standard setup based on a bulk crystal



▶ Standard setup based on a waveguide crystal



[Gisin's group, Geneva] S. Fasel *et al.*, NJP 6, 163 (2004)

[Tanzilli's group, Nice] O. Alibart *et al.*, Opt. Lett. 30, 1539 (2005)

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

- ▶ Retrieving $g^{(2)}(0)$ from the experimental data

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

- ▶ Retrieving $g^{(2)}(0)$ from the experimental data

$$\text{Single photon emission rate} = \eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$$

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

► Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate = $\eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$ ←..... heralding rate

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

► Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate = $\eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$ ←..... heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

▶ Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate $= \eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$
← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$
← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

▶ Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate $= \eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$
← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$
← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

with $\bar{n} = \eta_{coll}\mu\Delta T$

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

► Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate = $\eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$ ← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$ ← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

with $\bar{n} = \eta_{coll}\mu\Delta T$

Proba. of having 2 photons $P_2 \approx \eta_{coll}^2 \mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)$

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

► Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate = $\eta_{coll} \times \eta_{det} \times \left(N_{raw}^H - N_{DC}^H \right)$ ← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$ ← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

with $\bar{n} = \eta_{coll}\mu\Delta T$

Proba. of having 2 photons $P_2 \approx \eta_{coll}^2 \mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)$

Proba. of having 1 photon $P_1 \approx \eta_{coll} \left\{ \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right) \left[1 - 2\eta_{coll}\mu\Delta T \right] + \mu\Delta T \right\}$

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

► Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate = $\eta_{coll} \times \eta_{det} \times \left(N_{raw}^H - N_{DC}^H \right)$ ← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$ ← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

with $\bar{n} = \eta_{coll}\mu\Delta T$

Proba. of having 2 photons $P_2 \approx \eta_{coll}^2 \mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)$

Proba. of having 1 photon $P_1 \approx \eta_{coll} \left\{ \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right) \left[1 - 2\eta_{coll}\mu\Delta T \right] + \mu\Delta T \right\}$

Autocorr. function

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

▶ Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate = $\eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$ ← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$ ← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

with $\bar{n} = \eta_{coll}\mu\Delta T$

Proba. of having 2 photons $P_2 \approx \eta_{coll}^2 \mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)$

Proba. of having 1 photon $P_1 \approx \eta_{coll} \left\{ \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right) [1 - 2\eta_{coll}\mu\Delta T] + \mu\Delta T \right\}$

Autocorr. function $g^{(2)}(0) = \frac{2P_2}{P_1^2}$

Assuming a Poissonian distrib.

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

► Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate = $\eta_{coll} \times \eta_{det} \times (N_{raw}^H - N_{DC}^H)$ ← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$ ← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

with $\bar{n} = \eta_{coll}\mu\Delta T$

Proba. of having 2 photons $P_2 \approx \eta_{coll}^2 \mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)$

Proba. of having 1 photon $P_1 \approx \eta_{coll} \left\{ \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right) [1 - 2\eta_{coll}\mu\Delta T] + \mu\Delta T \right\}$

Autocorr. function $g^{(2)}(0) = \frac{2P_2}{P_1^2} \approx \frac{2\mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)}{\left(\mu\Delta T + [1 - 2\eta_{coll}\mu\Delta T] \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right) \right)^2}$

Assuming a Poissonian distrib.

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements

▶ Retrieving $g^{(2)}(0)$ from the experimental data

Single photon emission rate $= \eta_{coll} \times \eta_{det} \times \left(N_{raw}^H - N_{DC}^H \right)$ ← heralding rate

Proba. of having no photon $P_{\bar{n}}(n = 0) = e^{-\eta_{coll}\mu\Delta T}$ ← Poissonian distrib.
 μ : SPDC em. rate
 ΔT : det. time window

with $\bar{n} = \eta_{coll}\mu\Delta T$

Proba. of having 2 photons $P_2 \approx \eta_{coll}^2 \mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)$

Proba. of having 1 photon $P_1 \approx \eta_{coll} \left\{ \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right) \left[1 - 2\eta_{coll}\mu\Delta T \right] + \mu\Delta T \right\}$

Autocorr. function $g^{(2)}(0) = \frac{2P_2}{P_1^2} \approx \frac{2\mu \Delta T \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right)}{\left(\mu\Delta T + \left[1 - 2\eta_{coll}\mu\Delta T \right] \left(\frac{N_{raw}^H - N_{DC}^H}{N_{raw}^H} \right) \right)^2}$

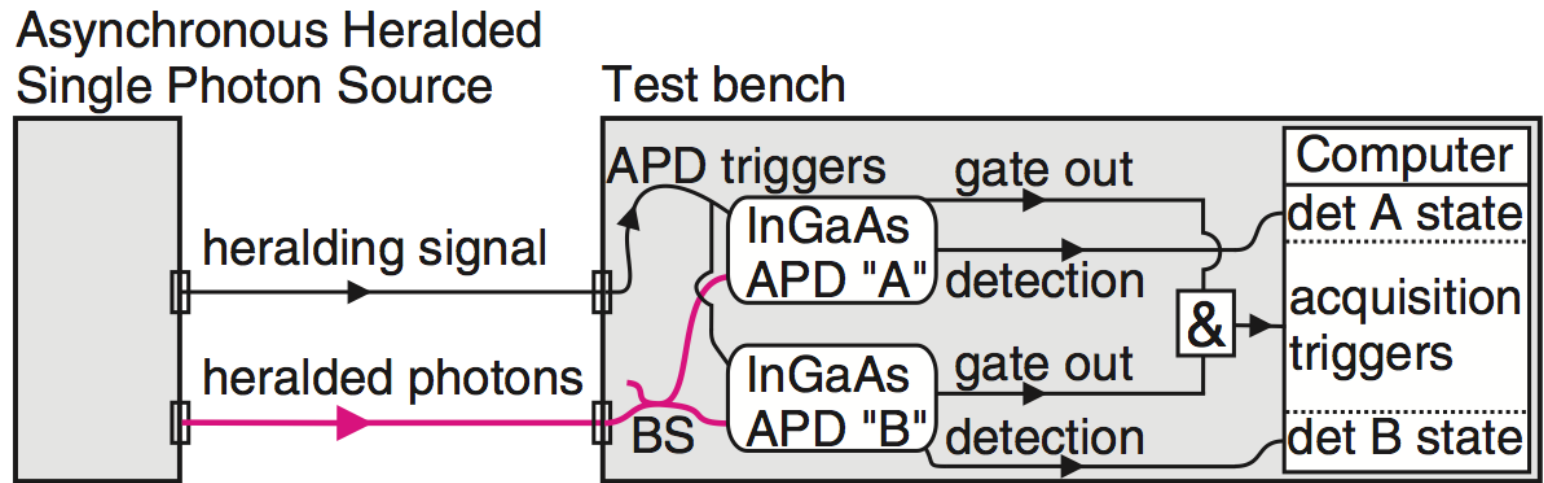
Assuming a Poissonian distrib.

Typical experimental parameters: $\eta_{coll} \sim 0.5$, $\Delta T \sim 1$ ns, $\mu \sim 10^6$ pairs/s

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements → a much clever approach

▶ Retrieving $g^{(2)}(0)$ from post data processing

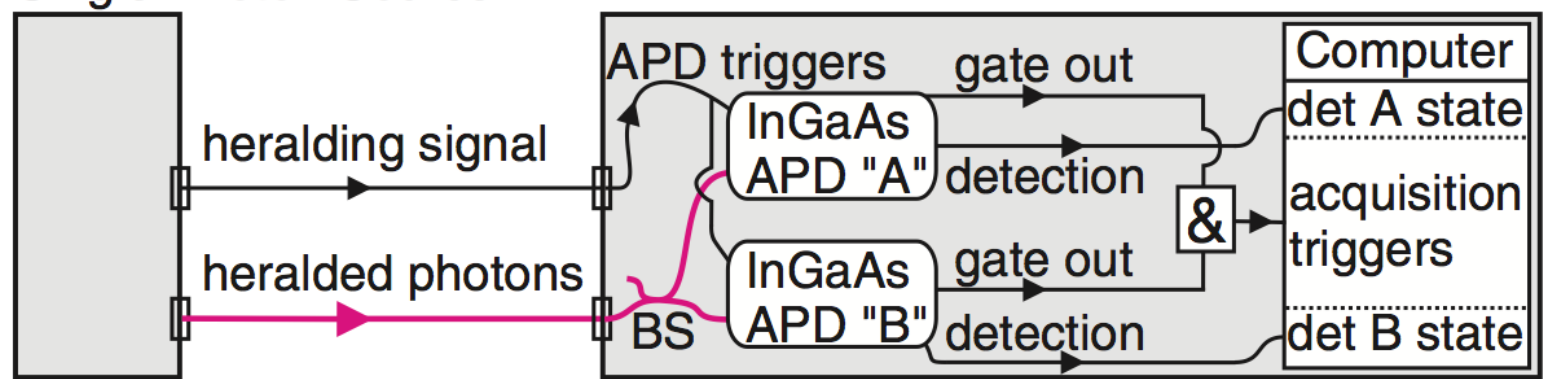


2.3 Heralded SPS based on nonlinear optics

Antibunching measurements → a much clever approach

► Retrieving $g^{(2)}(0)$ from post data processing

Asynchronous Heralded Single Photon Source



no detection
start at A
no detection
invalid start
no detection
stop at B

det A	det B
...	...
0	0
1	0
0	0
1	0
0	0
0	1
...	...

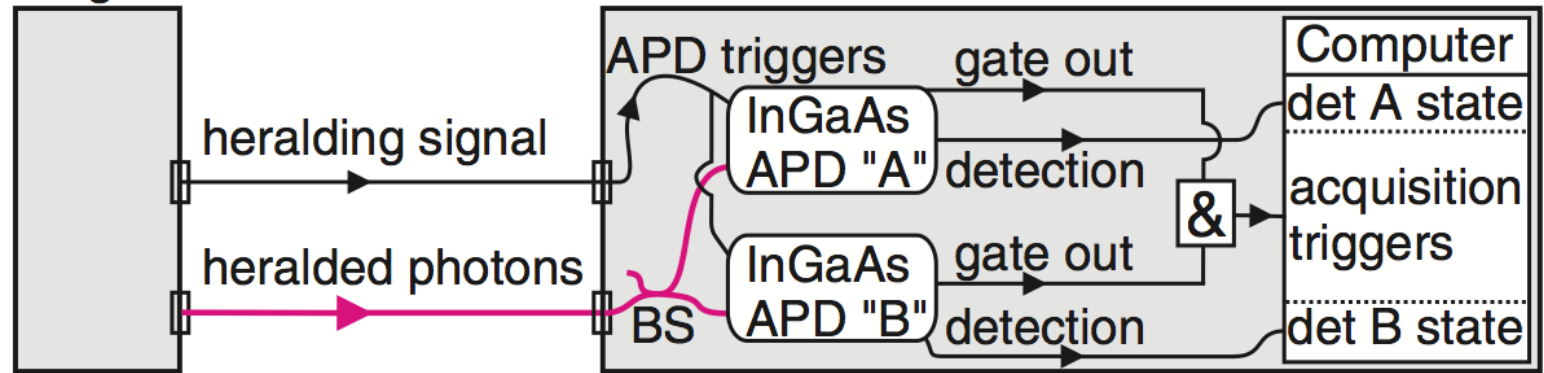
$n = -1$
 $n = 0$
 $n = +1$
 $n = +2$
 $n = +3$
 $n = +4$

2.3 Heralded SPS based on nonlinear optics

Antibunching measurements → a much clever approach

► Retrieving $g^{(2)}(0)$ from post data processing

Asynchronous Heralded Single Photon Source

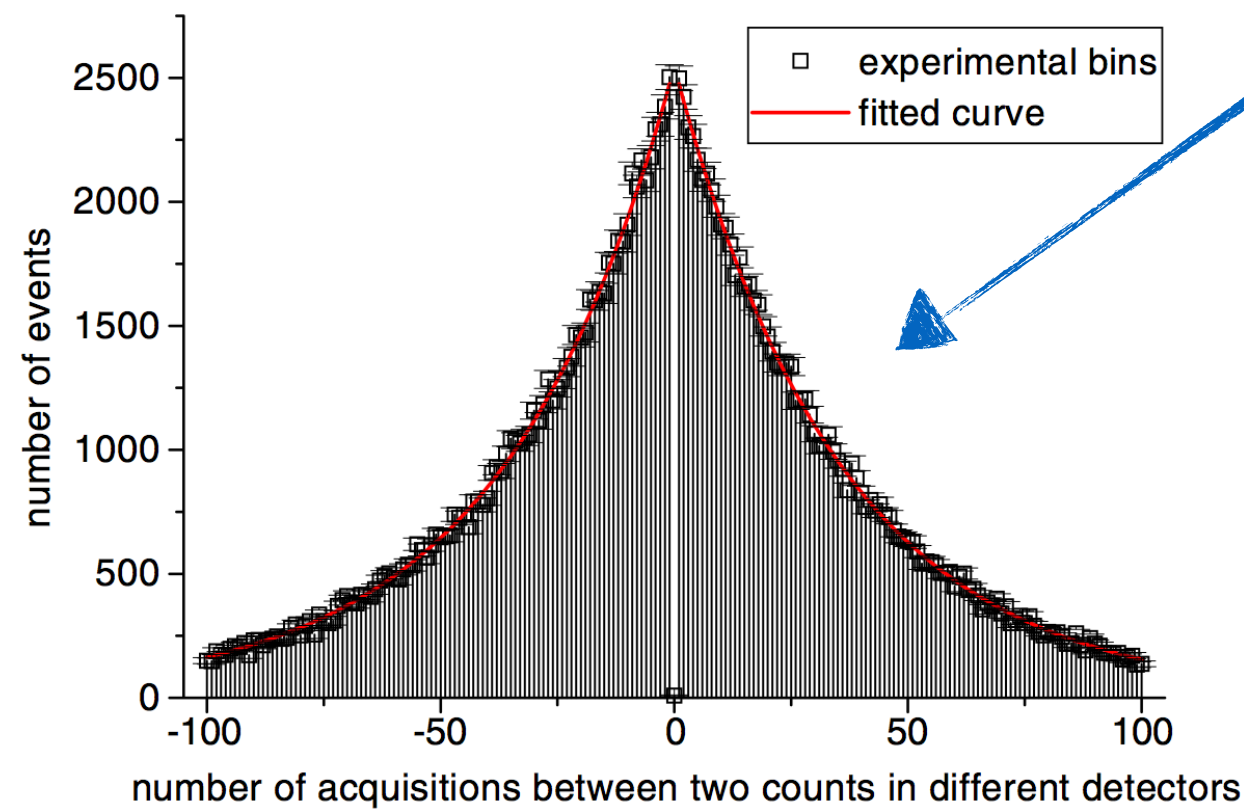


no detection
start at A
no detection
invalid start
no detection
stop at B

det A	det B
...	...
0	0
1	0
0	0
1	0
0	0
0	1
...	...

$n = -1$
 $n = 0$
 $n = +1$
 $n = +2$
 $n = +3$
 $n = +4$

► Results

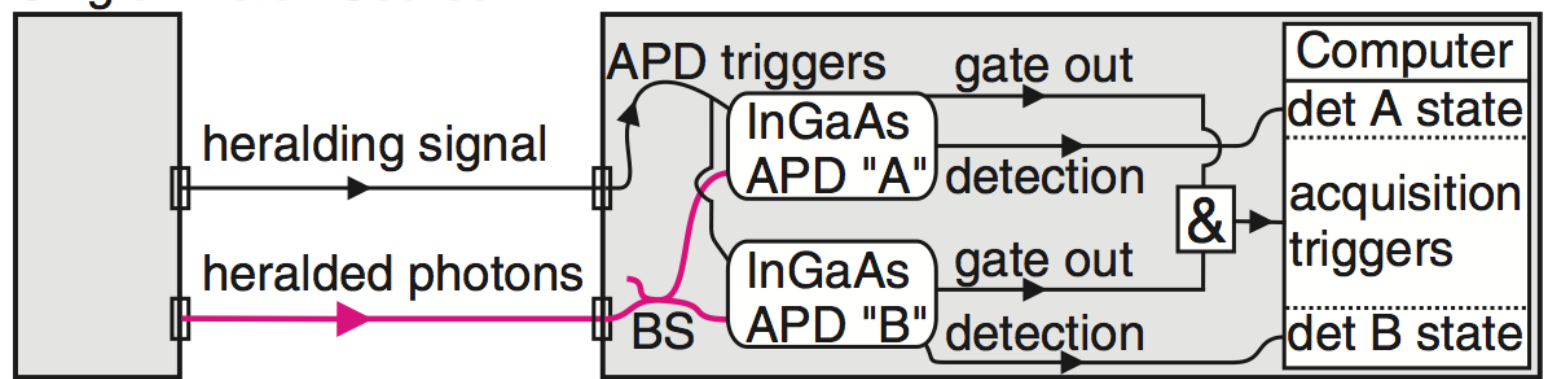


2.3 Heralded SPS based on nonlinear optics

Antibunching measurements → a much clever approach

▶ Retrieving $g^{(2)}(0)$ from post data processing

Asynchronous Heralded Single Photon Source



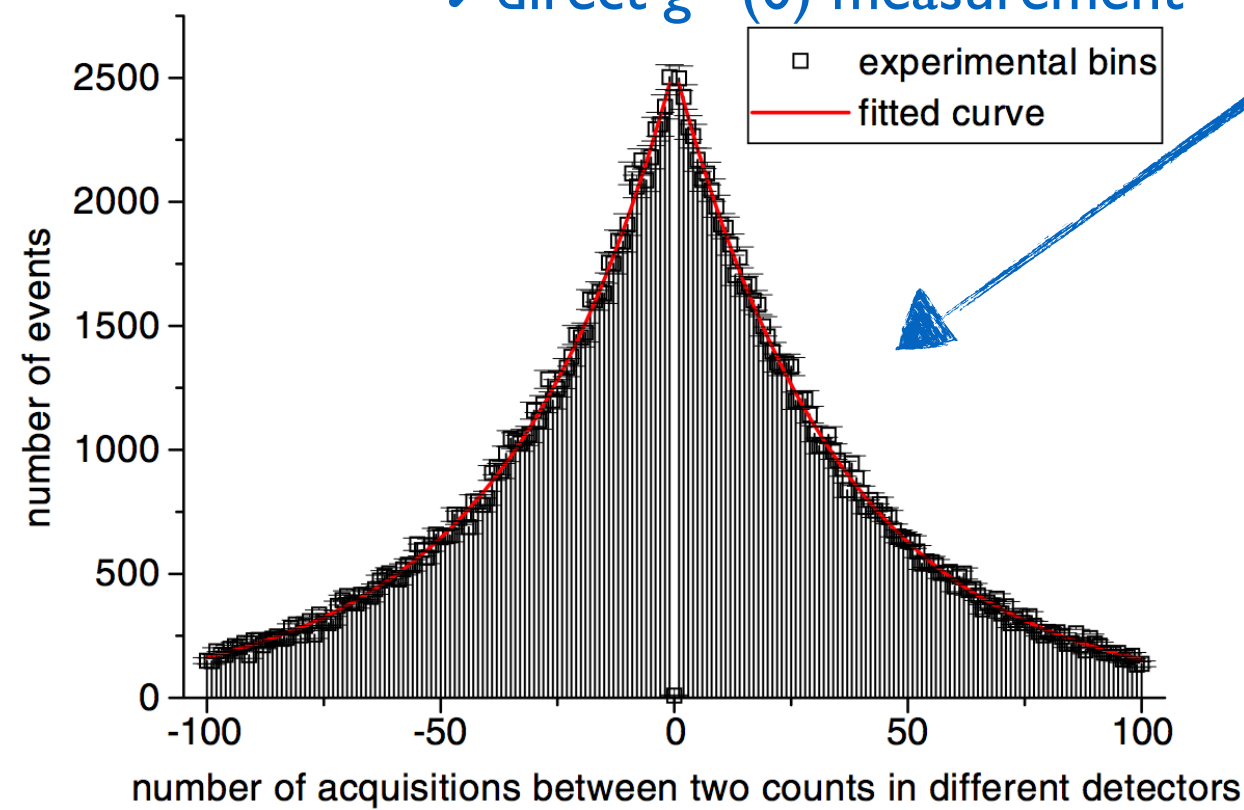
no detection
start at A
no detection
invalid start
no detection
stop at B

det A	det B
...	...
0	0
1	0
0	0
1	0
0	0
0	1
...	...

$n = -1$
 $n = 0$
 $n = +1$
 $n = +2$
 $n = +3$
 $n = +4$

▶ Results Same as obtained with "true" SPS

→ direct $g^{(2)}(0)$ measurement

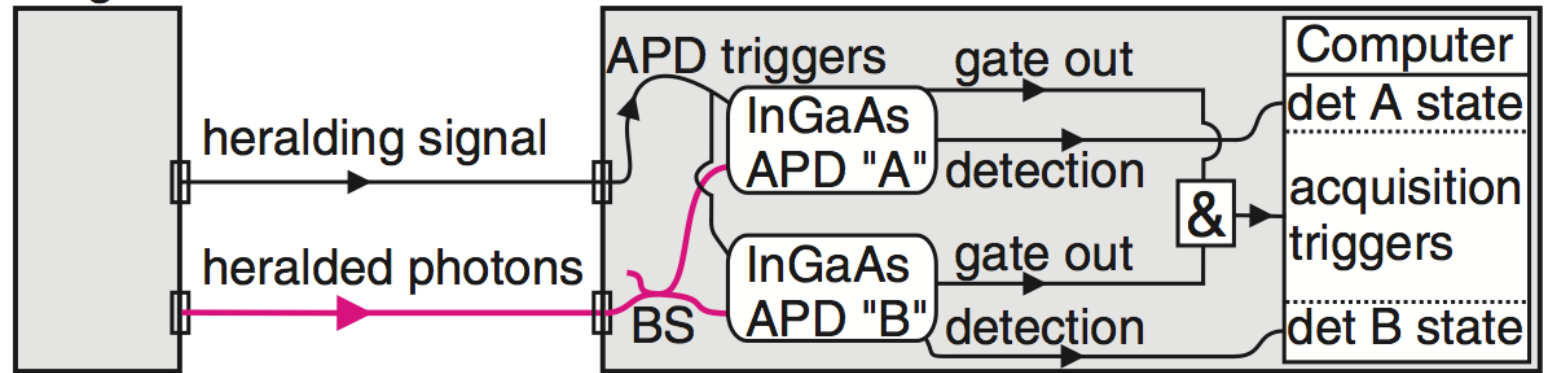


2.3 Heralded SPS based on nonlinear optics

Antibunching measurements → a much clever approach

▶ Retrieving $g^{(2)}(0)$ from post data processing

Asynchronous Heralded Single Photon Source

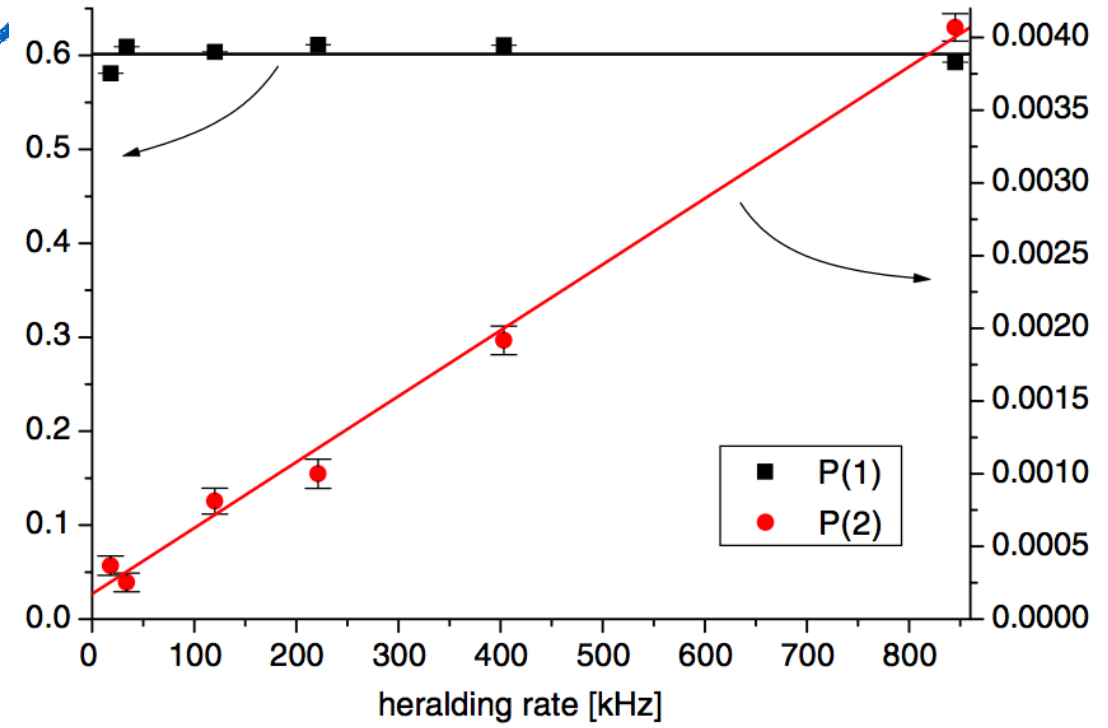
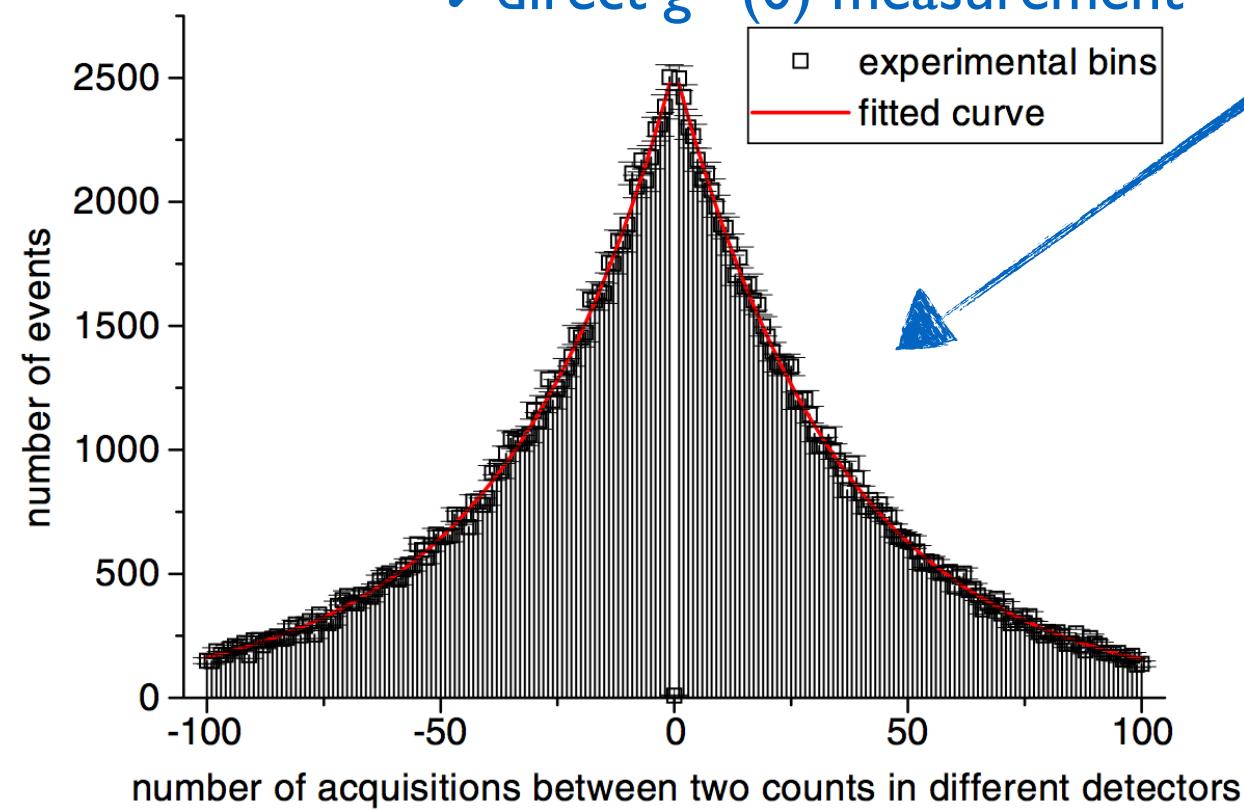


no detection
start at A
no detection
invalid start
no detection
stop at B

det A	det B
...	...
0	0
1	0
0	0
1	0
0	0
0	1
...	...

$n = -1$
 $n = 0$
 $n = +1$
 $n = +2$
 $n = +3$
 $n = +4$

▶ Results Same as obtained with "true" SPS → direct $g^{(2)}(0)$ measurement



2.3 Heralded SPS based on nonlinear optics

Summary and future directions

► Summary

Source	λ (nm)	P0	P1	P2	g
Waveguide	1550 / 1310	0.63	0.37	7.10	0.08
Bulk	1550 / 810	0.39	0.61	2.10	2.10

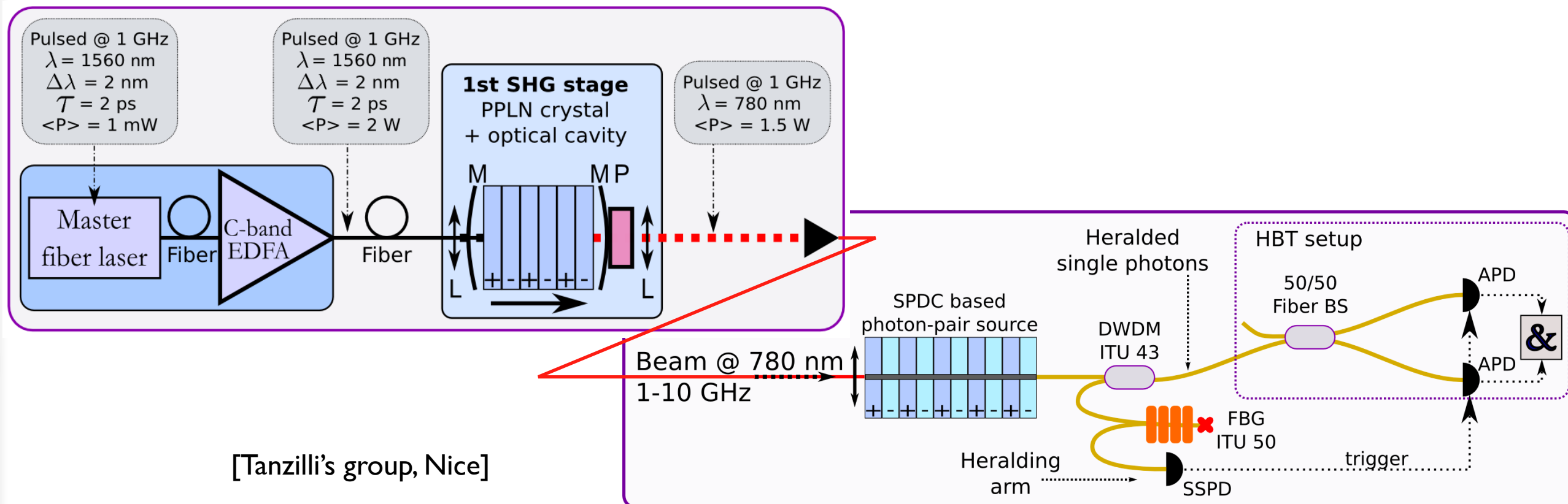
2.3 Heralded SPS based on nonlinear optics

Summary and future directions

► Summary

Source	λ (nm)	P0	P1	P2	g
Waveguide	1550 / 1310	0.63	0.37	7.10	0.08
Bulk	1550 / 810	0.39	0.61	2.10	2.10

► Towards Ultrafast heralding rates \rightarrow 10 MHz



2.3 Heralded SPS based on nonlinear optics

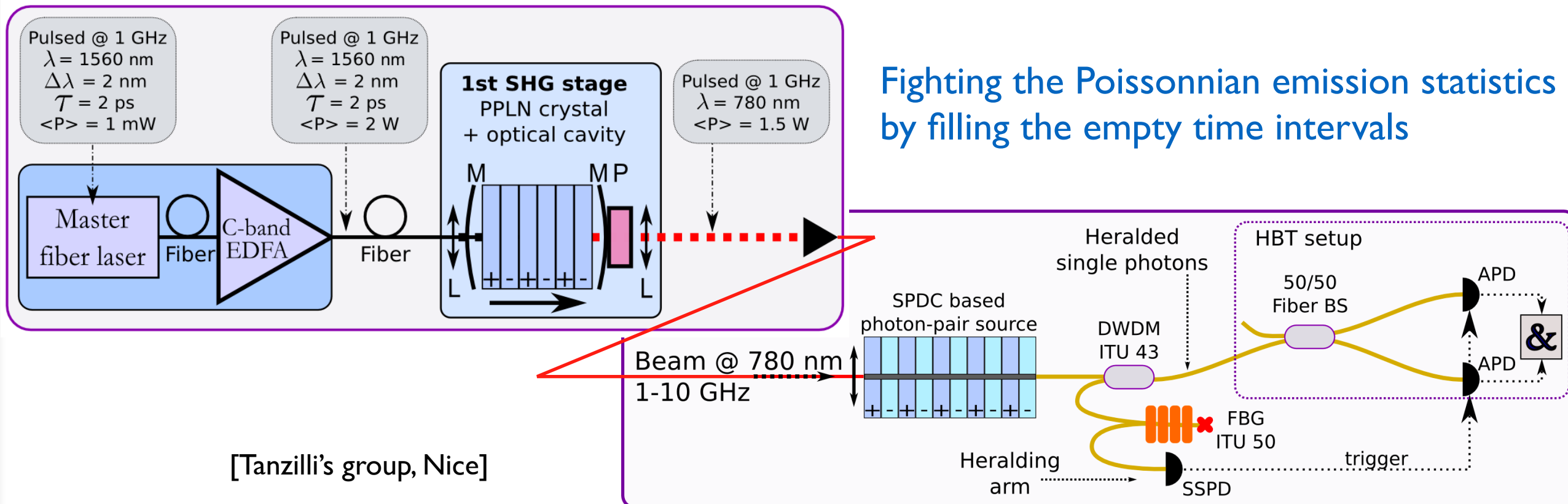
Summary and future directions

► Summary

Source	λ (nm)	P0	P1	P2	g
Waveguide	1550 / 1310	0.63	0.37	7.10	0.08
Bulk	1550 / 810	0.39	0.61	2.10	2.10

► Towards Ultrafast heralding rates \rightarrow 10 MHz

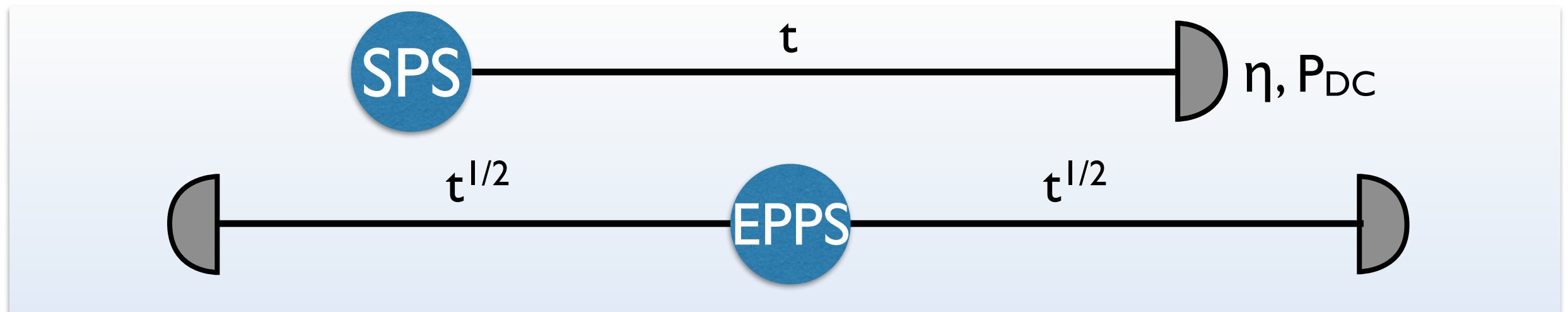
Fighting the Poissonian emission statistics by filling the empty time intervals



2.4 Connecting HSPs for quantum networks

► The context of today's quantum communication

Distribution of quantum bits of information
using single photon and entangled photon pair sources

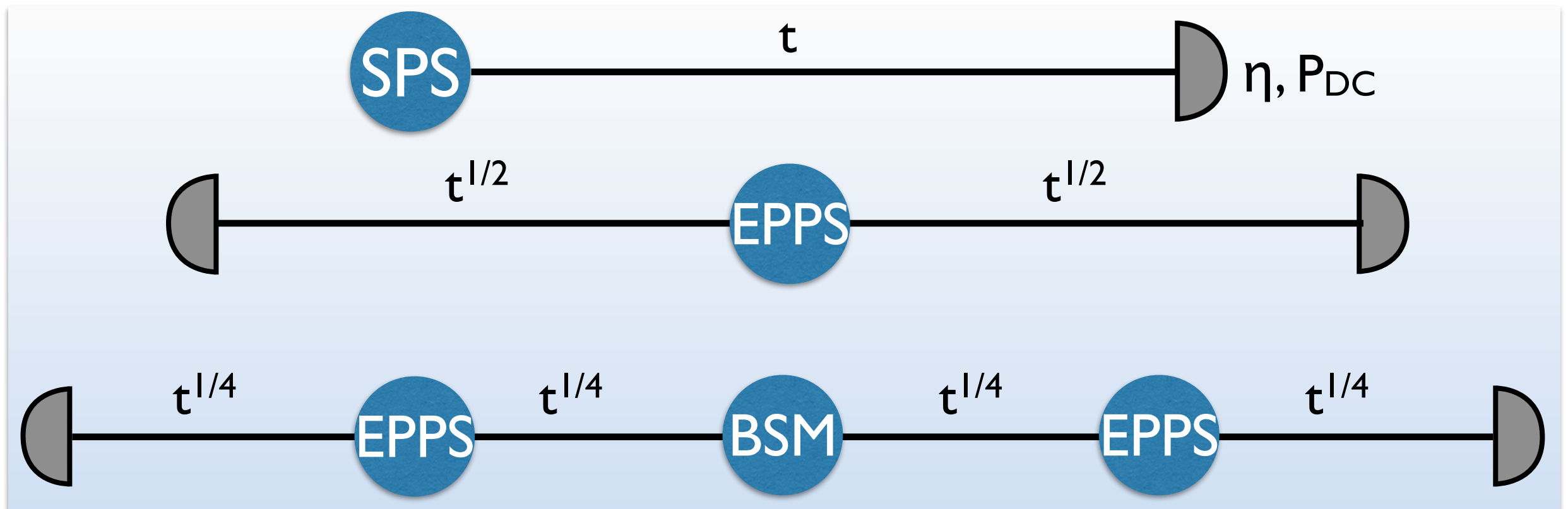


2.4 Connecting HSPs for quantum networks

► The context of today's quantum communication

Distribution of quantum bits of information
using single photon and entangled photon pair sources

over long distances

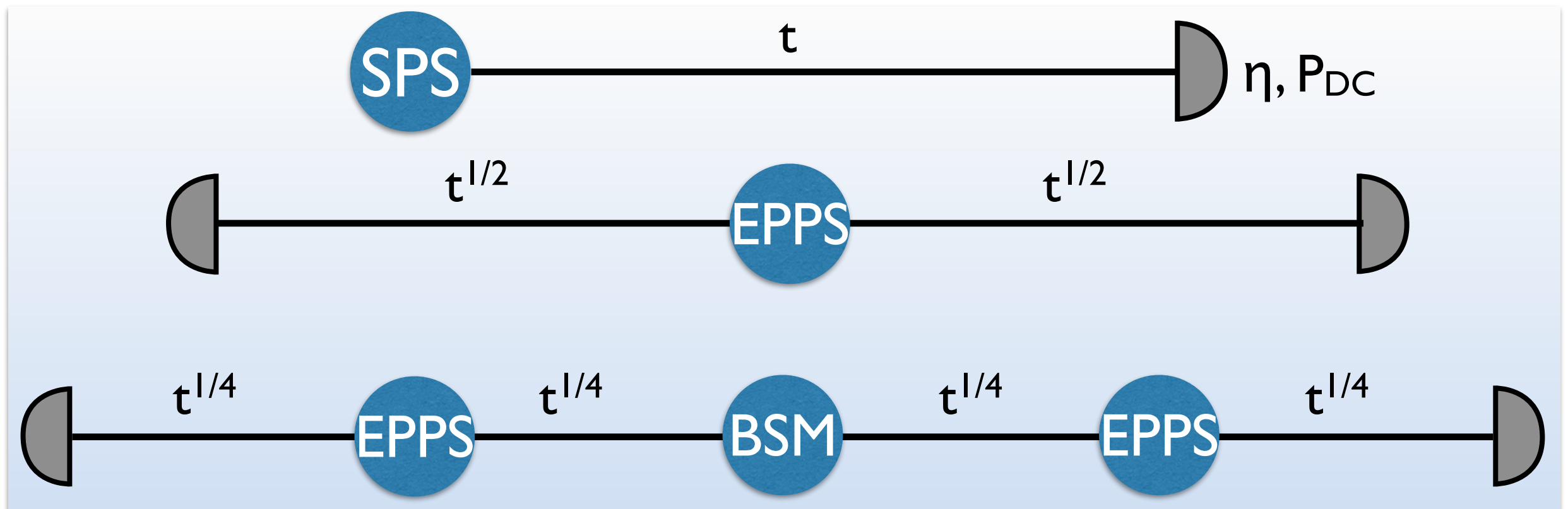


2.4 Connecting HSPs for quantum networks

► The context of today's quantum communication

Distribution of quantum bits of information
using single photon and entangled photon pair sources

over long distances



► BSM based on a HBT type setup → Hong-Ou-Mandel (HOM)

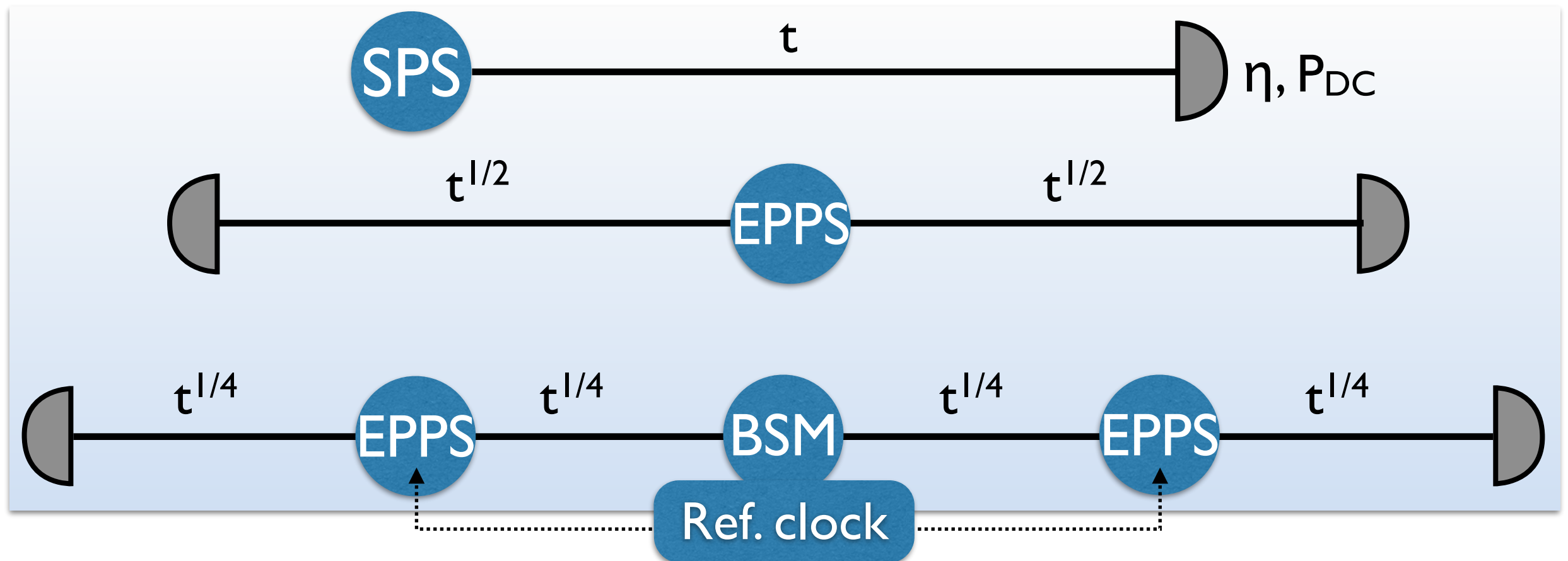
Referred to as two-photon interference

2.4 Connecting HSPs for quantum networks

► The context of today's quantum communication

Distribution of quantum bits of information
using single photon and entangled photon pair sources

over long distances



► BSM based on a HBT type setup → Hong-Ou-Mandel (HOM)

Referred to as two-photon interference

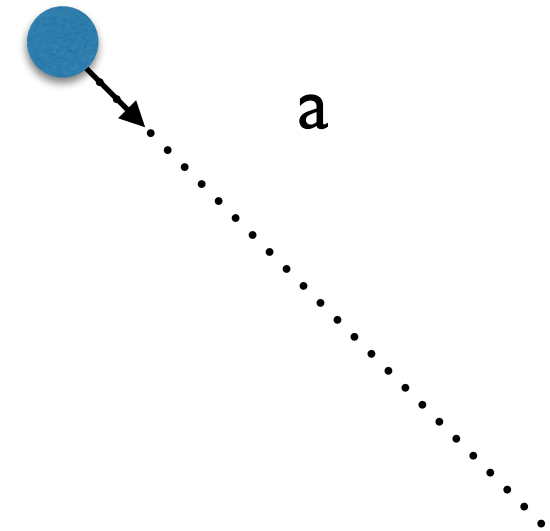
→ Needs a synchronization procedure, and identical photons (!)

2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode a →

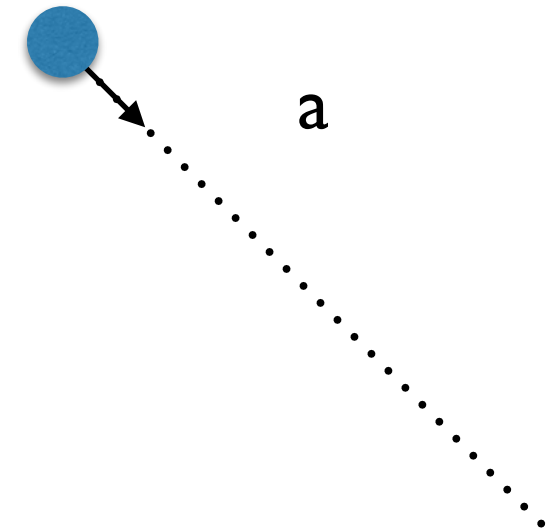
2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode a →



2.4 Connecting HSPs: how does the HOM effect work ?

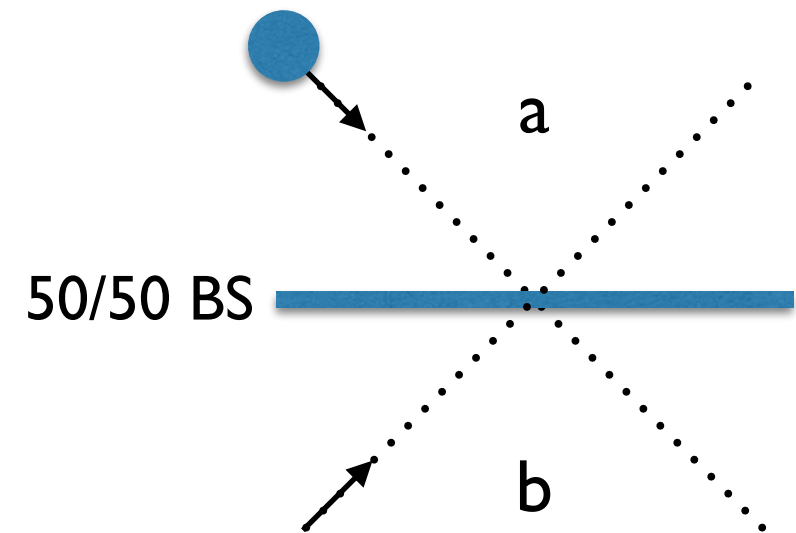
One photon in spatial mode $a \rightarrow a^\dagger$



2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode $a \rightarrow a^\dagger$

Action of a BS for a single photon

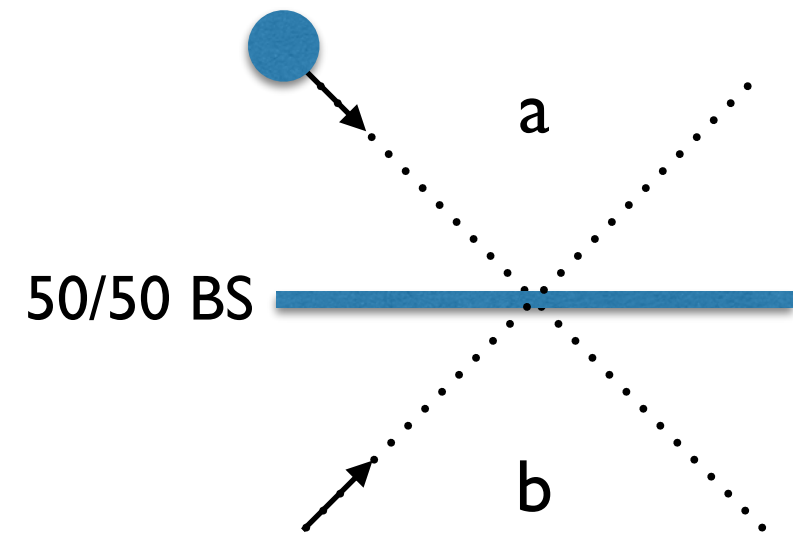


2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode $a \rightarrow a^\dagger$

Action of a BS for a single photon

$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger)$$

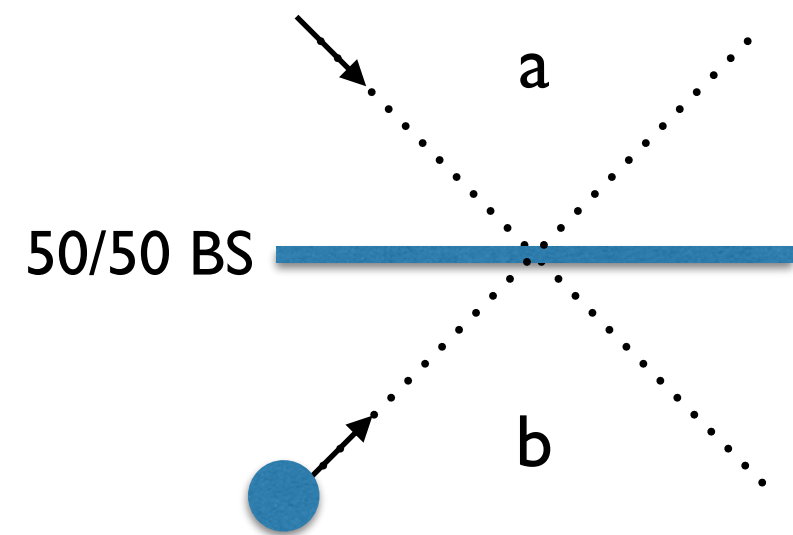


2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode $a \rightarrow a^\dagger$

Action of a BS for a single photon

$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$



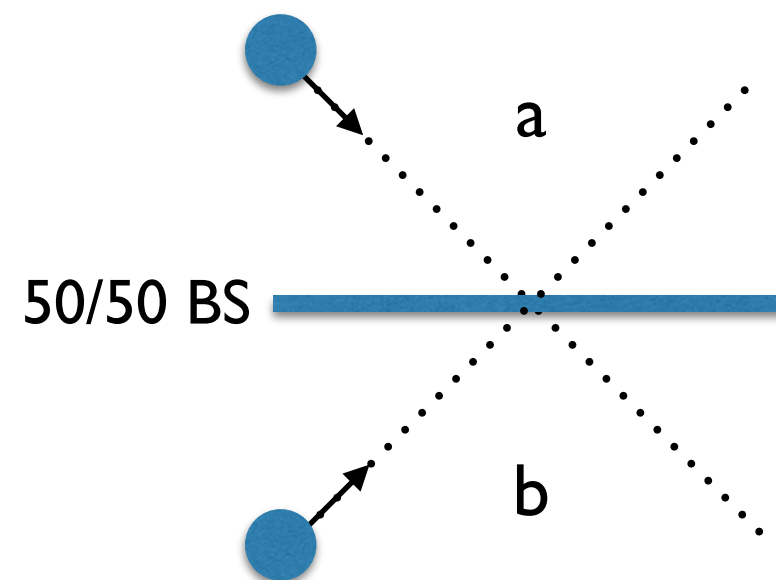
2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode $a \rightarrow a^\dagger$

Action of a BS for a single photon

$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

Action of a BS for two single photons



2.4 Connecting HSPs: how does the HOM effect work ?

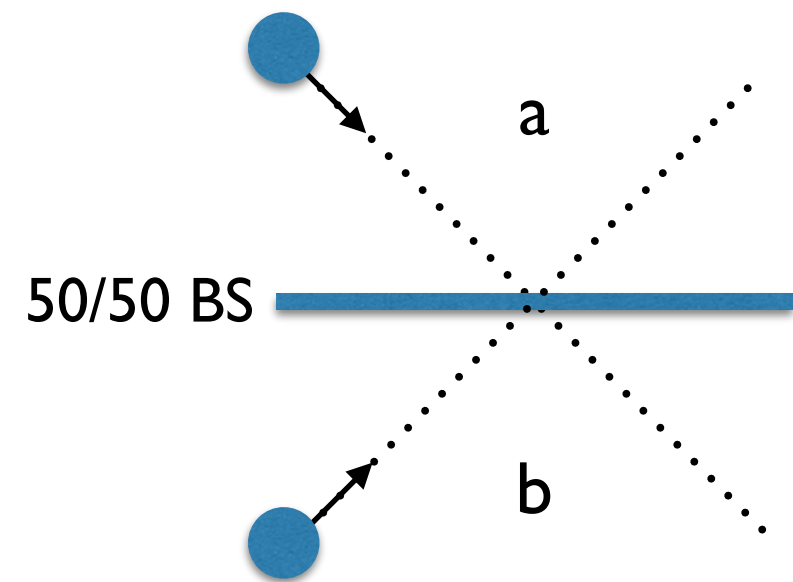
One photon in spatial mode $a \rightarrow a^\dagger$

Action of a BS for a single photon

$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

Action of a BS for two single photons

$$a^\dagger b^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$



2.4 Connecting HSPs: how does the HOM effect work ?

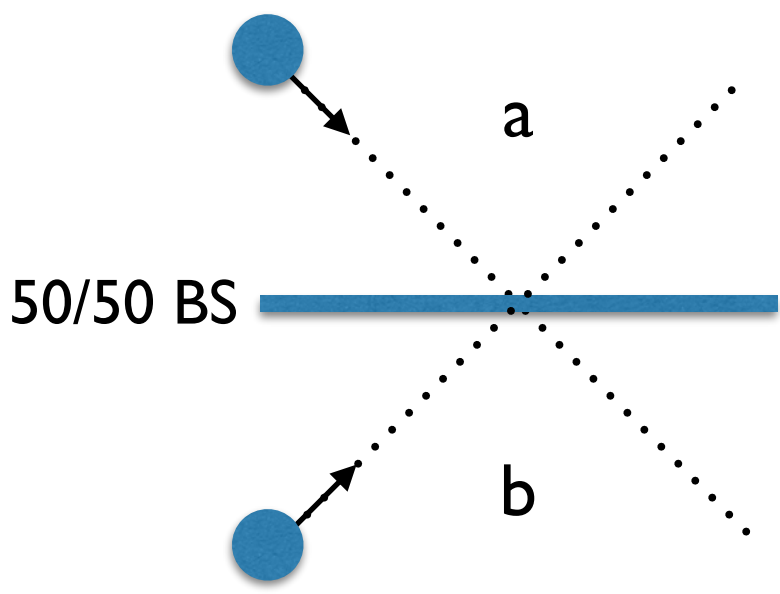
One photon in spatial mode a $\rightarrow a^\dagger$

Action of a BS for a single photon

$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

Action of a BS for two single photons

$$\begin{aligned}
 a^\dagger b^\dagger &\mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger) \\
 &\mapsto \frac{1}{2} (ia^{\dagger 2} + ib^{\dagger 2} + b^\dagger a^\dagger - a^\dagger b^\dagger)
 \end{aligned}$$



2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode a $\rightarrow a^\dagger$

Action of a BS for a single photon

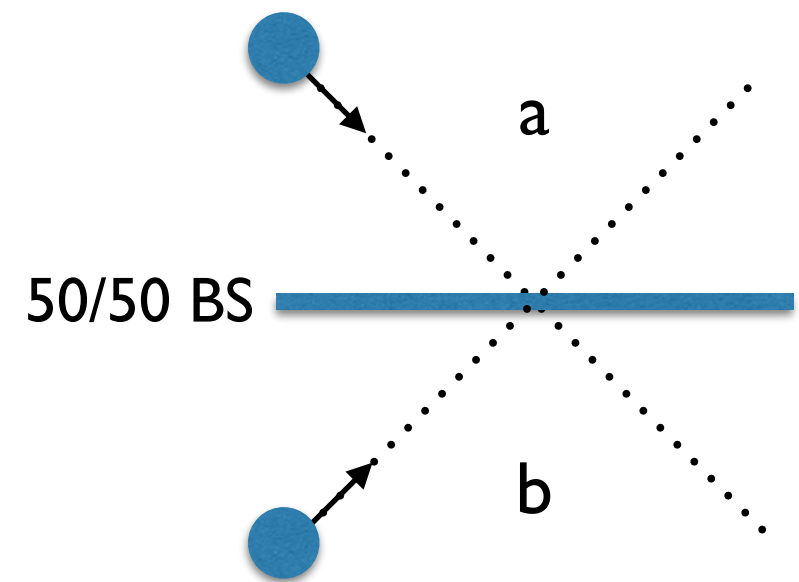
$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

Action of a BS for two single photons

$$a^\dagger b^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

$$\mapsto \frac{1}{2} (ia^{\dagger 2} + ib^{\dagger 2} + \underbrace{b^\dagger a^\dagger}_{\text{Indistinguishable}} - \underbrace{a^\dagger b^\dagger}_{\text{Indistinguishable}})$$

Indistinguishable \rightarrow a,b commute \rightarrow cancel



2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode a $\rightarrow a^\dagger$

Action of a BS for a single photon

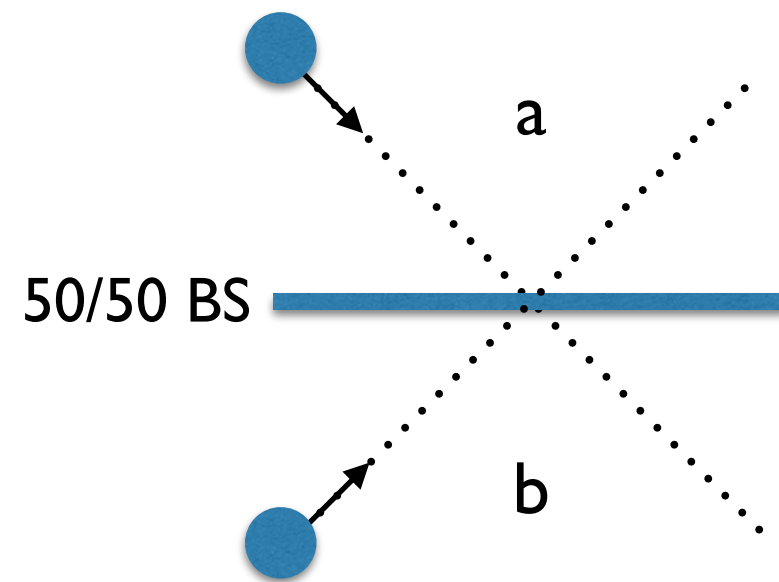
$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

Action of a BS for two single photons

$$a^\dagger b^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

$$\mapsto \frac{1}{2} (ia^{\dagger 2} + ib^{\dagger 2} + \underbrace{b^\dagger a^\dagger}_{\text{Indistinguishable}} - \underbrace{a^\dagger b^\dagger}_{\text{Indistinguishable}})$$

$$\mapsto \frac{1}{\sqrt{2}} (|2_a, 0_b\rangle - |0_a, 2_b\rangle)$$



Indistinguishable \rightarrow a,b commute \rightarrow cancel

2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode a $\rightarrow a^\dagger$

Action of a BS for a single photon

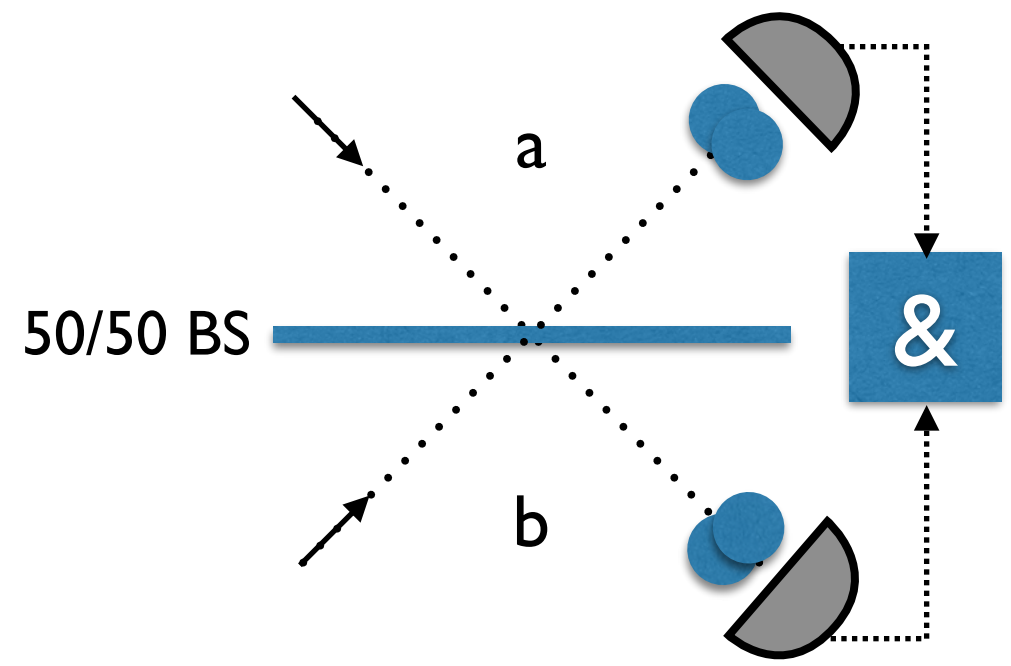
$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

Action of a BS for two single photons

$$\begin{aligned}
 a^\dagger b^\dagger &\mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger) \\
 &\mapsto \frac{1}{2} (ia^{\dagger 2} + ib^{\dagger 2} + \underbrace{b^\dagger a^\dagger}_{\text{Indistinguishable}} - \underbrace{a^\dagger b^\dagger}_{\text{Indistinguishable}})
 \end{aligned}$$

$$\mapsto \frac{1}{\sqrt{2}} (|2_a, 0_b\rangle - |0_a, 2_b\rangle)$$

\rightarrow 2 photons either in a or b \rightarrow path entanglement



Indistinguishable \rightarrow a,b commute \rightarrow cancel

2.4 Connecting HSPs: how does the HOM effect work ?

One photon in spatial mode a $\rightarrow a^\dagger$

Action of a BS for a single photon

$$a^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \quad b^\dagger \mapsto \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

Action of a BS for two single photons

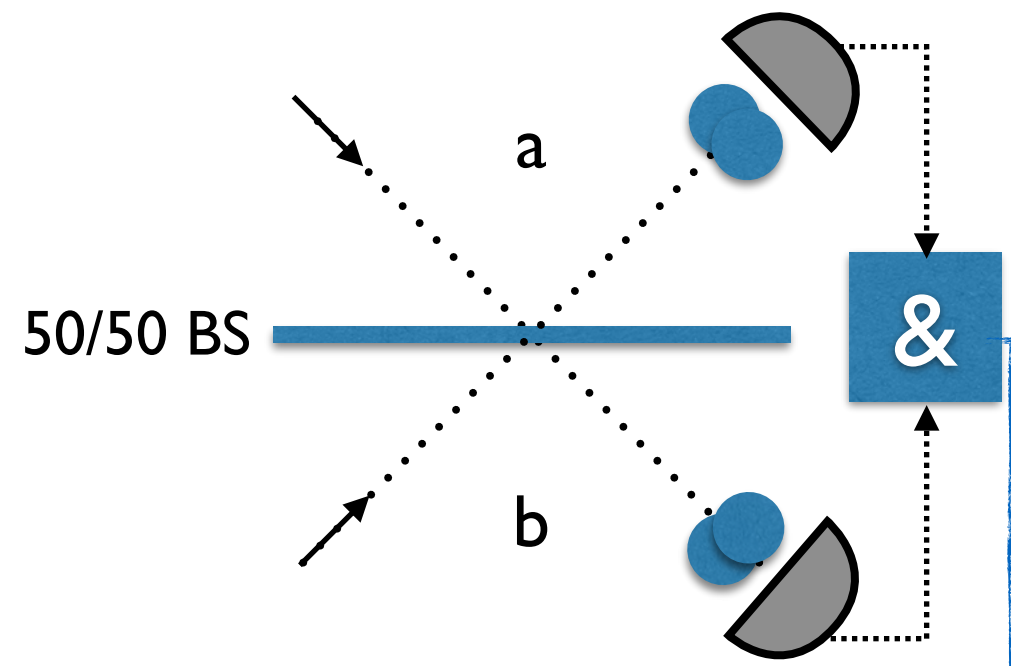
$$a^\dagger b^\dagger \mapsto \frac{1}{\sqrt{2}} (ia^\dagger + b^\dagger) \frac{1}{\sqrt{2}} (a^\dagger + ib^\dagger)$$

$$\mapsto \frac{1}{2} (ia^{\dagger 2} + ib^{\dagger 2} + \underbrace{b^\dagger a^\dagger}_{\text{Indistinguishable}} - \underbrace{a^\dagger b^\dagger}_{\text{Indistinguishable}})$$

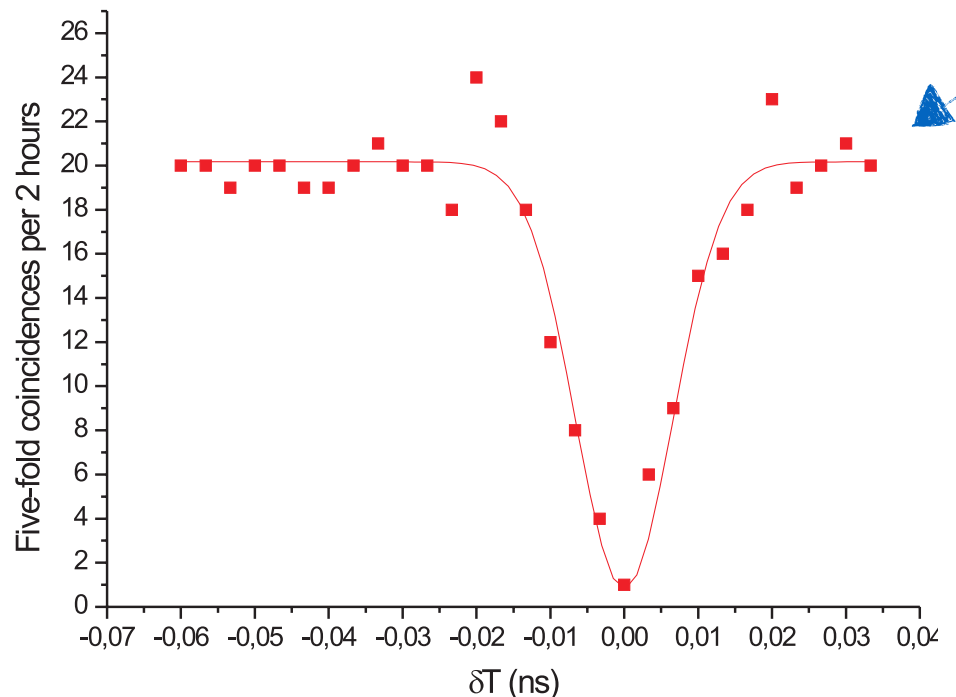
$$\mapsto \frac{1}{\sqrt{2}} (|2_a, 0_b\rangle - |0_a, 2_b\rangle)$$

\rightarrow 2 photons either in a or b \rightarrow path entanglement

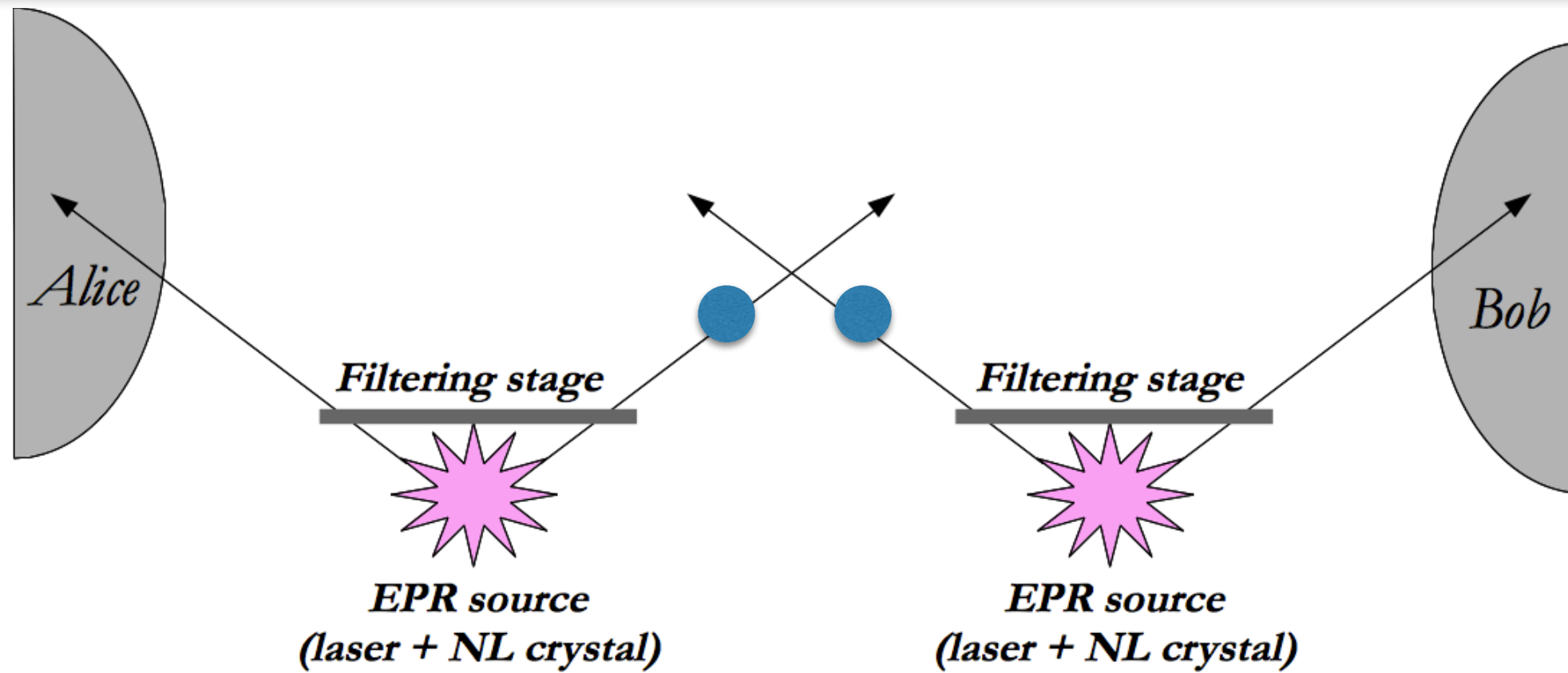
HOM effect: dip in the coincidence counts



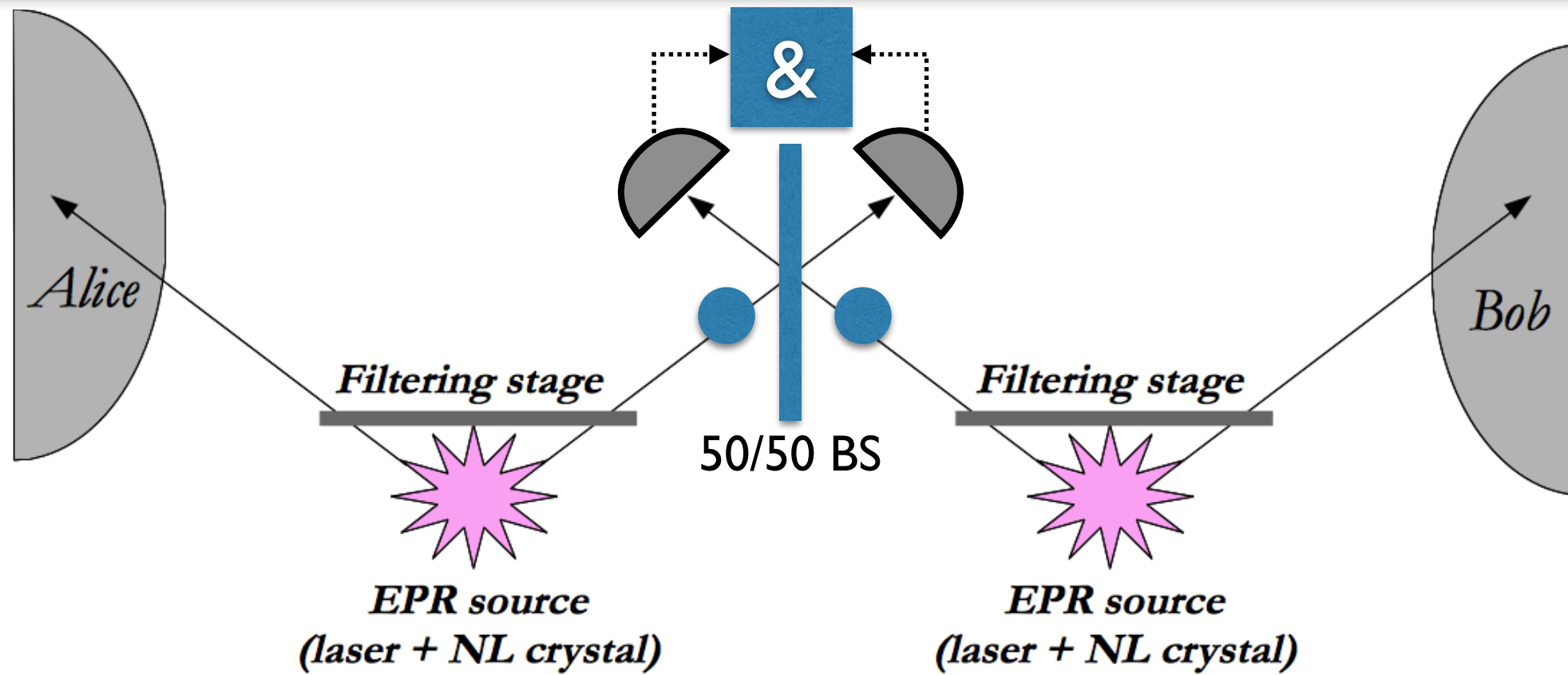
Indistinguishable \rightarrow a,b commute \rightarrow cancel



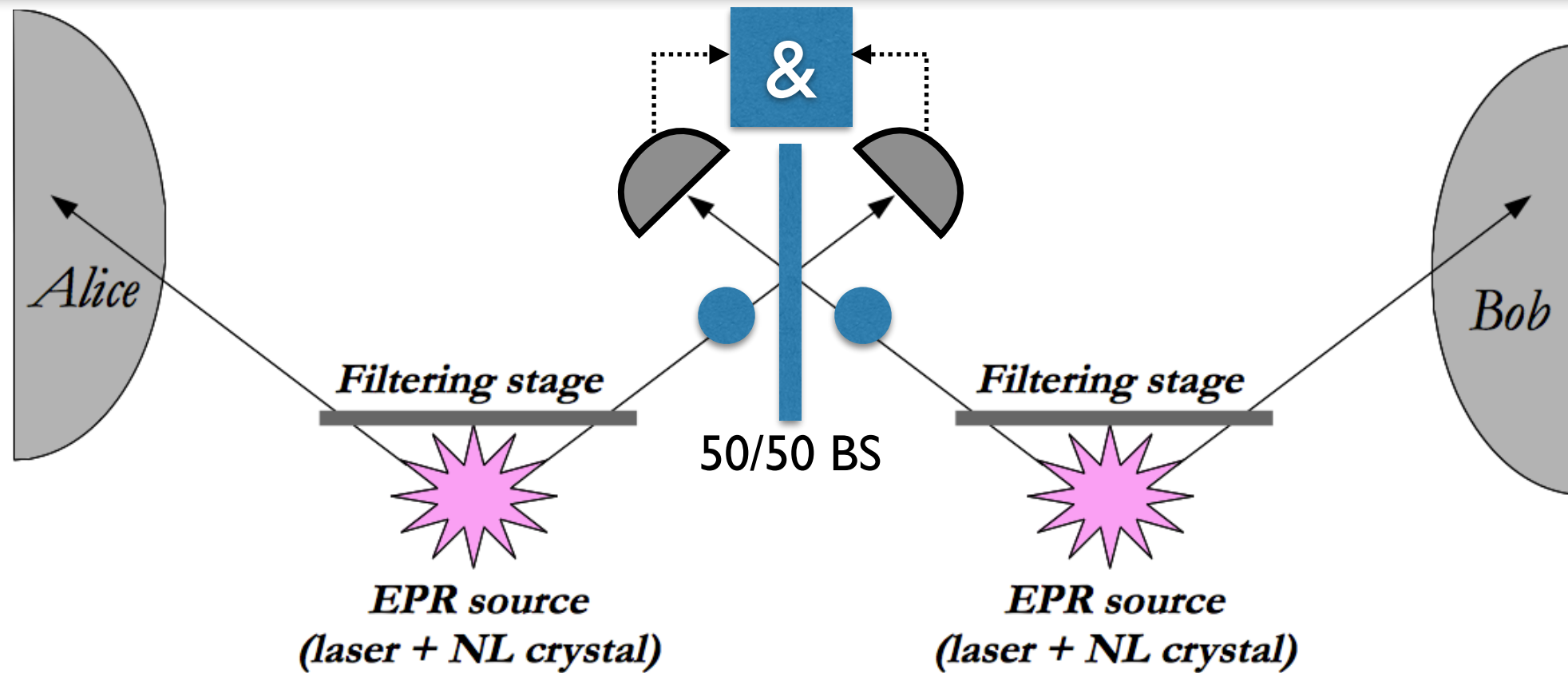
2.4 Connecting HSPs: indistinguishable single photons



2.4 Connecting HSPs: indistinguishable single photons



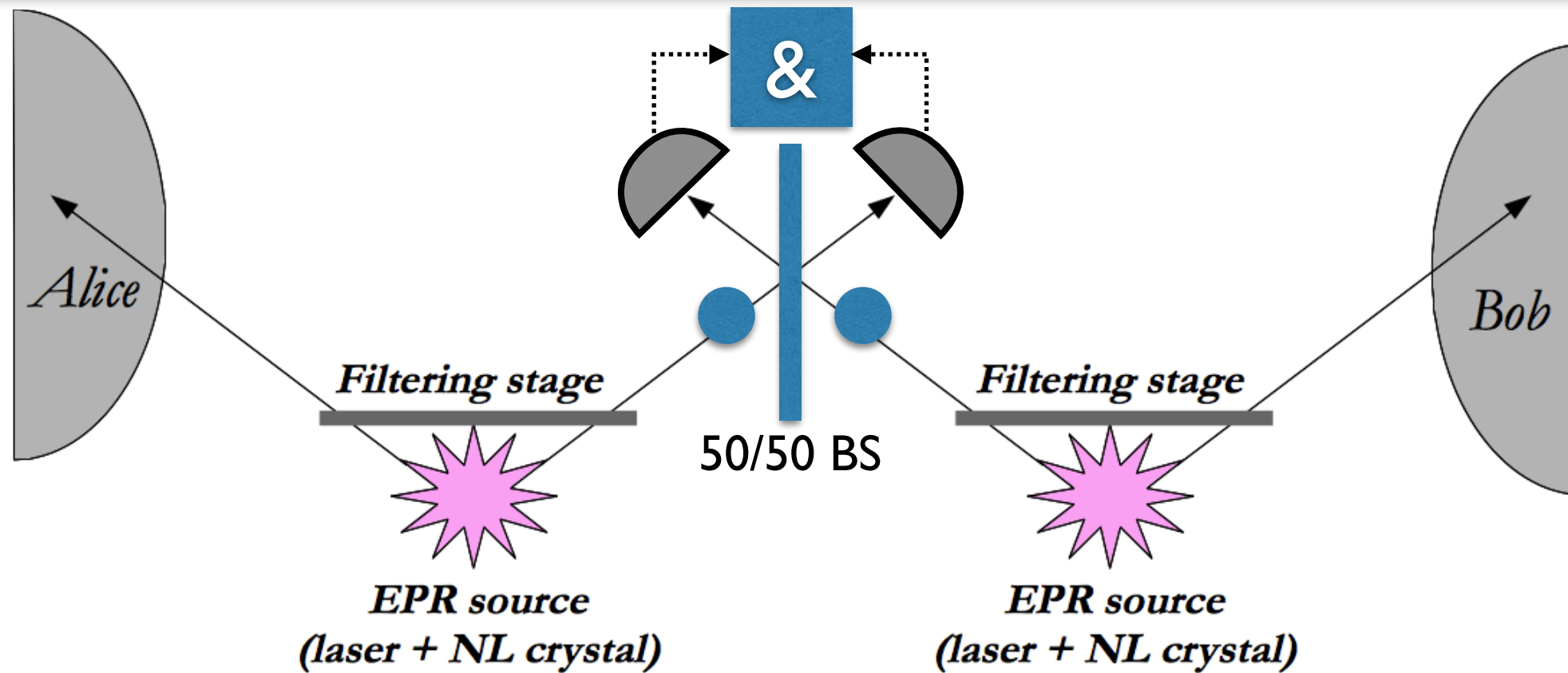
2.4 Connecting HSPs: indistinguishable single photons



Indistinguishable photons should have identical

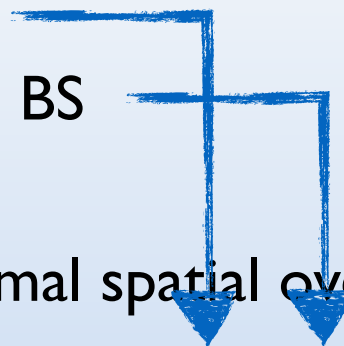
- Wavelengths (λ)
- Spectra ($\Delta\lambda$)
- Arrival times at the BS
- Polarization modes
- Spatial mode (maximal spatial overlap at the BS)

2.4 Connecting HSPs: indistinguishable single photons

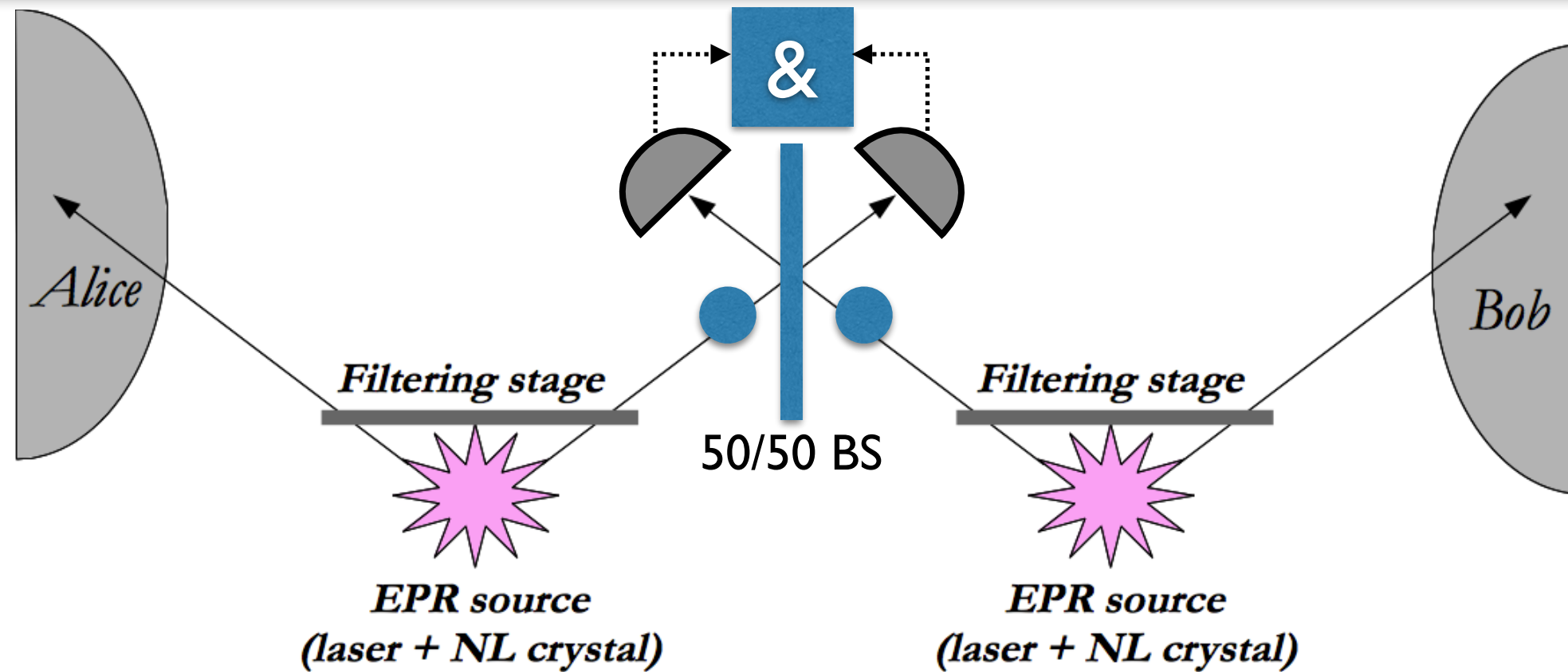


Indistinguishable photons should have identical

- Wavelengths (λ)
- Spectra ($\Delta\lambda$)
- Arrival times at the BS
- Polarization modes
- Spatial mode (maximal spatial overlap at the BS)
- Both in a single temporal mode



2.4 Connecting HSPs: indistinguishable single photons



Indistinguishable photons should have identical

- Wavelengths (λ)
- Spectra ($\Delta\lambda$)
- Arrival times at the BS
- Polarization modes
- Spatial mode (maximal spatial overlap at the BS)
- Both in a single temporal mode

Achieved by

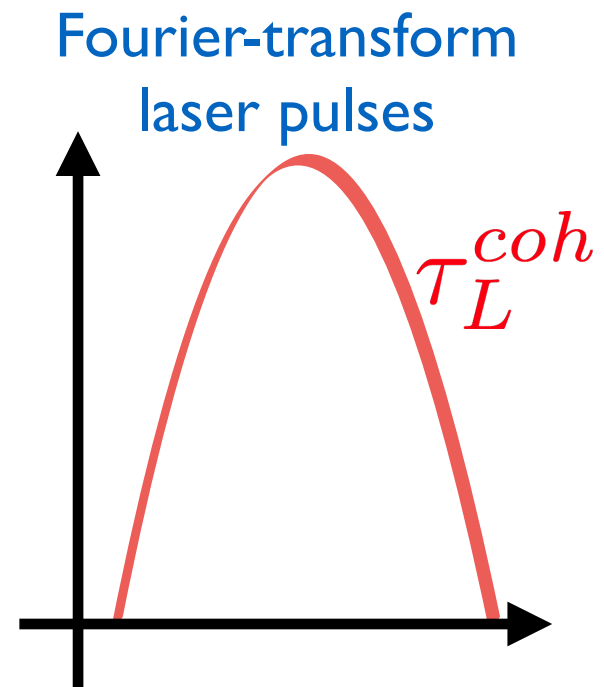
- Phase matching in the Xtals
- Filtering stages
- Synchronization
- Polarization controllers
- Single mode fibers
- HBT measurement !

2.4 Connecting HSPs: single temporal mode

- ▶ The photons are created in a multimode manner

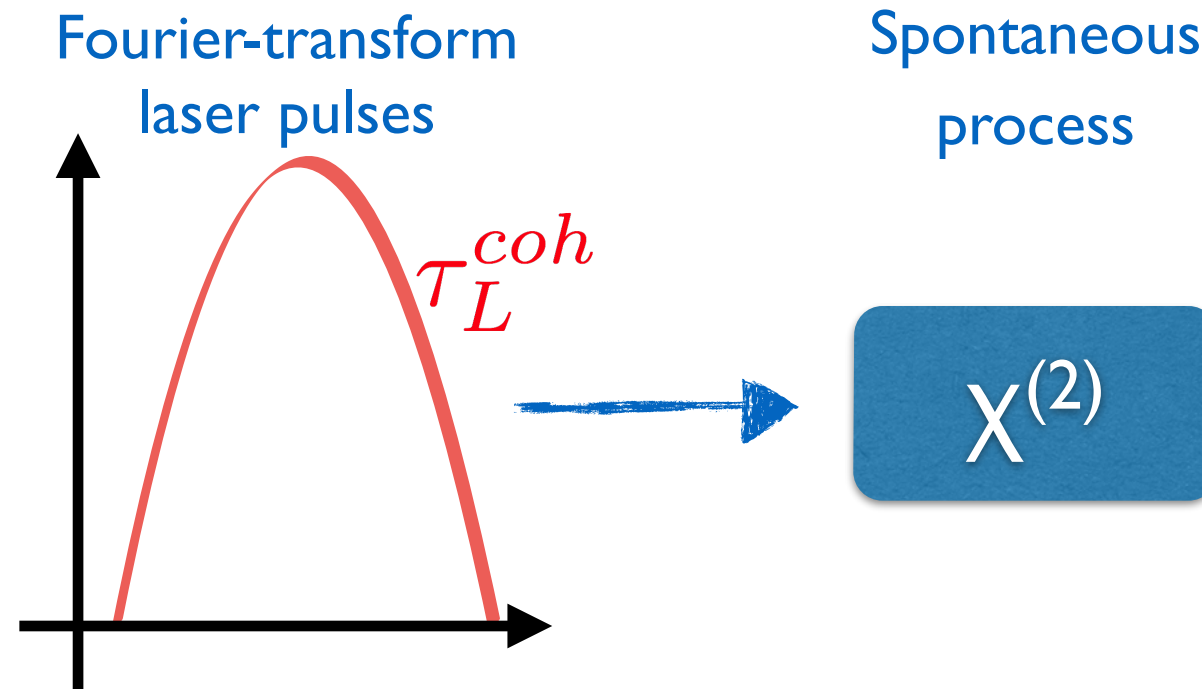
2.4 Connecting HSPs: single temporal mode

- ▶ The photons are created in a multimode manner



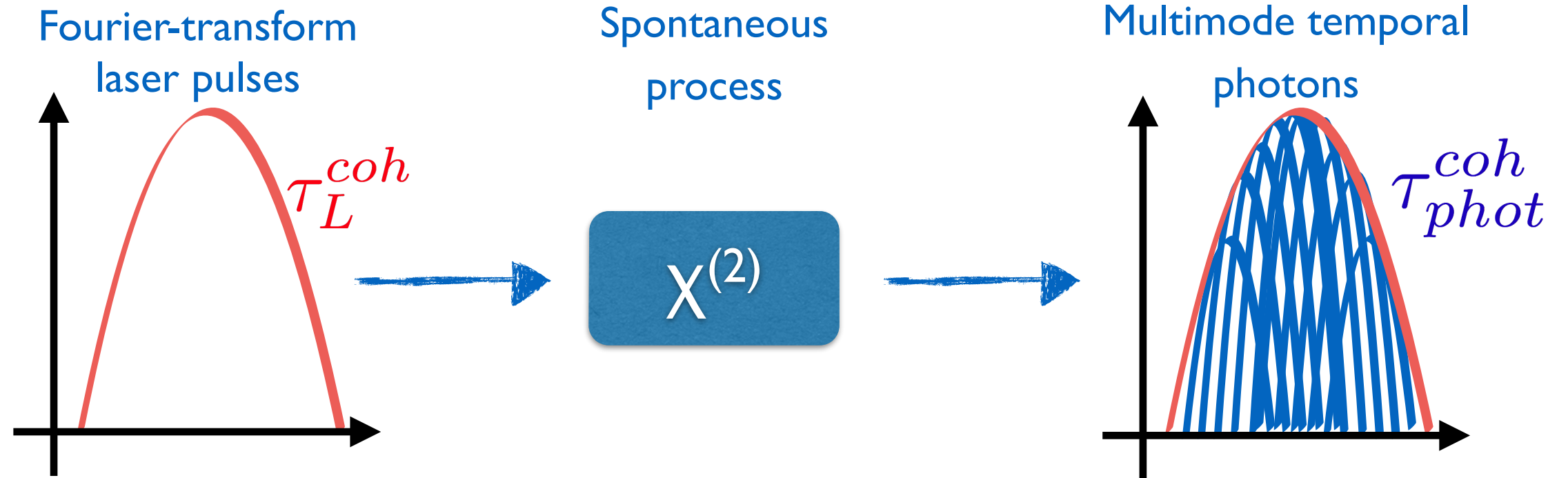
2.4 Connecting HSPs: single temporal mode

- ▶ The photons are created in a multimode manner



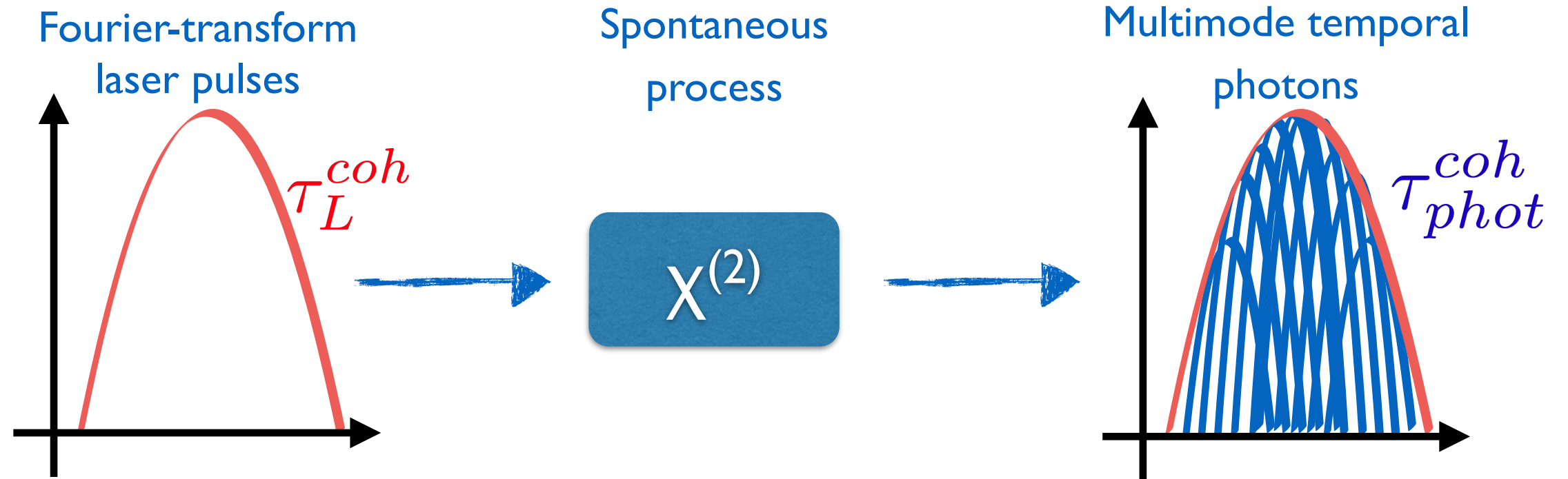
2.4 Connecting HSPs: single temporal mode

- ▶ The photons are created in a multimode manner



2.4 Connecting HSPs: single temporal mode

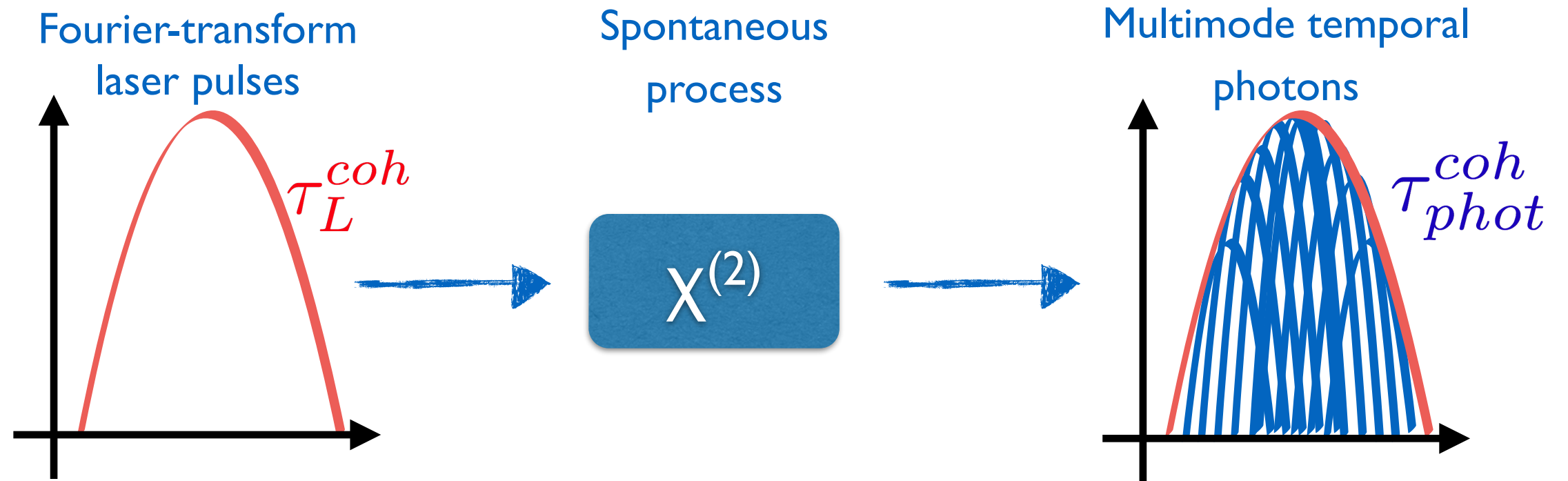
- ▶ The photons are created in a multimode manner



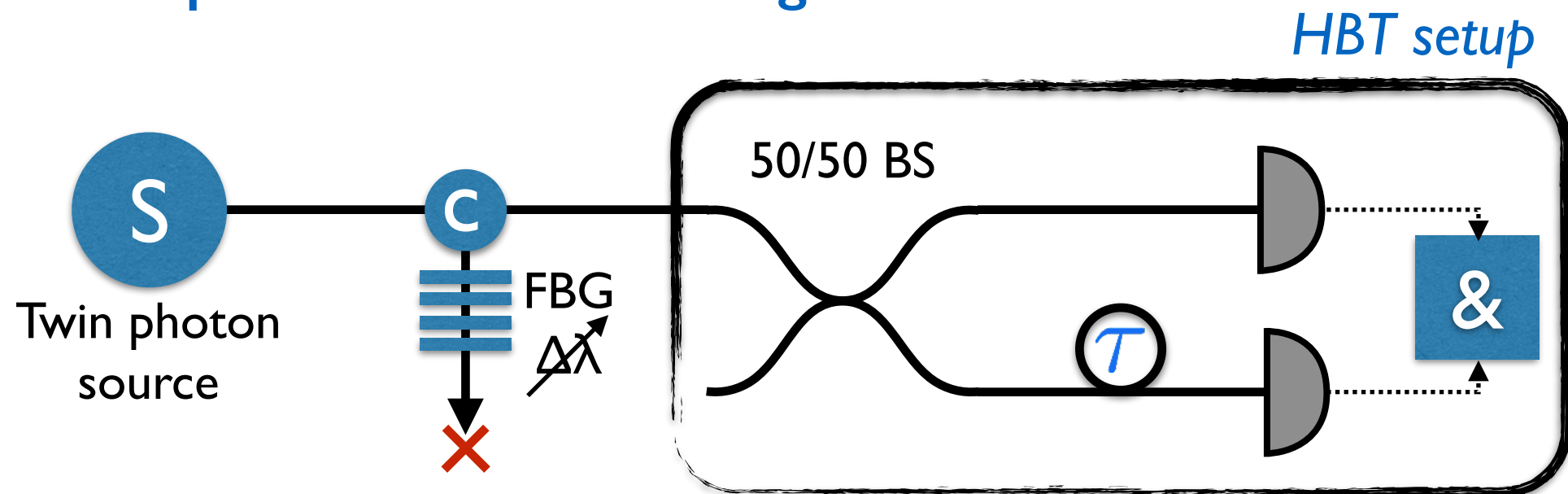
- ▶ HBT setup without the heralding mode

2.4 Connecting HSPs: single temporal mode

- ▶ The photons are created in a multimode manner



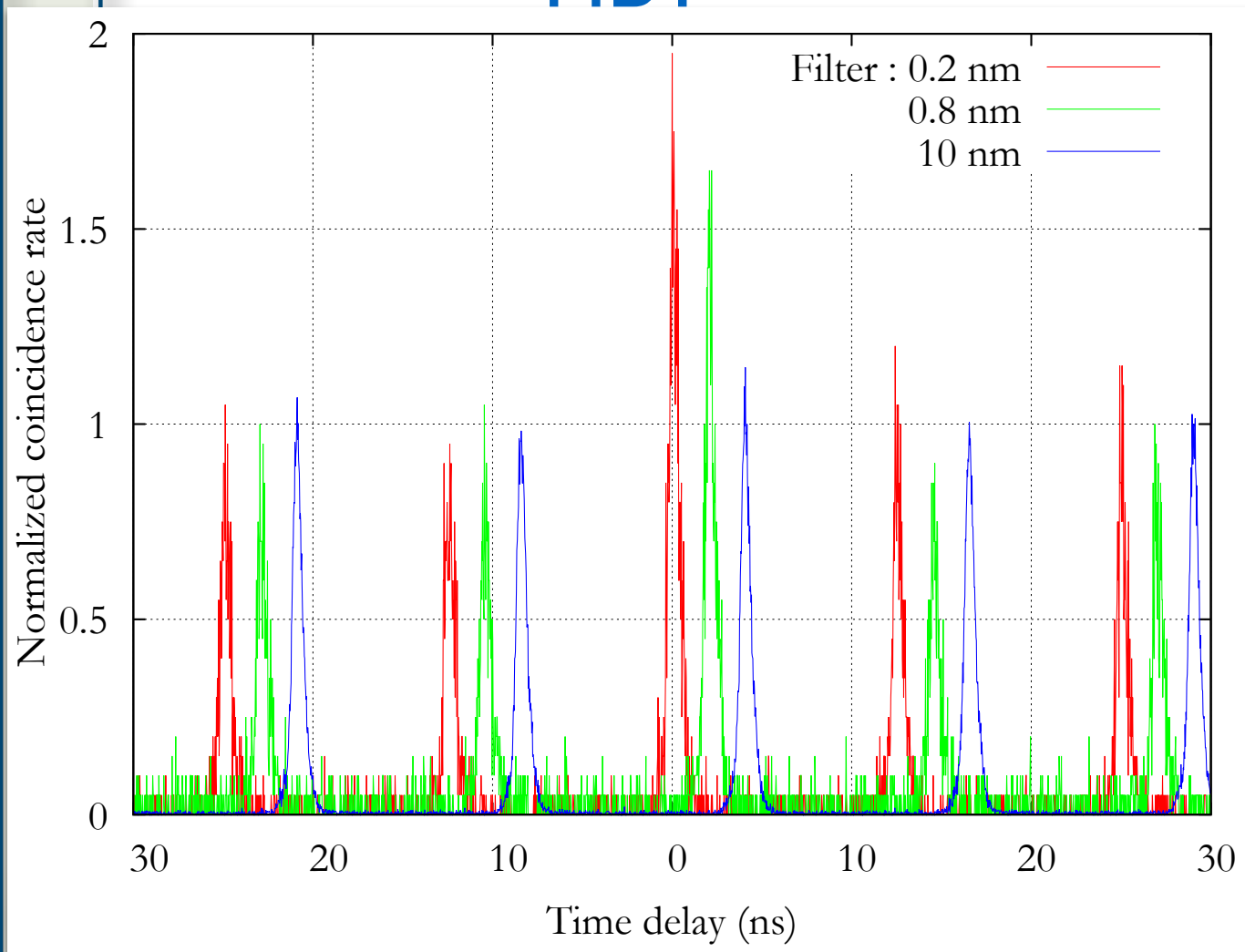
- ▶ HBT setup without the heralding mode



2.4 Connecting HSPs: single temporal mode

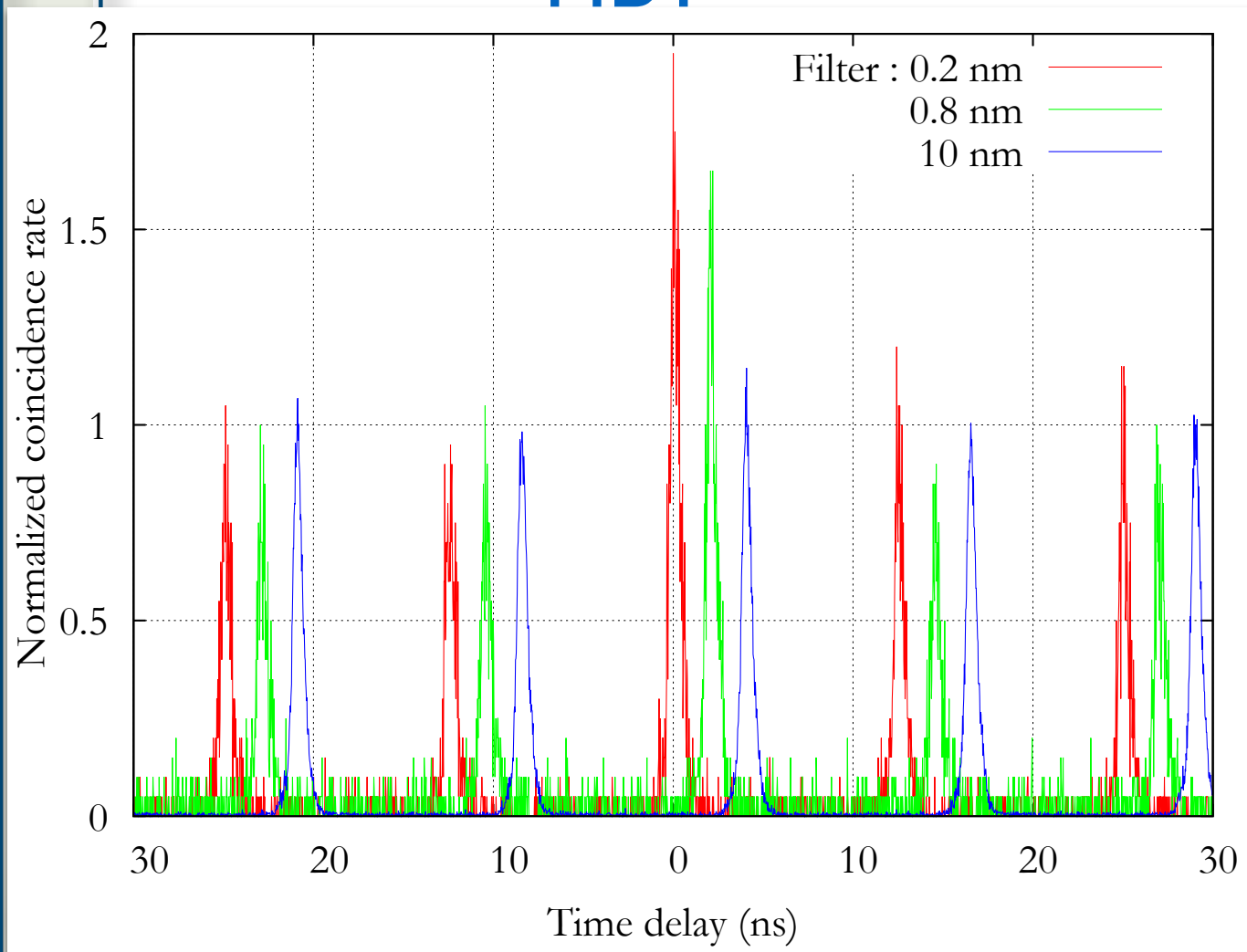
2.4 Connecting HSPs: single temporal mode

HBT



2.4 Connecting HSPs: single temporal mode

HBT

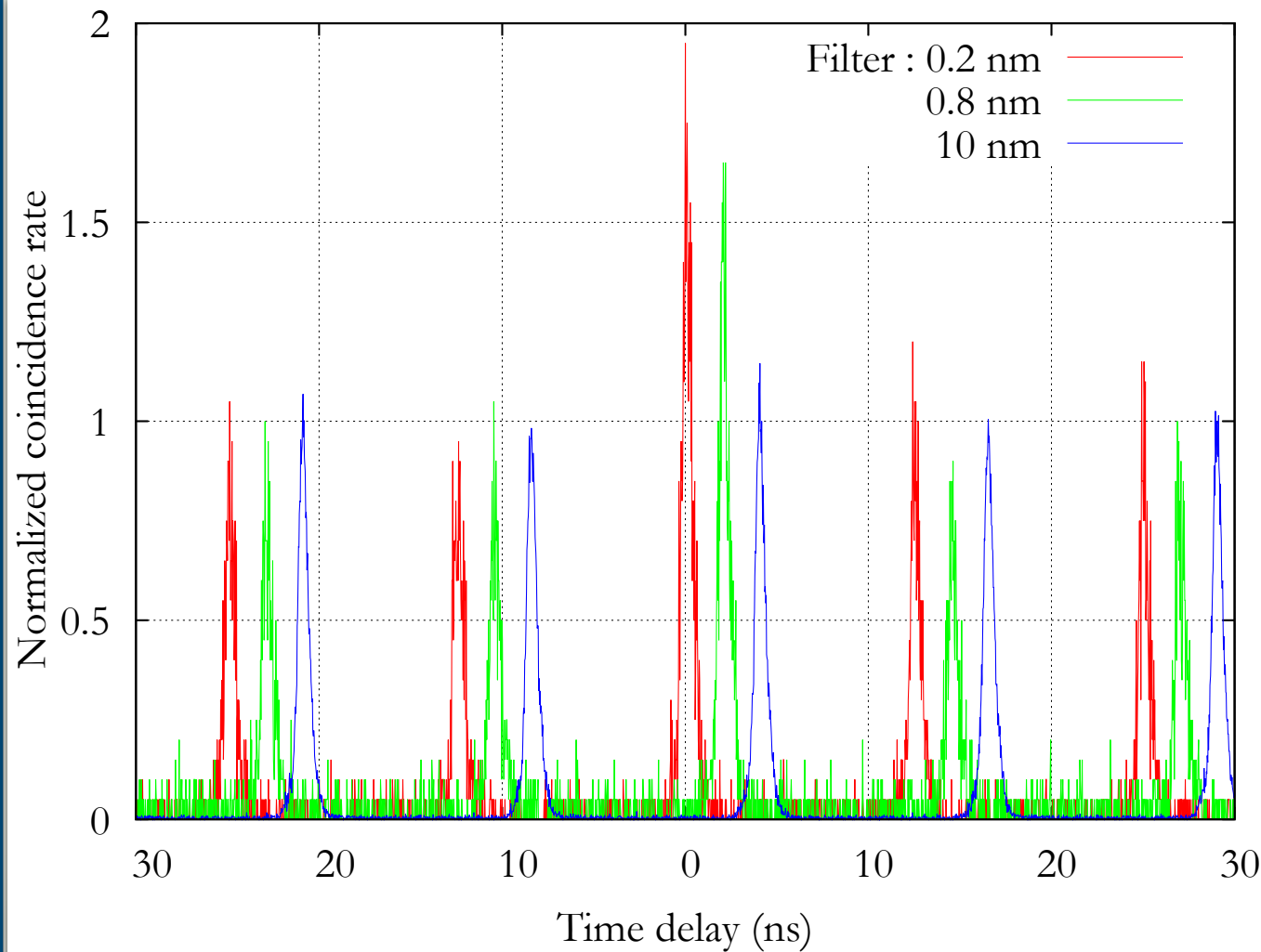


Thermal stat. → bunching effect

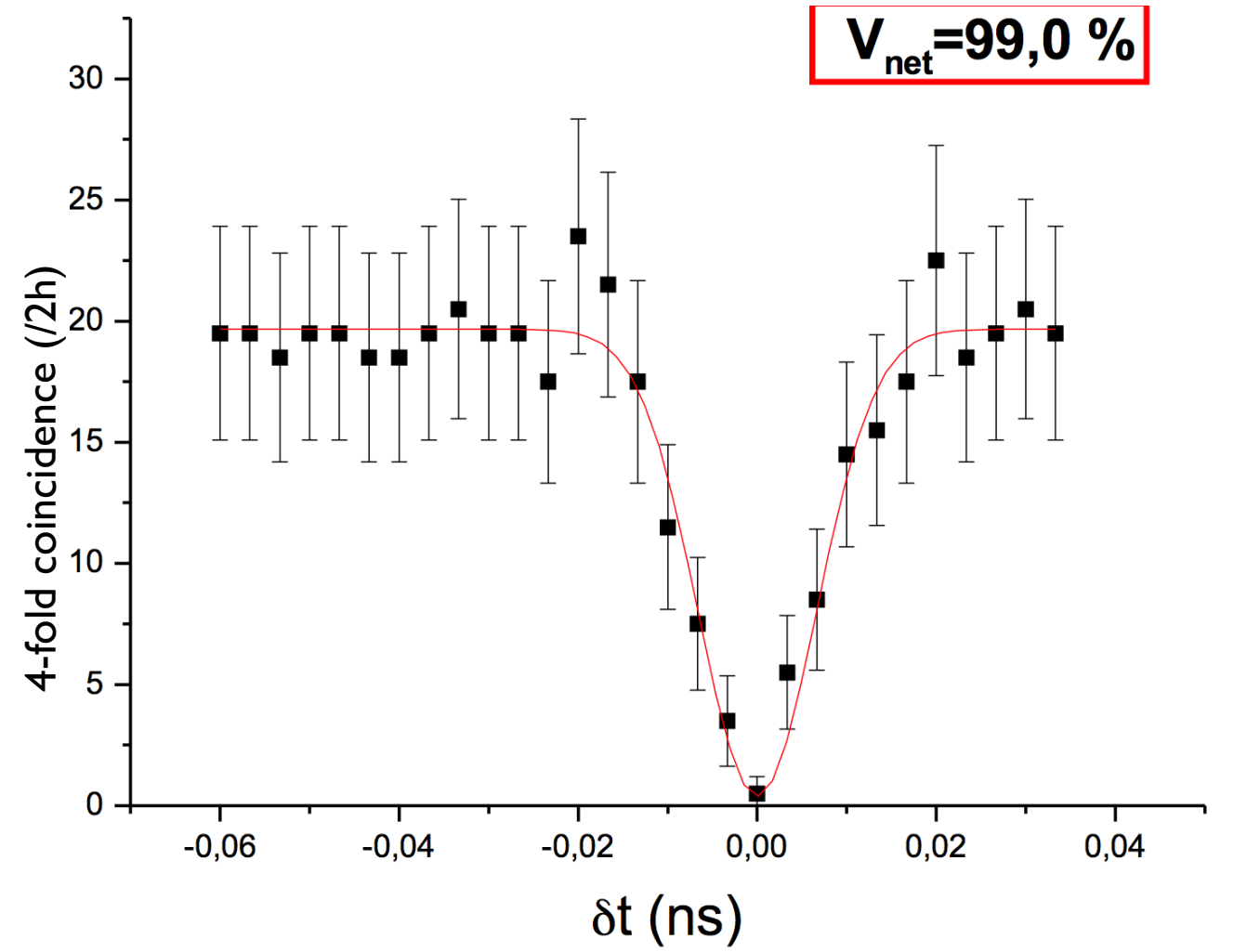
Poissonian stat. → coherent state

2.4 Connecting HSPs: single temporal mode

HBT



HOM

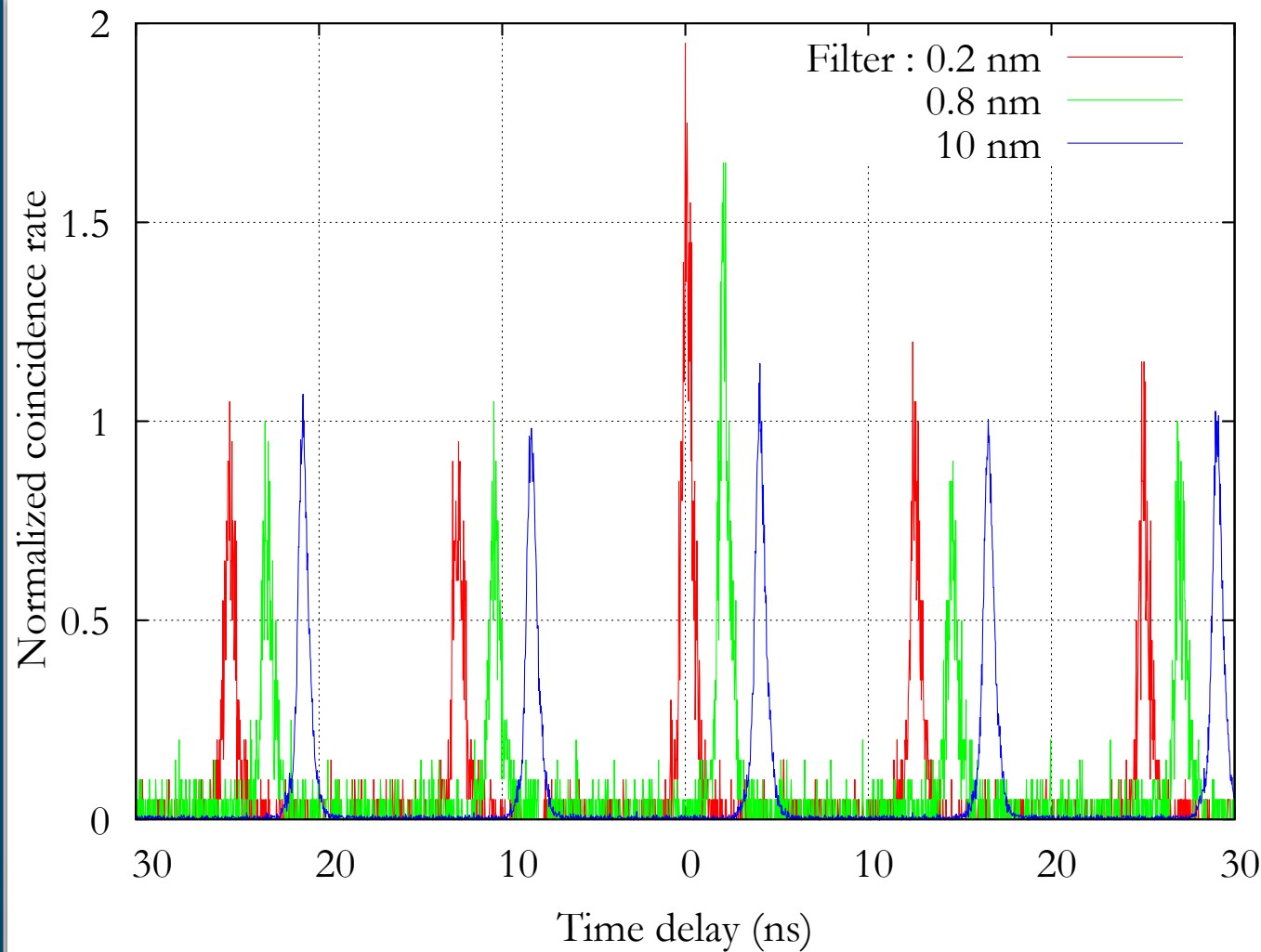


Thermal stat. → bunching effect

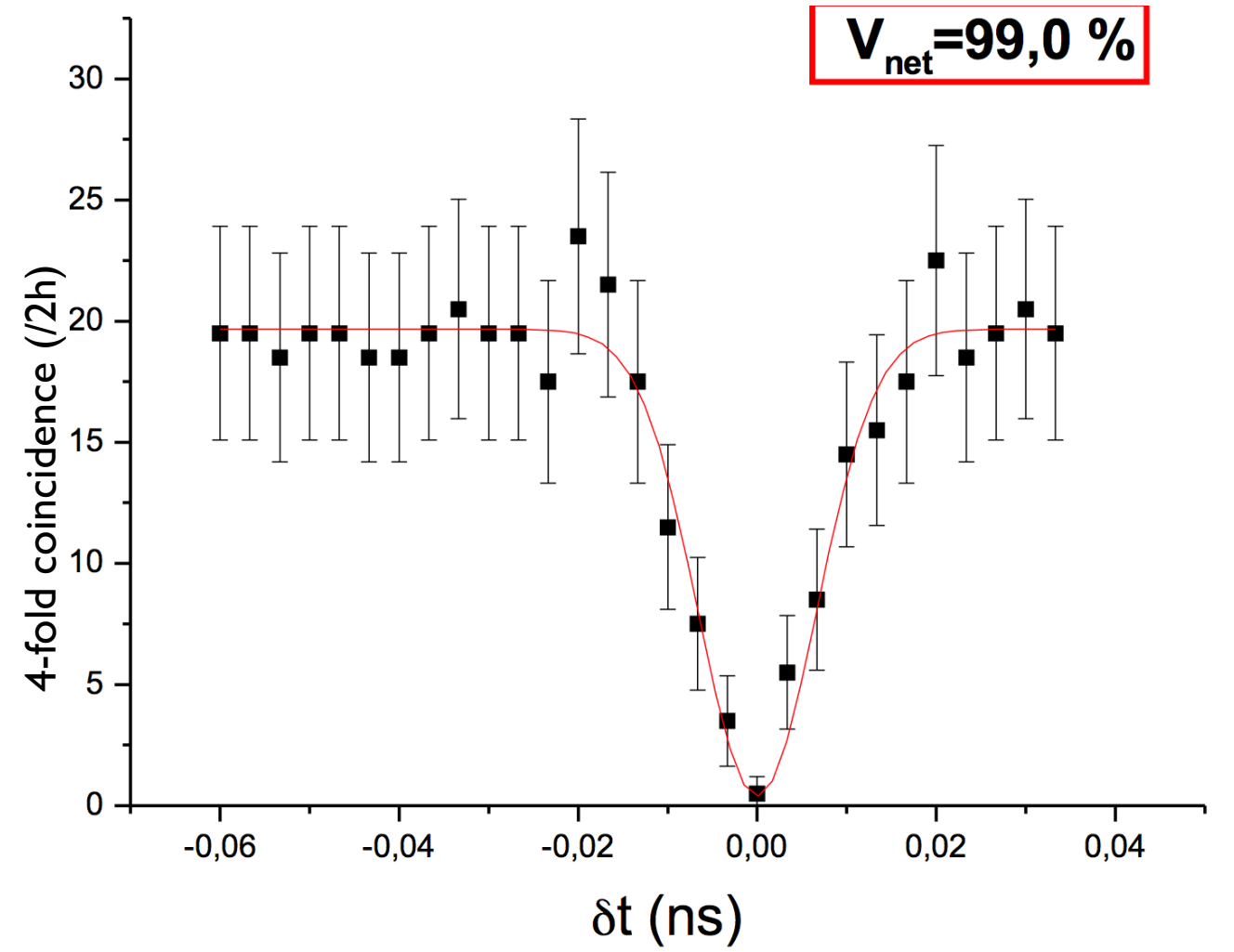
Poissonian stat. → coherent state

2.4 Connecting HSPs: single temporal mode

HBT



HOM



Thermal stat. \rightarrow bunching effect

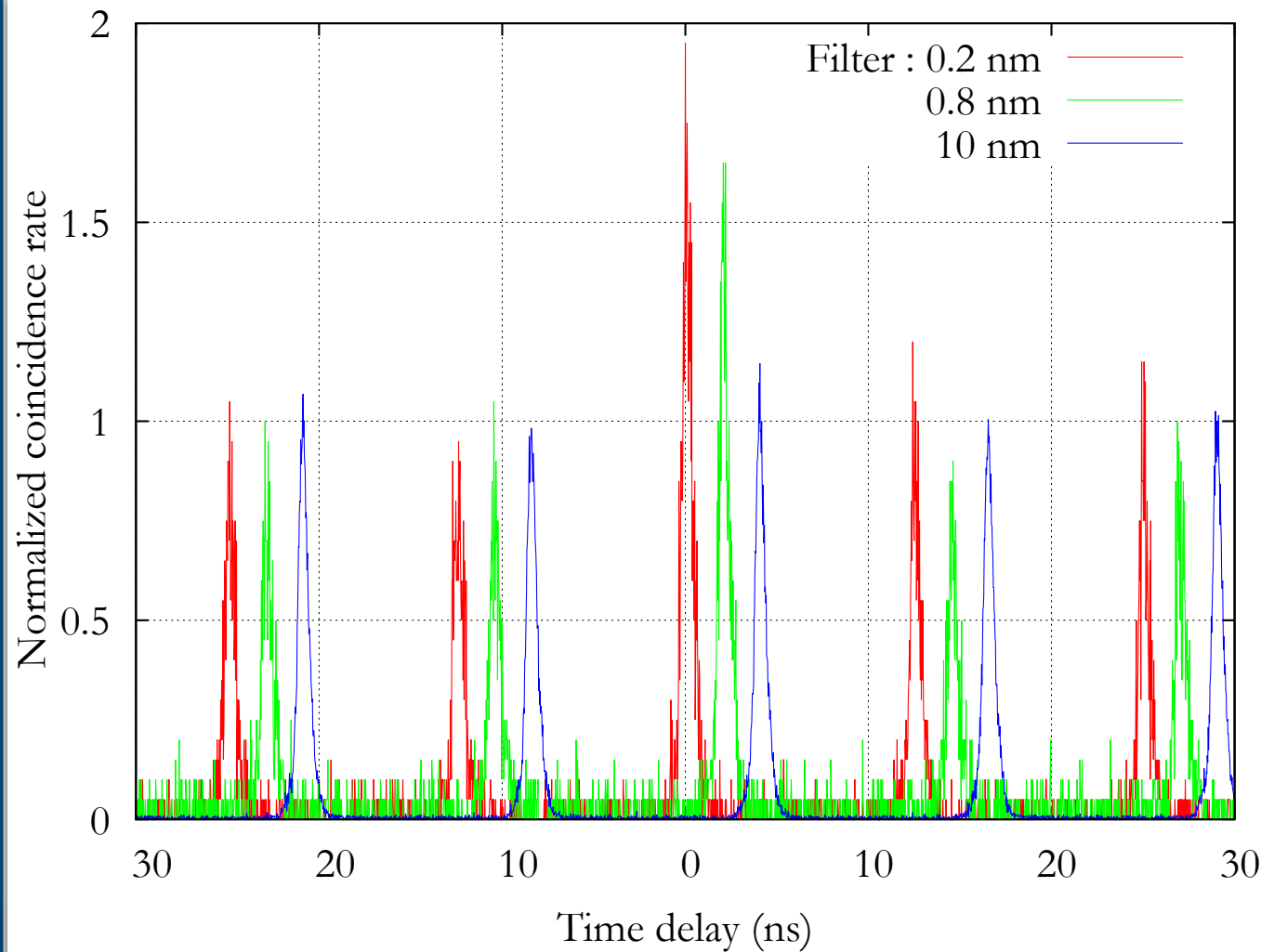
Poissonian stat. \rightarrow coherent state

Max HOM interf. VIS (pure photons)

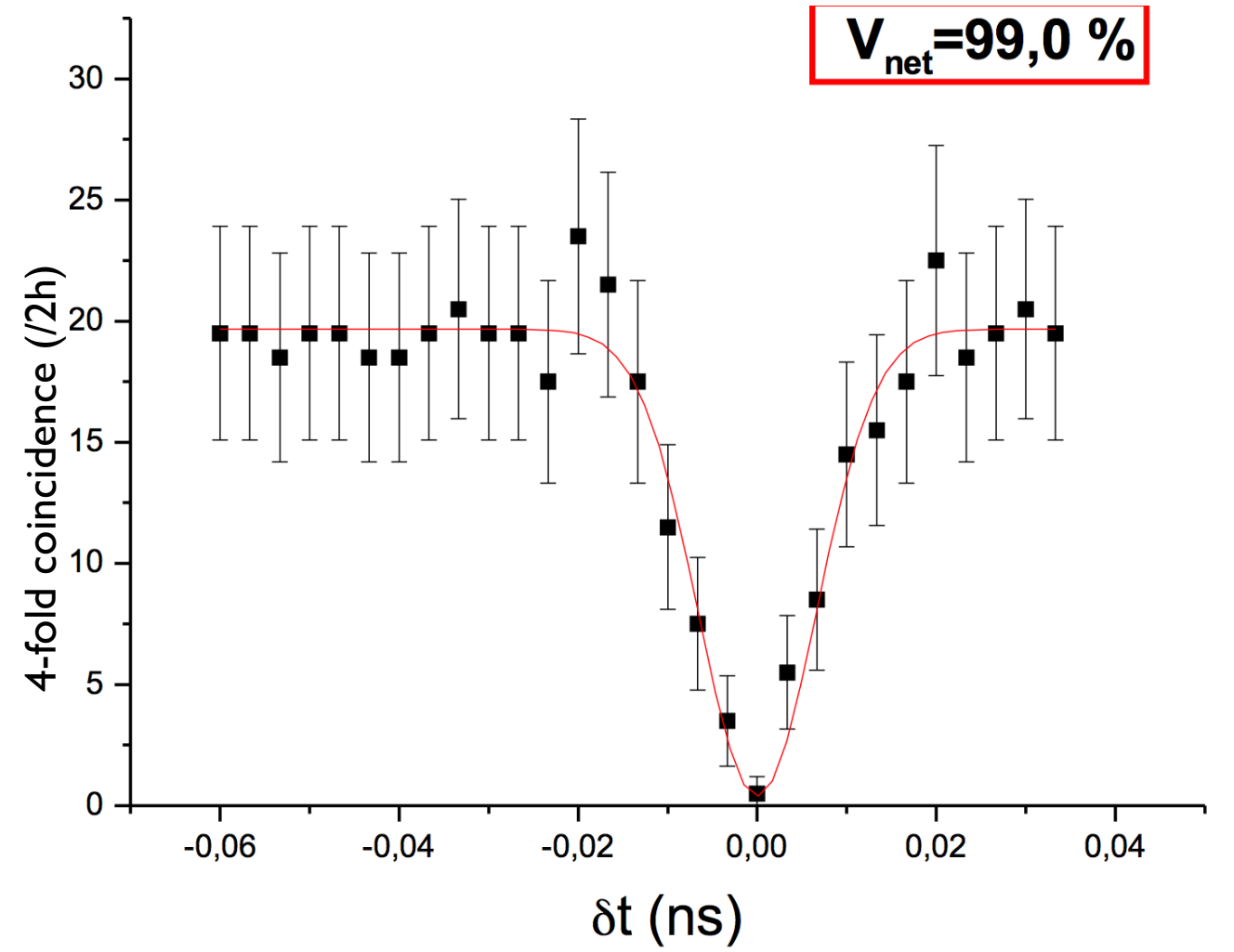
Visibility limited to 33%

2.4 Connecting HSPs: single temporal mode

HBT



HOM



Thermal stat. → bunching effect

Poissonian stat. → coherent state

Max HOM interf. VIS (pure photons)

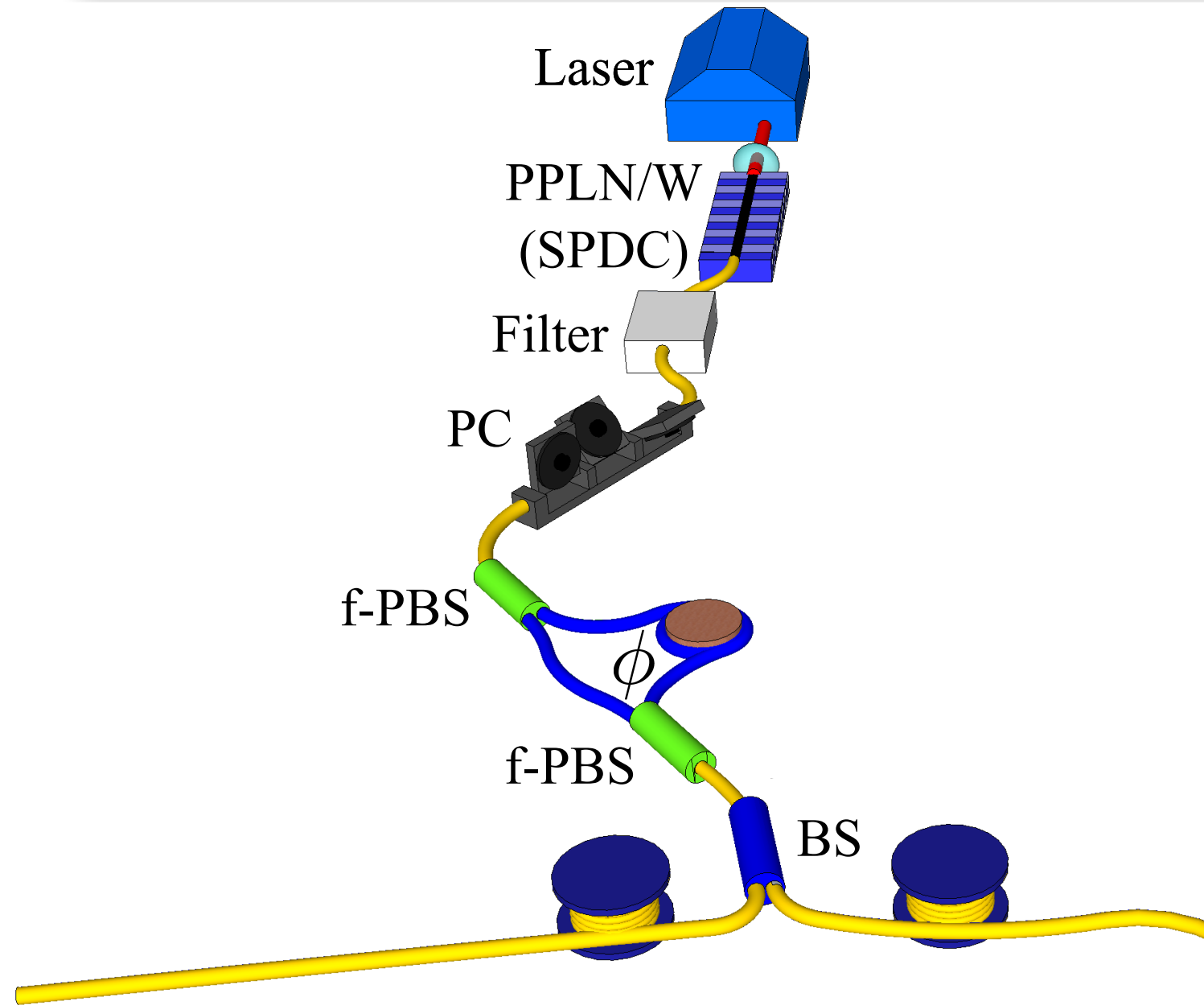
Visibility limited to 33%



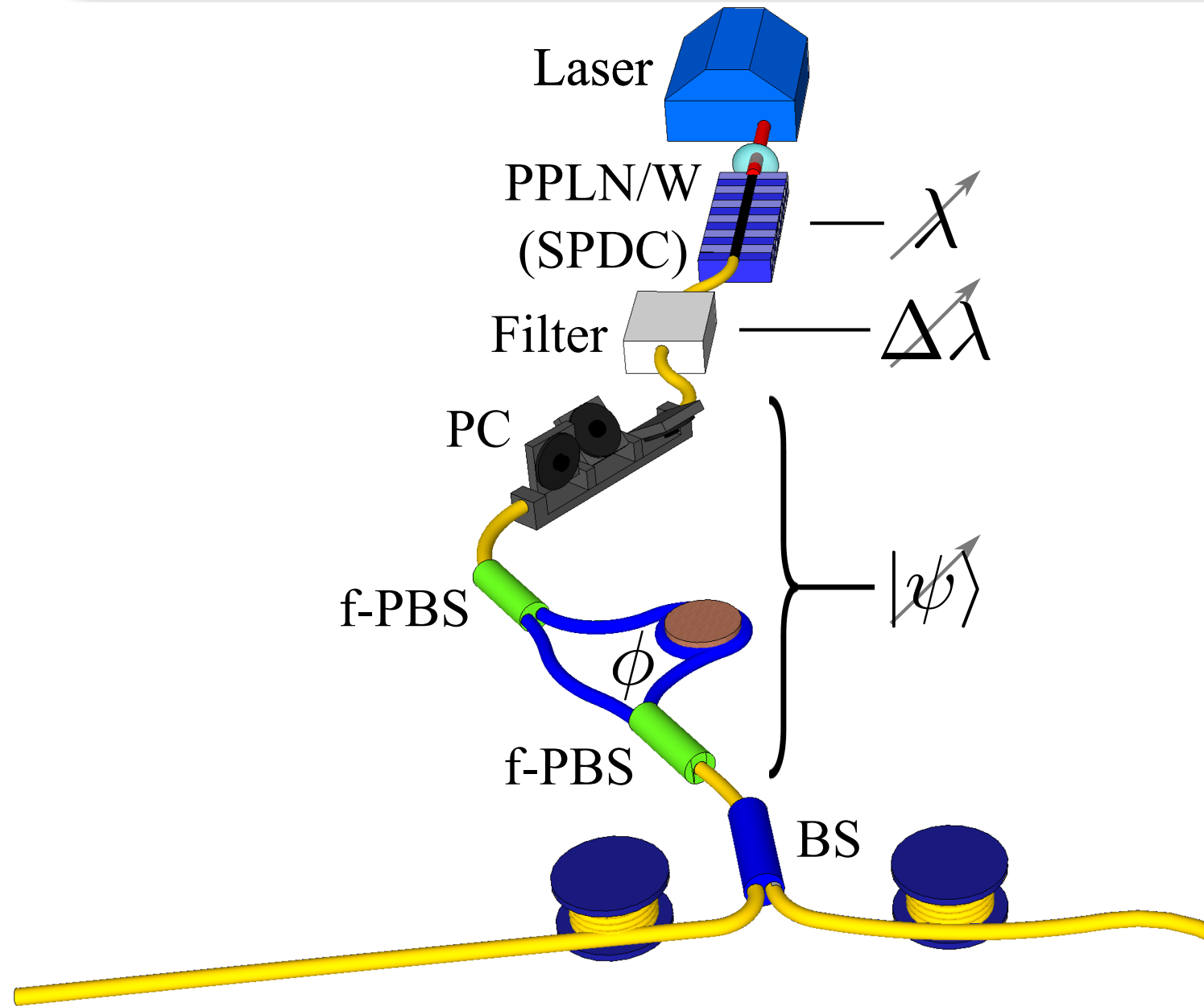
Outline

1. Introduction
2. HBT for characterizing single photon sources (SPS)
3. HBT for characterizing photon pair sources
 1. A versatile source of polarization entangled photons
 2. Cross-correlation function & coherence time of the emitted photons
 3. Bell inequality type measurements

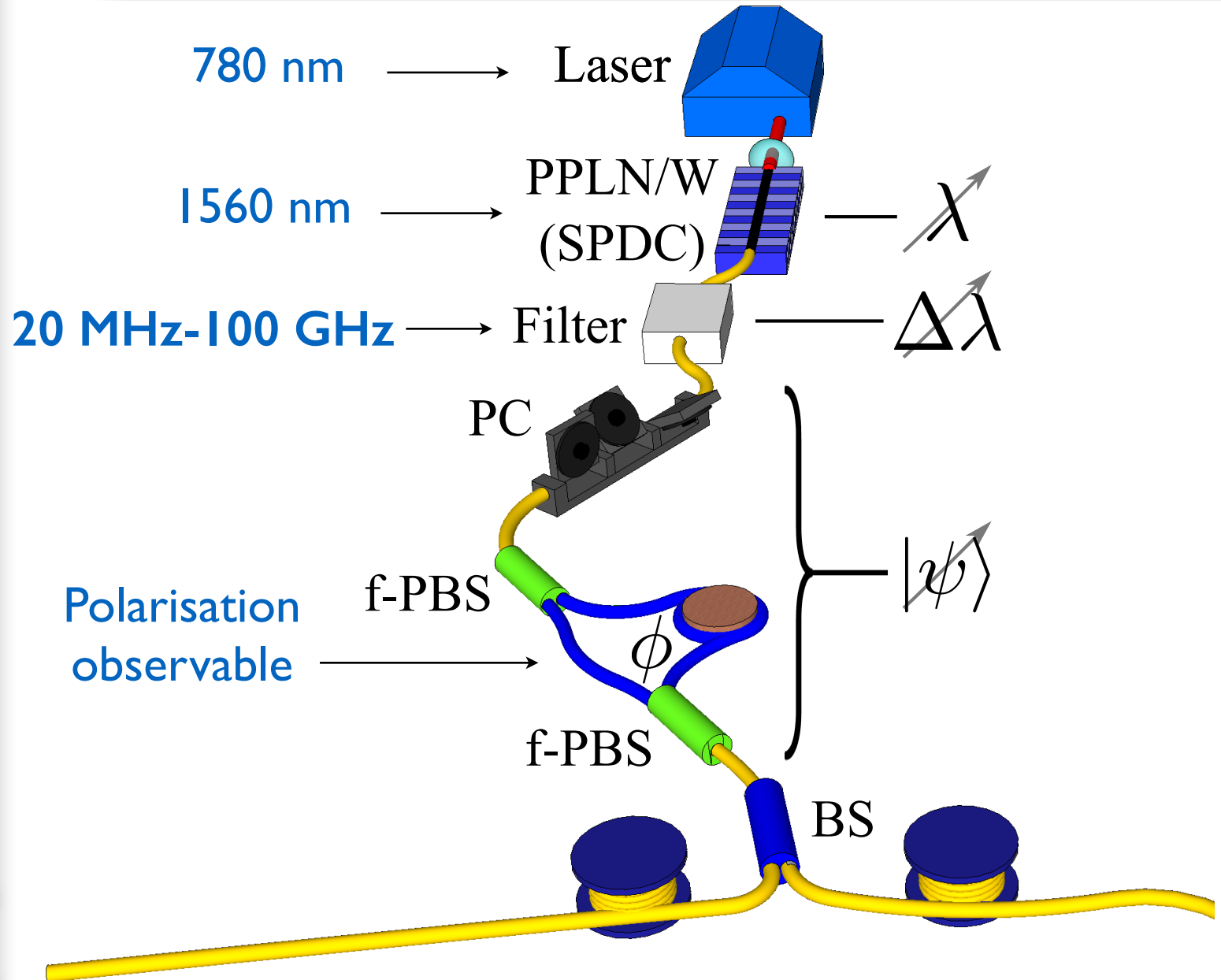
3.1 A versatile polarization entangled photon pair source



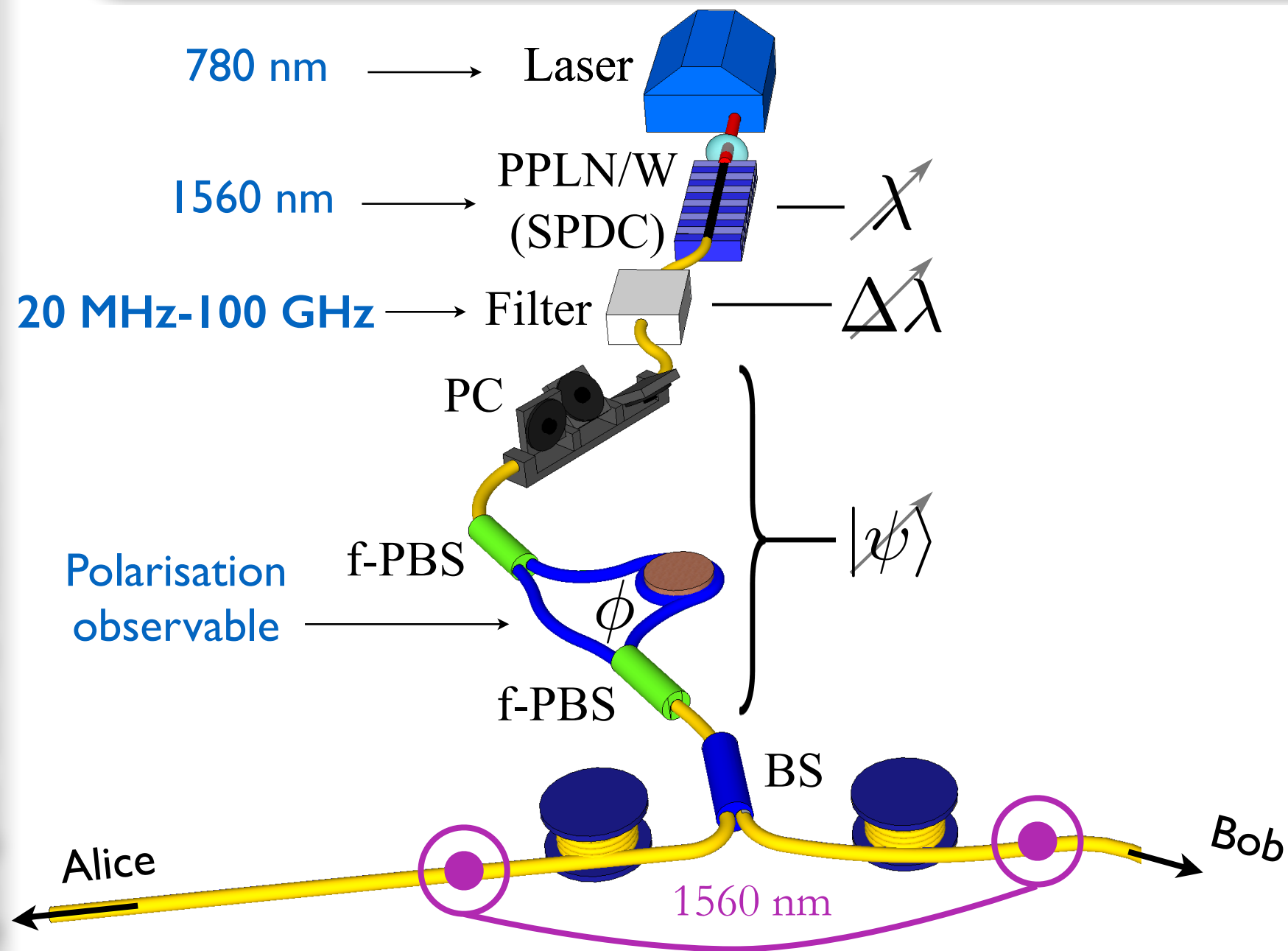
3.1 A versatile polarization entangled photon pair source



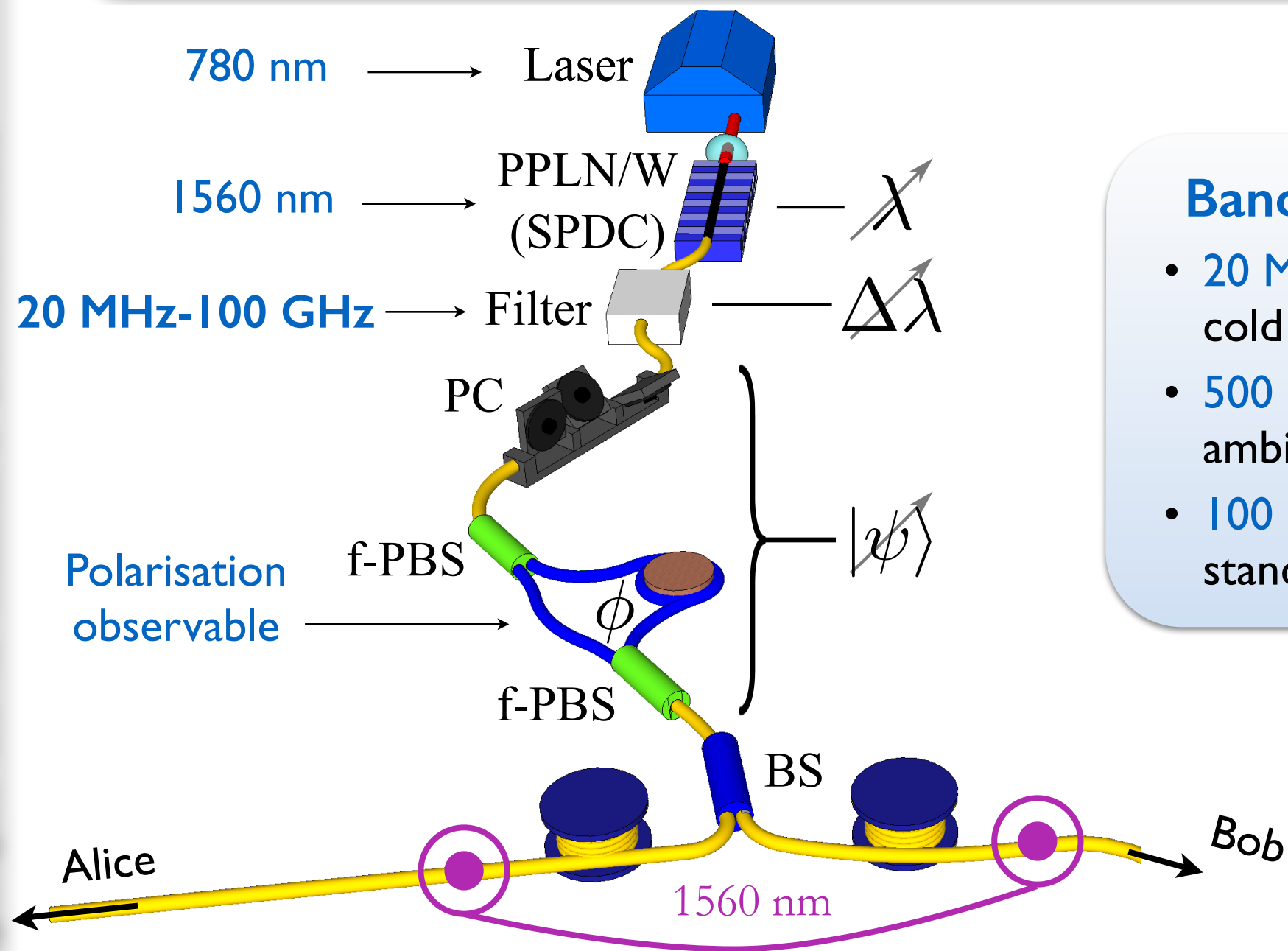
3.1 A versatile polarization entangled photon pair source



3.1 A versatile polarization entangled photon pair source



3.1 A versatile polarization entangled photon pair source

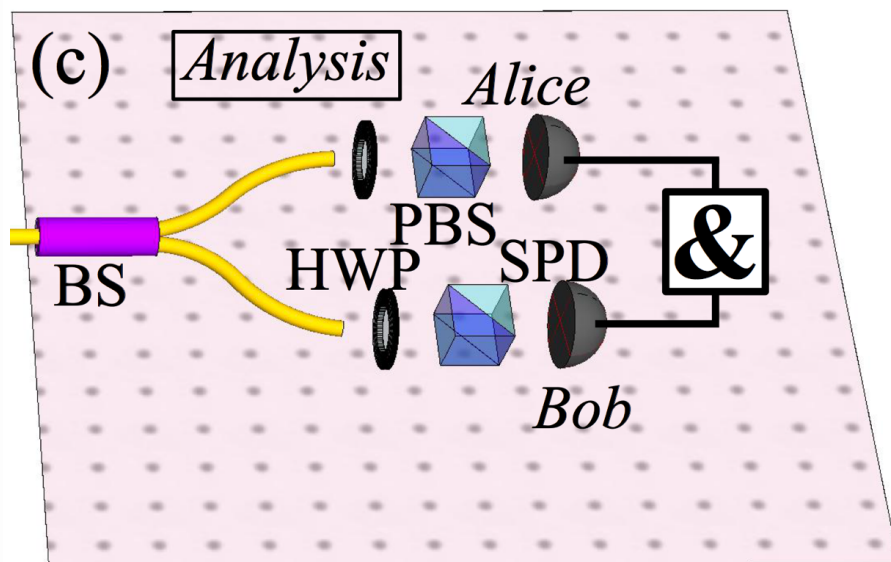


- Bandwidth versatility**
- 20 MHz → Interface with cold atoms Q-memory
 - 500 MHz → Interface with ambient T atomic vapors
 - 100 GHz → Q-crypto in standard telecom channels

F. Kaiser *et al.*, Science [2012]
 F. Kaiser *et al.*, Laser Phys. Lett. [2013]
 F. Kaiser *et al.*, Optics Comm. [2014]

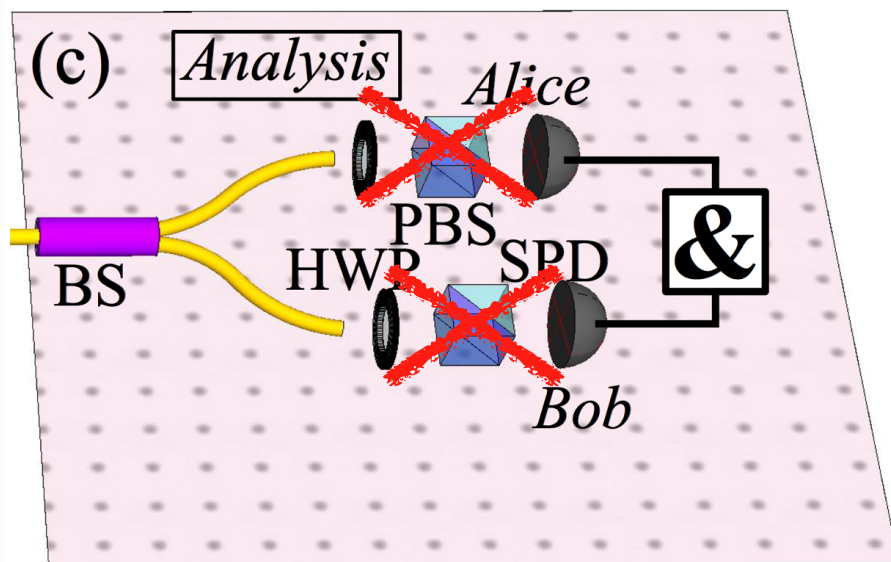
3.2 Cross-correlation function via coincidence measurement

- ▶ Tests with 3 FBG filters : 80 GHz, 540 MHz, 25 MHz



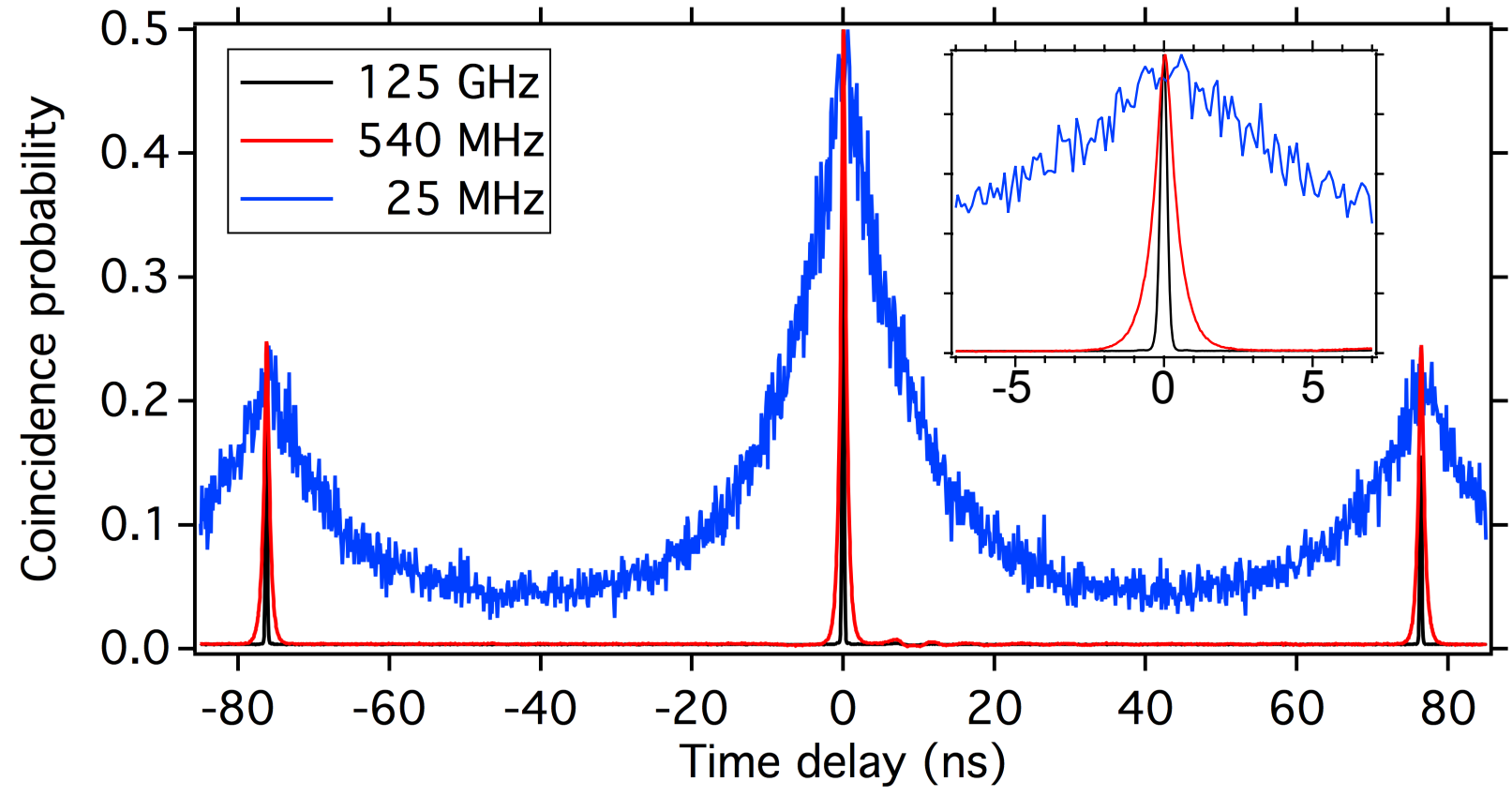
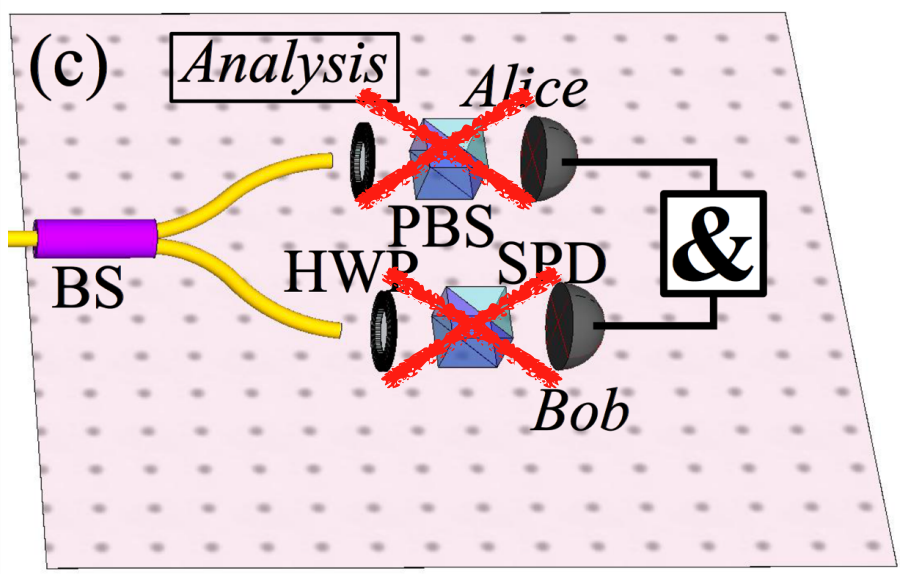
3.2 Cross-correlation function via coincidence measurement

- ▶ Tests with 3 FBG filters : 80 GHz, 540 MHz, 25 MHz



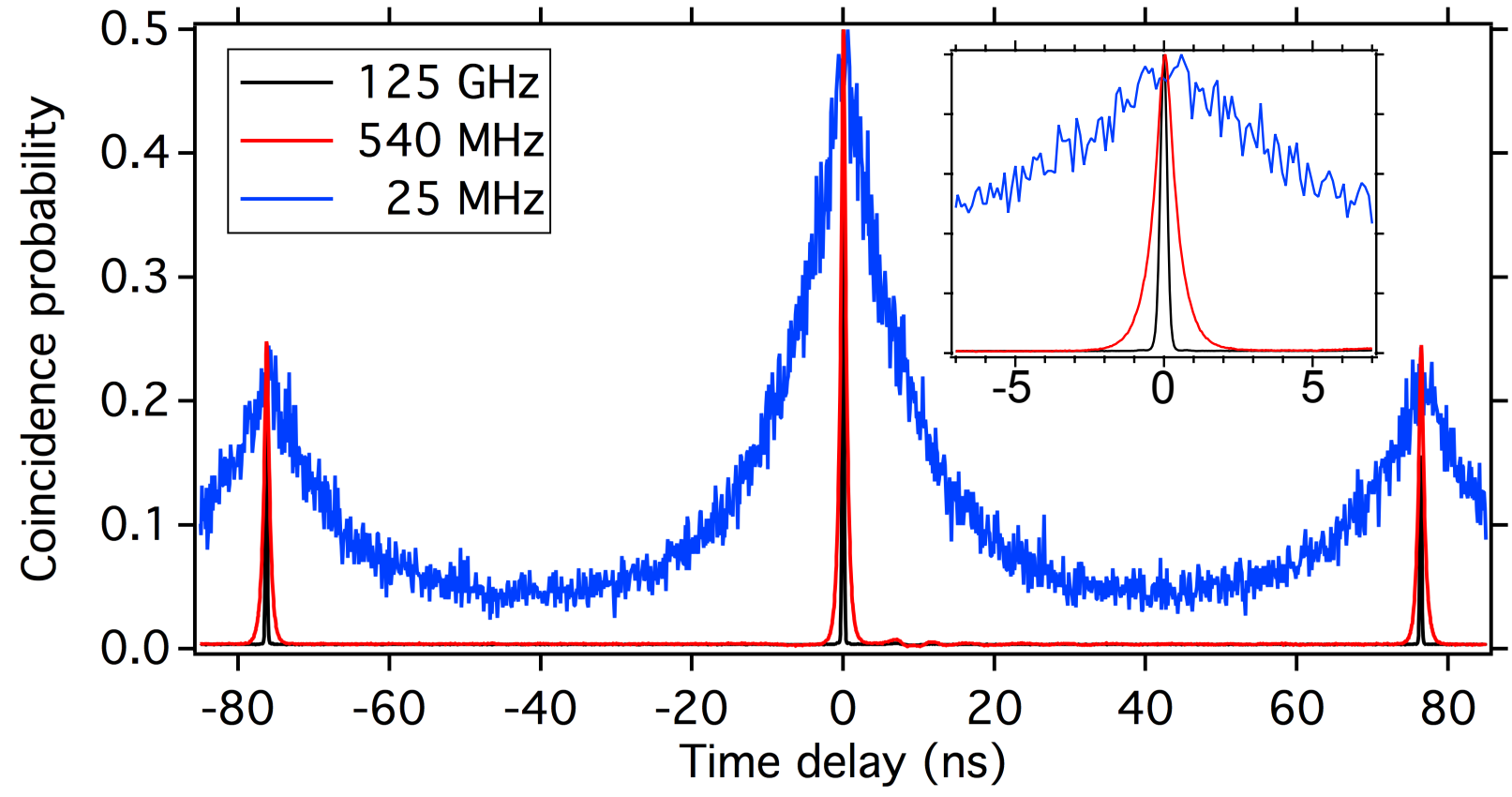
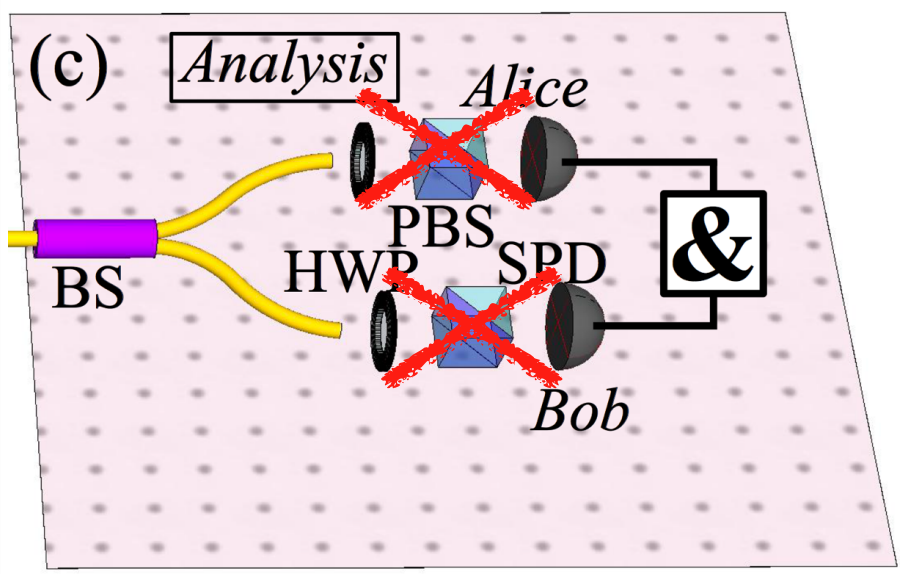
3.2 Cross-correlation function via coincidence mst

► Tests with 3 FBG filters : 80 GHz, 540 MHz, 25 MHz



3.2 Cross-correlation function via coincidence mst

▶ Tests with 3 FBG filters : 80 GHz, 540 MHz, 25 MHz

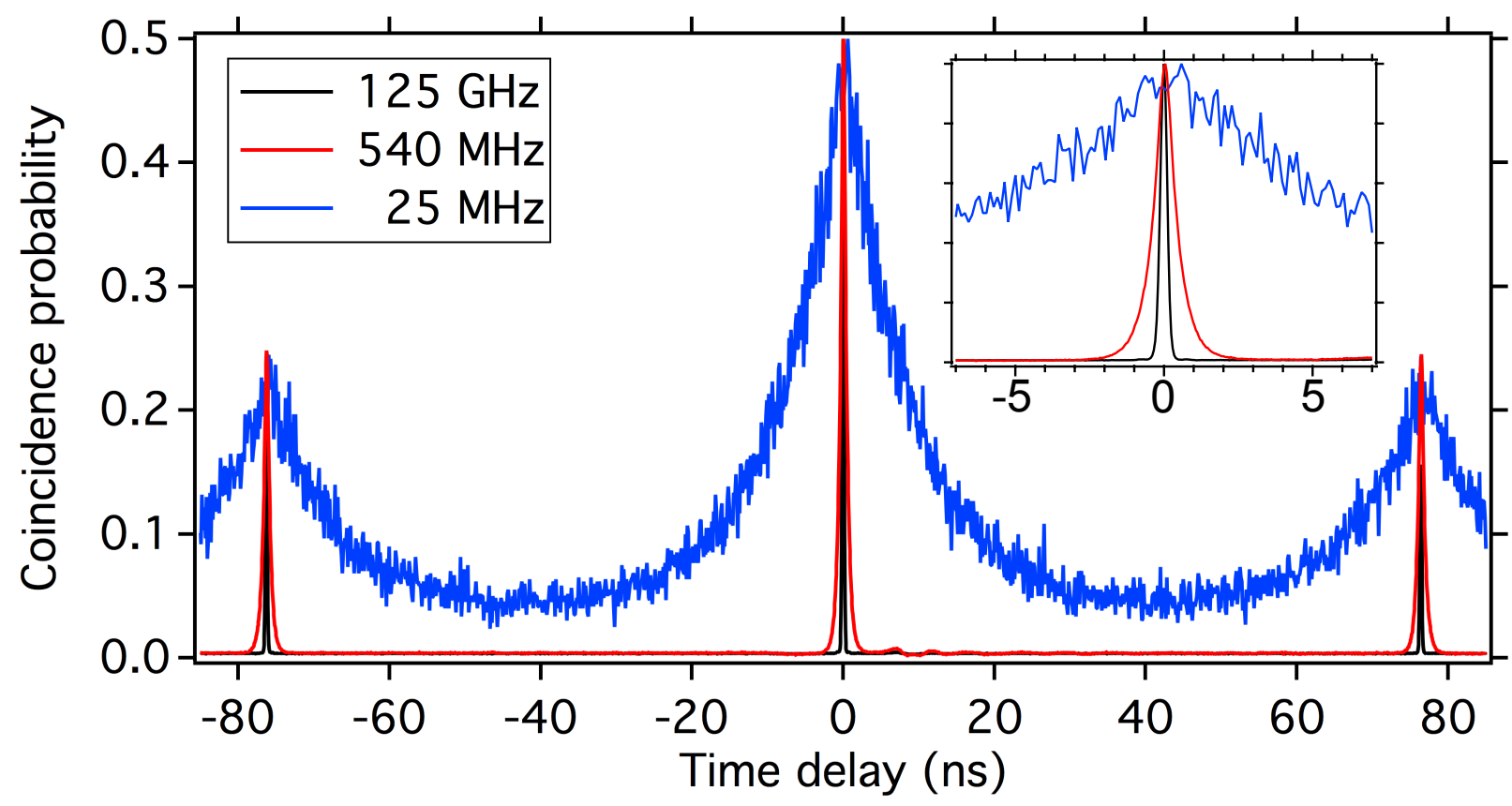
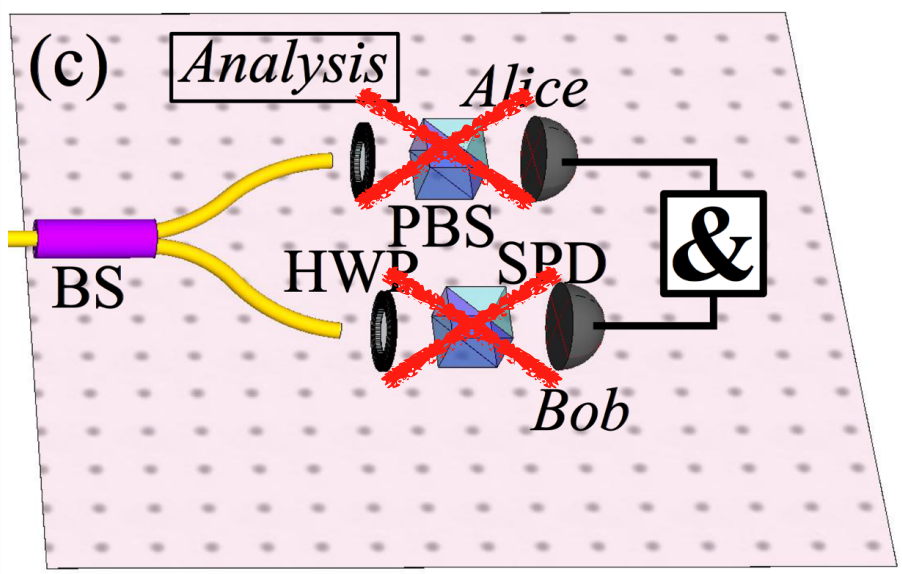


	τ_{phot}	$\Delta\nu$
—	440 ps	125 GHz
—	0.8 ns	540 MHz
—	15.6 ns	25 MHz

↔

3.2 Cross-correlation function via coincidence mst

▶ Tests with 3 FBG filters : 80 GHz, 540 MHz, 25 MHz



	τ_{phot}	$\Delta\nu$
—	440 ps	125 GHz
—	0.8 ns	540 MHz
—	15.6 ns	25 MHz

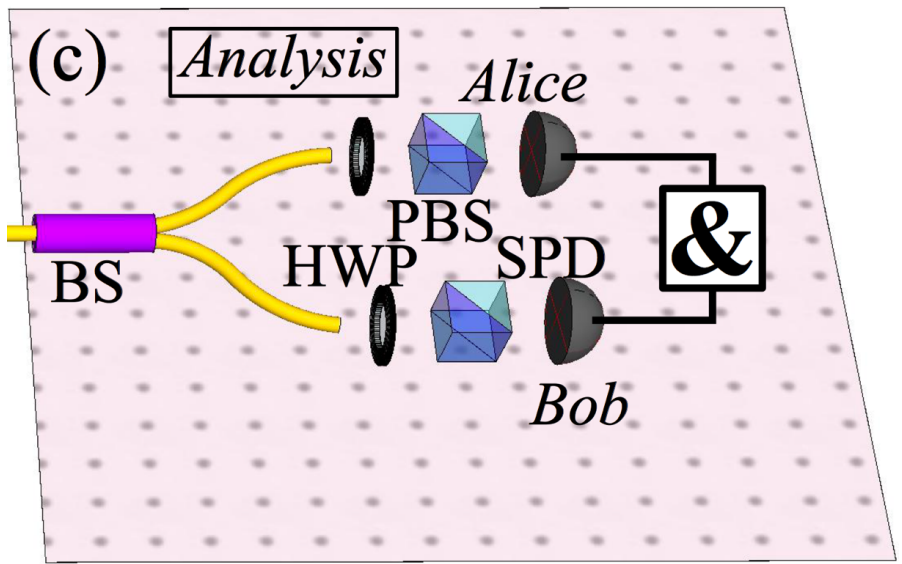
\leftrightarrow

From the CAR

- Clear violation of the Cauchy-Schwartz inequality
- Twin photons are emitted simultaneously, at the quantum level

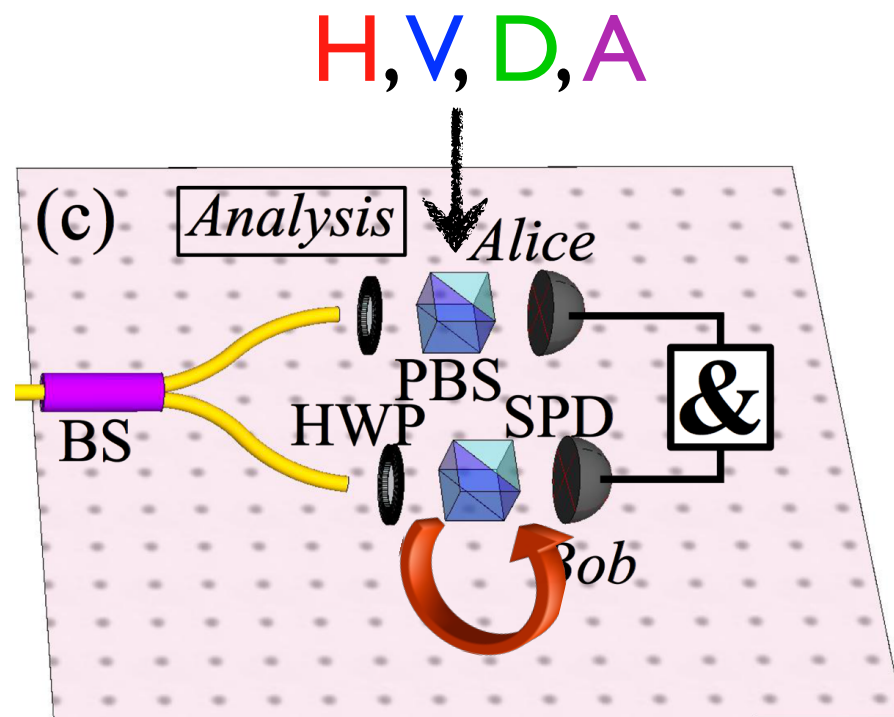
3.3 Bell inequality type mst

▶ Tests with 3 FBG filters



3.3 Bell inequality type mst

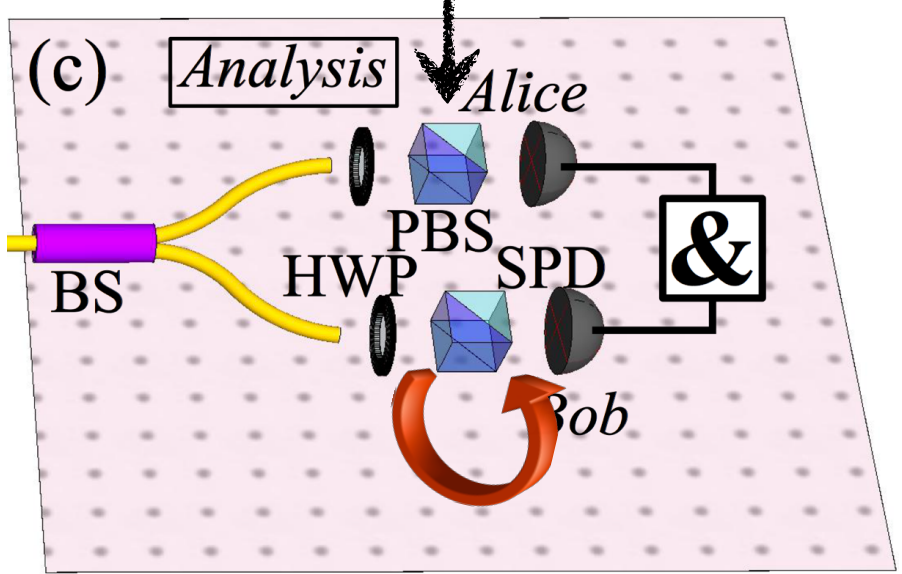
▶ Tests with 3 FBG filters



3.3 Bell inequality type mst

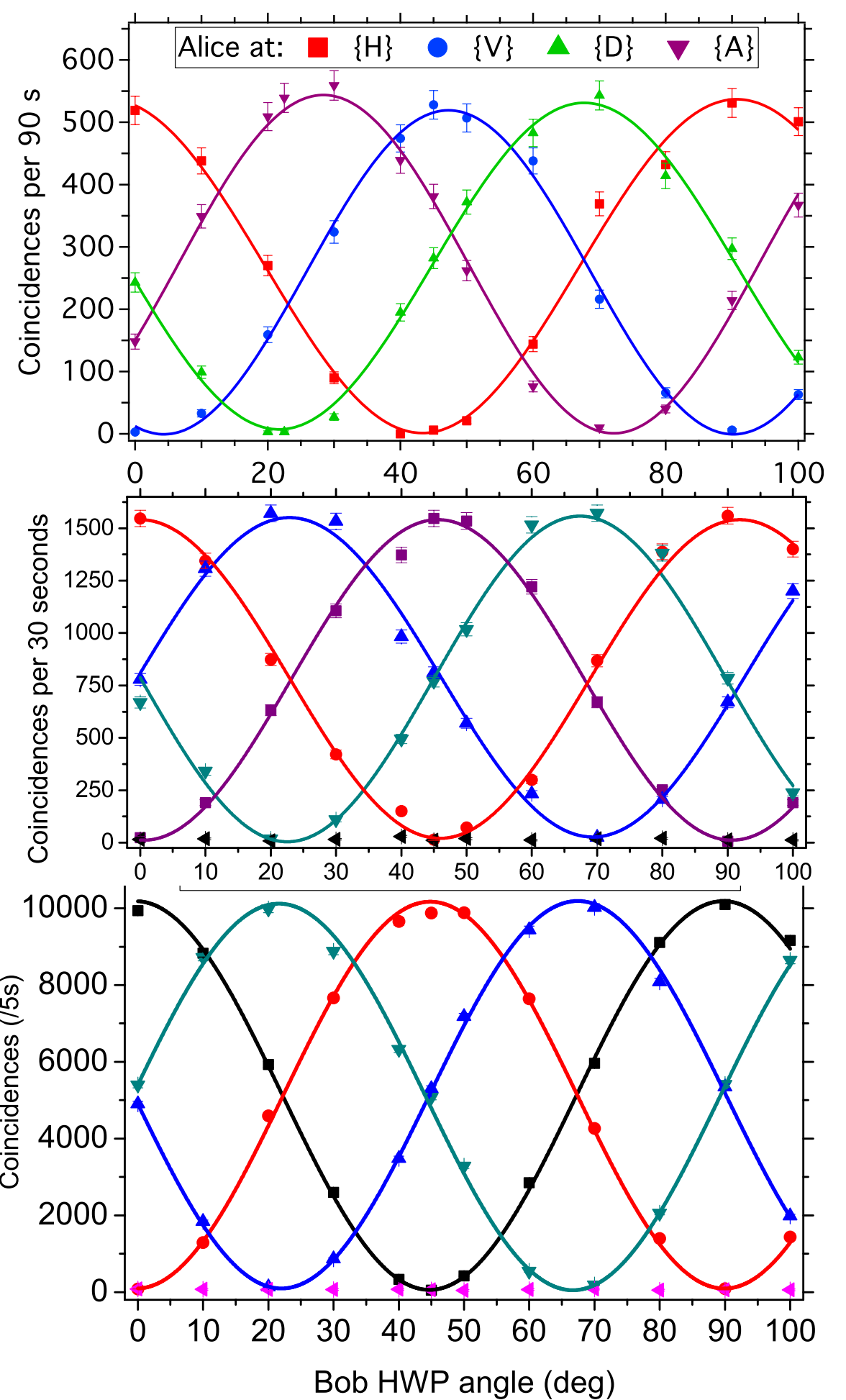
▶ Tests with 3 FBG filters 80 GHz

H, V, D, A



540 MHz

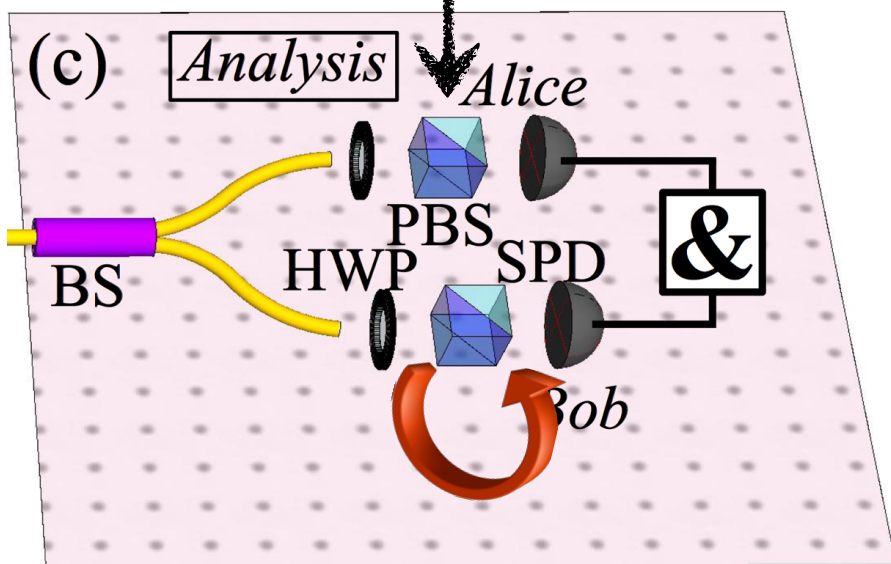
25 MHz



3.3 Bell inequality type mst

▶ Tests with 3 FBG filters 80 GHz

H, V, D, A

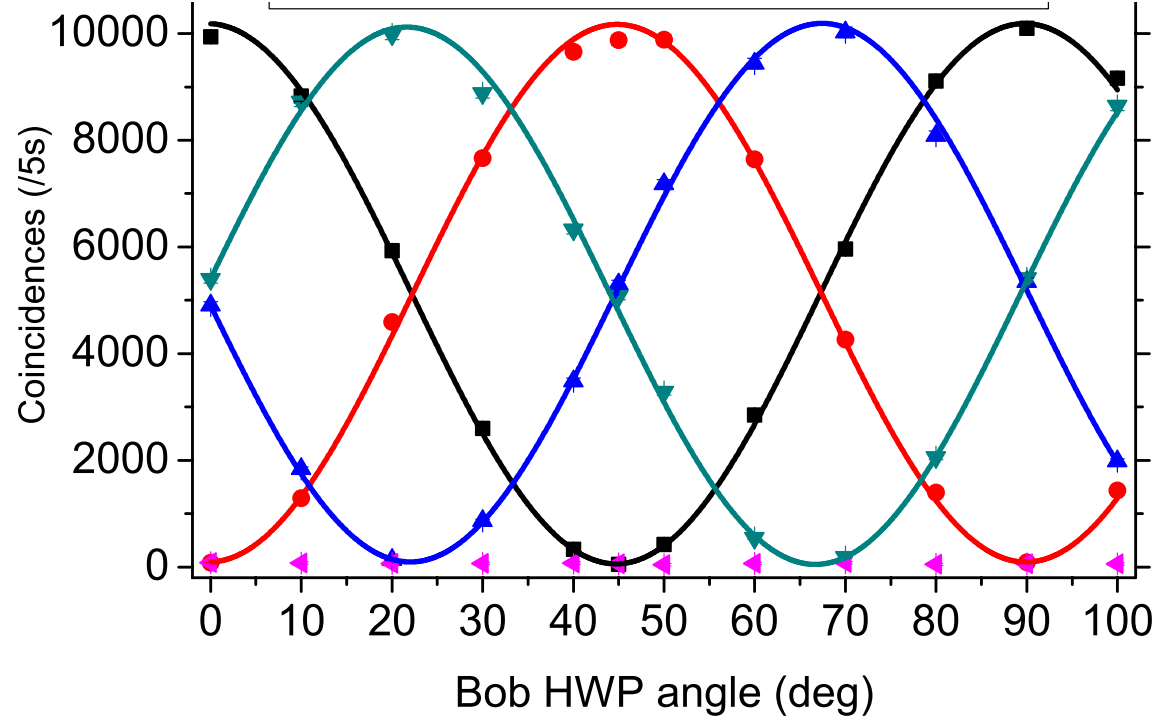
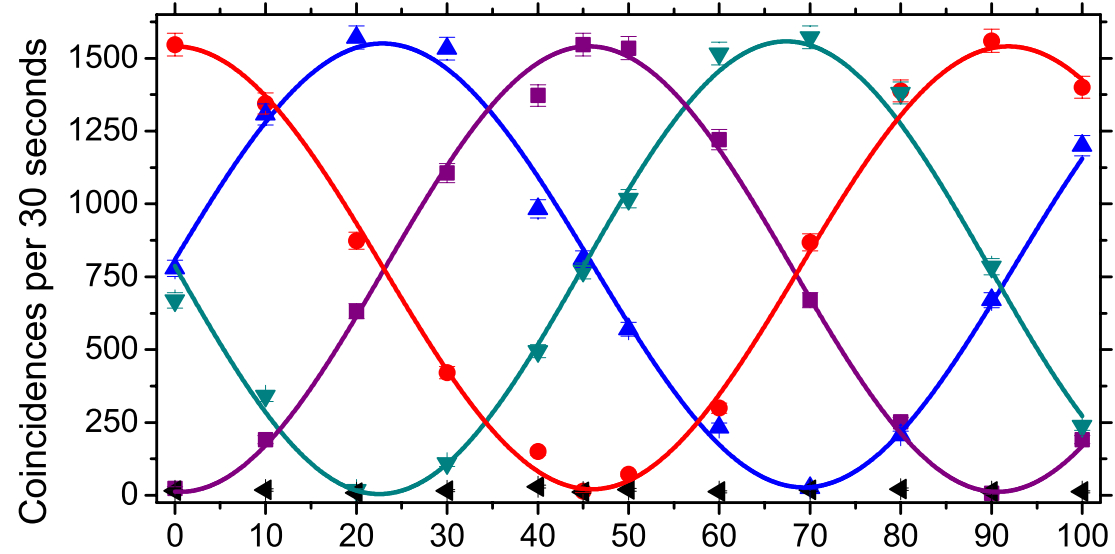
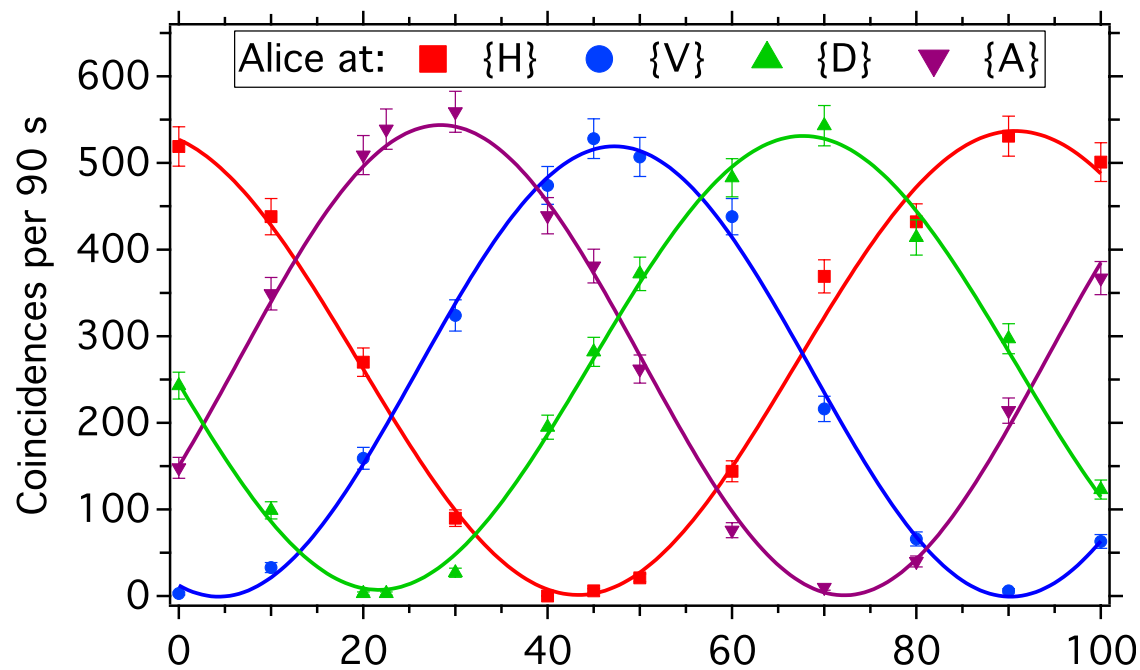


540 MHz

	Standard H,V,D,A
Vraw	> 99% all

25 MHz

Excellent quality entanglement



Conclusion

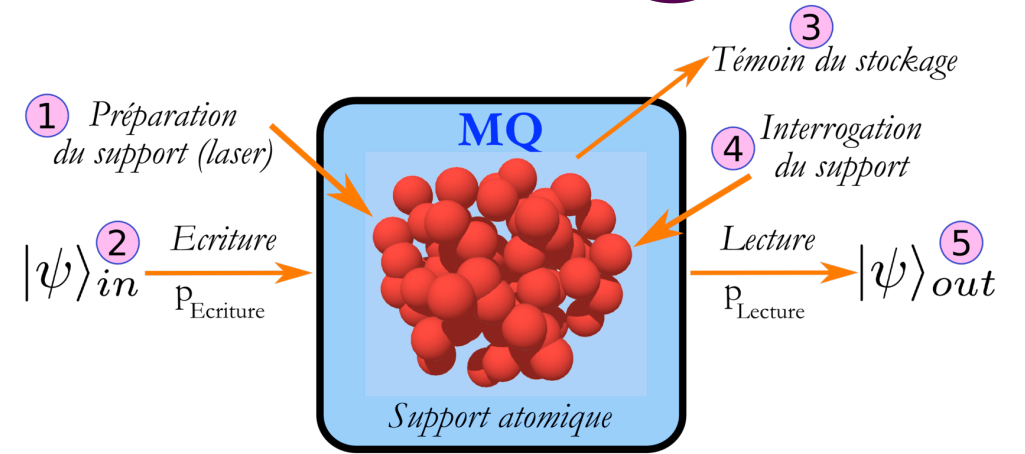
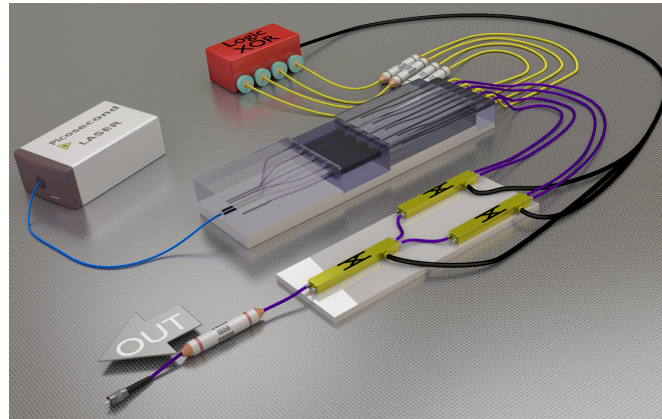
1. **HBT-type setups play an important role in Q optics & com.**
 1. Help characterizing single photon and paired photon sources
 2. Permit connecting heralded SPSs in a quantum network
 3. Offer direct measurement of the photons' coherence time
2. **Particular attention has to be paid to**
 1. Single photon detectors (noise, timing jitter)
 2. HOM effect requires photons in a pure temporal mode
 3. Losses over the quantum channels

L'équipe QILM + OptiNiL

UMR 7336



Laboratoire Physique de la Matière Condensée



Laurent Labonté Amandine Issautier Virginia D'Auria Djeylan Aktas

Lutfi Ngah



Marc De Micheli

Florent Doutre

Olivier Alibart

ST

Anders Kastberg